

# Constraining Dark Matter with Background Light

Sam McDermott

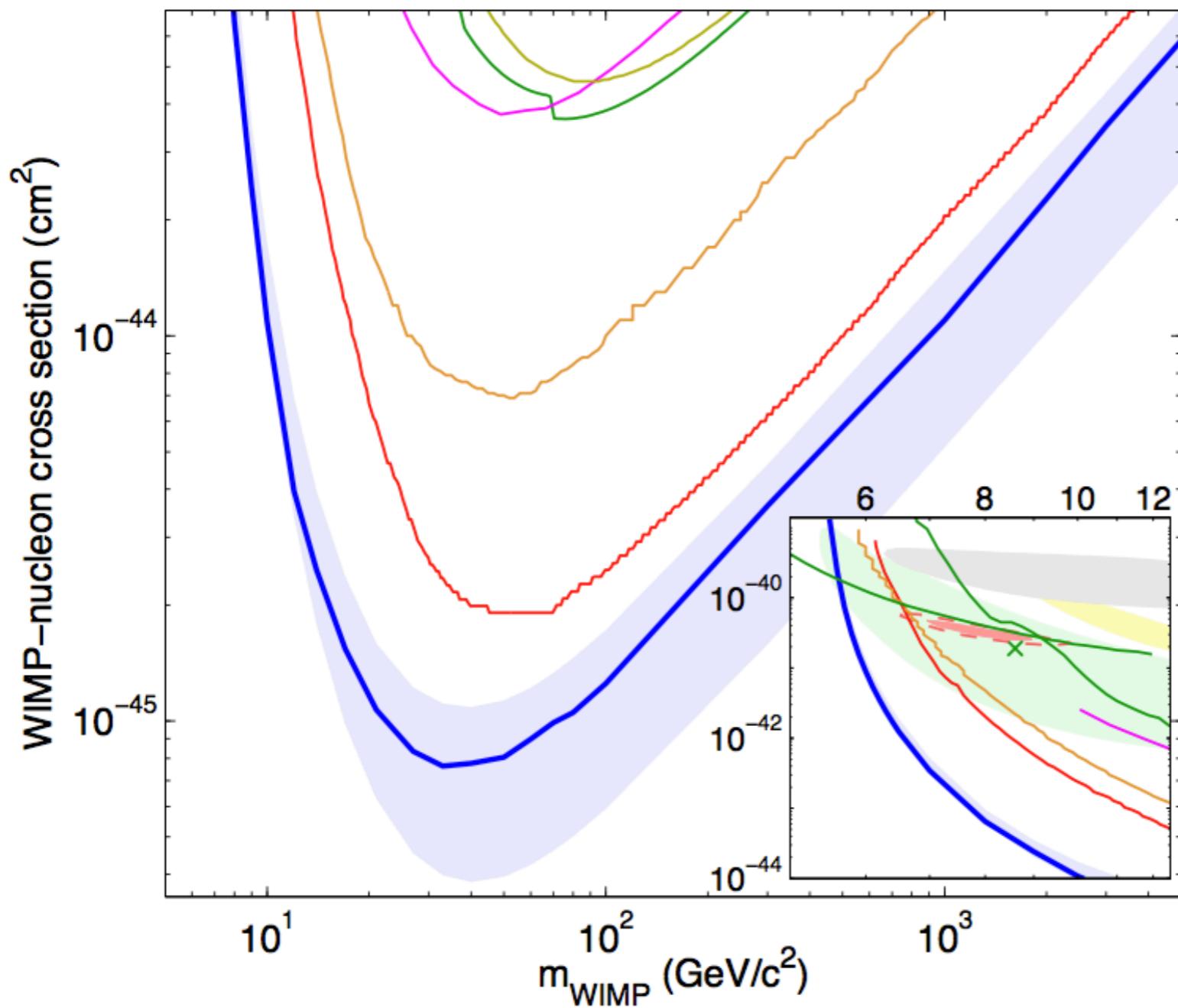
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LANL



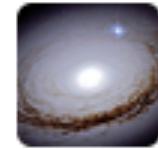
from 1309.4091,  
with Rouven Essig, Eric Kuflik, Tomer Volansky, and Kathryn Zurek

and 1312.0608  
with Ilias Cholis and Dan Hooper

# Prelude



LUX  
sees  
nothing



Matthew Buckley @physicsmatt

#LUXdm 20 times better than previous experiments. And light dark matter just got dragged out back and beaten with a stick.

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30 Oct

Is this really a big problem? See e.g.:  
Cirigliano, Graesser, Ovanesyan, and  
Shoemaker 1311.5886;  
Gresham and Zurek 1311.2082

# Motivation

- More and more of our favorite DM parameter space seems to be getting ruled out
- (Infinitely) large swaths of well-motivated DM parameter space are currently up for grabs
- We need complementary probes to
  - test assumptions of bounds
  - explore different parameter space
- This talk will focus on indirect detection

# Motivation

- Photons (produced directly, from FSR, from cascades, etc.) are generic in DM annihilation and decay
- Data are “just sitting there” ready to use

This talk

two very different methods  
of approaching this problem

# Outline

- Light decaying dark matter:
  - models
  - statistical methodology
  - results
- “Light” annihilating dark matter:
  - isotropic gamma rays from DM
  - astrophysical backgrounds
  - results

Part I:

Light (MeV-ish)  
Decaying DM

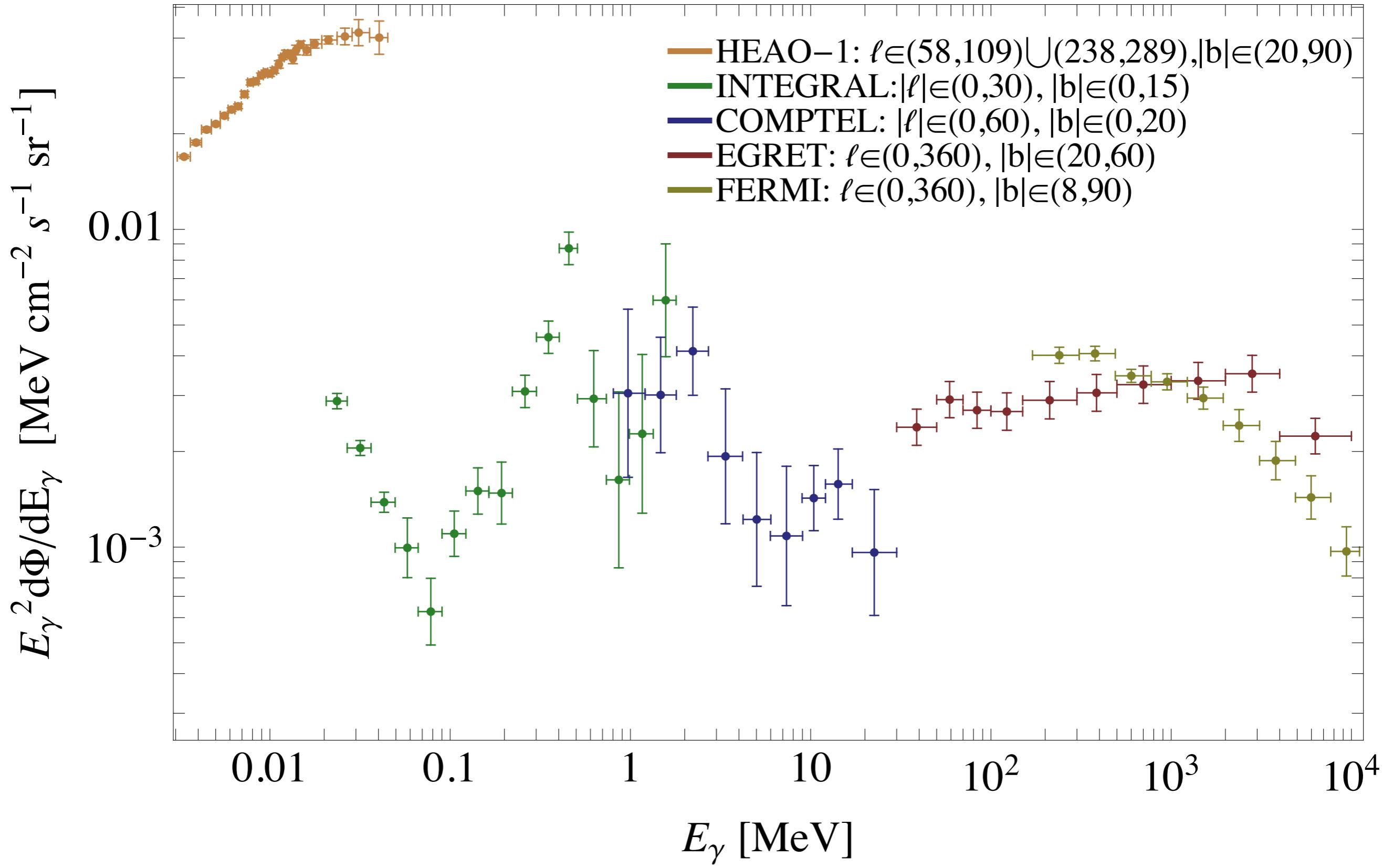
# “Light” Dark Matter

- still cold
  - not warm, not an ALP that forms a galactic scale BEC, etc.
  - few keV to few GeV mass
- we assume that all (model-dependent) observables of standard cosmology are taken care of (i.e., asymmetric or thermal production where appropriate)

# Data

- “Diffuse” X-Rays and Gamma-Rays
  - HEAO-1 (1977)
  - INTEGRAL (2008)
  - COMPTEL (1998)
  - EGRET (2003)
  - Fermi (2012) (21 months)
- Some observations near the galactic poles, some near the galactic center

# Data



# Analysis methods

- There are many different methods:
  - spectral fit plus power law in sliding energy window (cf. Weniger)
  - precise background modeling (cf. Siegal-Gaskins; newer work here; others)
  - “on-off” or template analyses (cf. Koushiappas + Geringer-Sameth; Finkbeiner + Slatyer; Hooper + Slatyer; Tempel, Hektor, + Raidal; others)

# Analysis methods

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  - requires exceptional energy resolution
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# Analysis methods

- There are many different methods:
  - spectral fit plus power law in ~~solid state~~ energy window (cf. Weniger)
  - precise background modeling (cf. Siegal-Gaskins; newer work here; others)
  - “on-off” or template analyses (cf. Koushiappas + Geringer-Sameth; Finn + Slatyer; Hooper + Slatyer; Tempel, Hektor, Raidal; others)

requires detailed knowledge of astrophysics over very wide ranges of position and energy space

# Analysis methods

- There are many different methods:
    - spectral fit plus power law in sliding energy window (cf. Weniger)
    - precise background modeling (cf. Siegal-Gaskins; newer work here; others)
    - “on-off” or template analyses (cf. Koushiappas + Geringer-Sameth; Finkbeiner + Slatyer; Hooper + Slatyer; Tempel, Hektor, + Raidal; others)
- requires exceptional angular resolution

# Analysis methods

- There are many different methods:
  - ~~spectral fit plus power law in sliding energy window (cf. Weniger)~~
  - ~~precise background modeling (cf. Siegal-Gaskins; newer work here; others)~~
  - ~~“on-off” or template analyses (cf. Koushiappas + Geringer-Sameth; Finkbeiner + Slatyer; Hooper + Slatyer; Tempel, Hektor, + Raidal; others)~~

# Our analysis (1309.4091)

- Only direct photon production and primary FSR; no ICS, synchrotron, etc.
- We simply require (for every energy bin):  
$$\text{Flux}_{\text{predicted}} \leq \text{Flux}_{\text{observed}} + 2 \times \text{Error Bar}_{\text{observed}}$$
- Robust results!

# Our analysis (1309.4091)

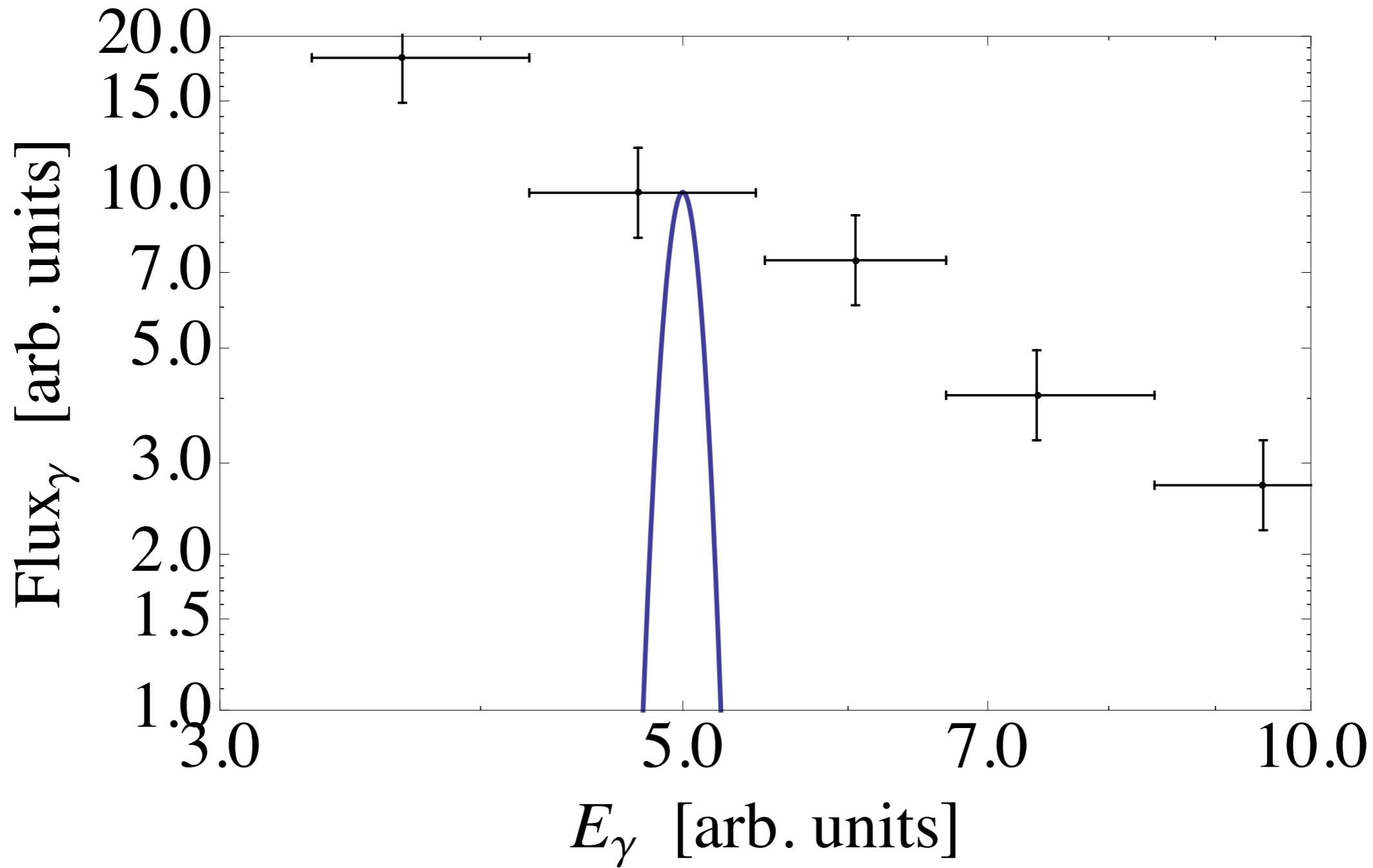
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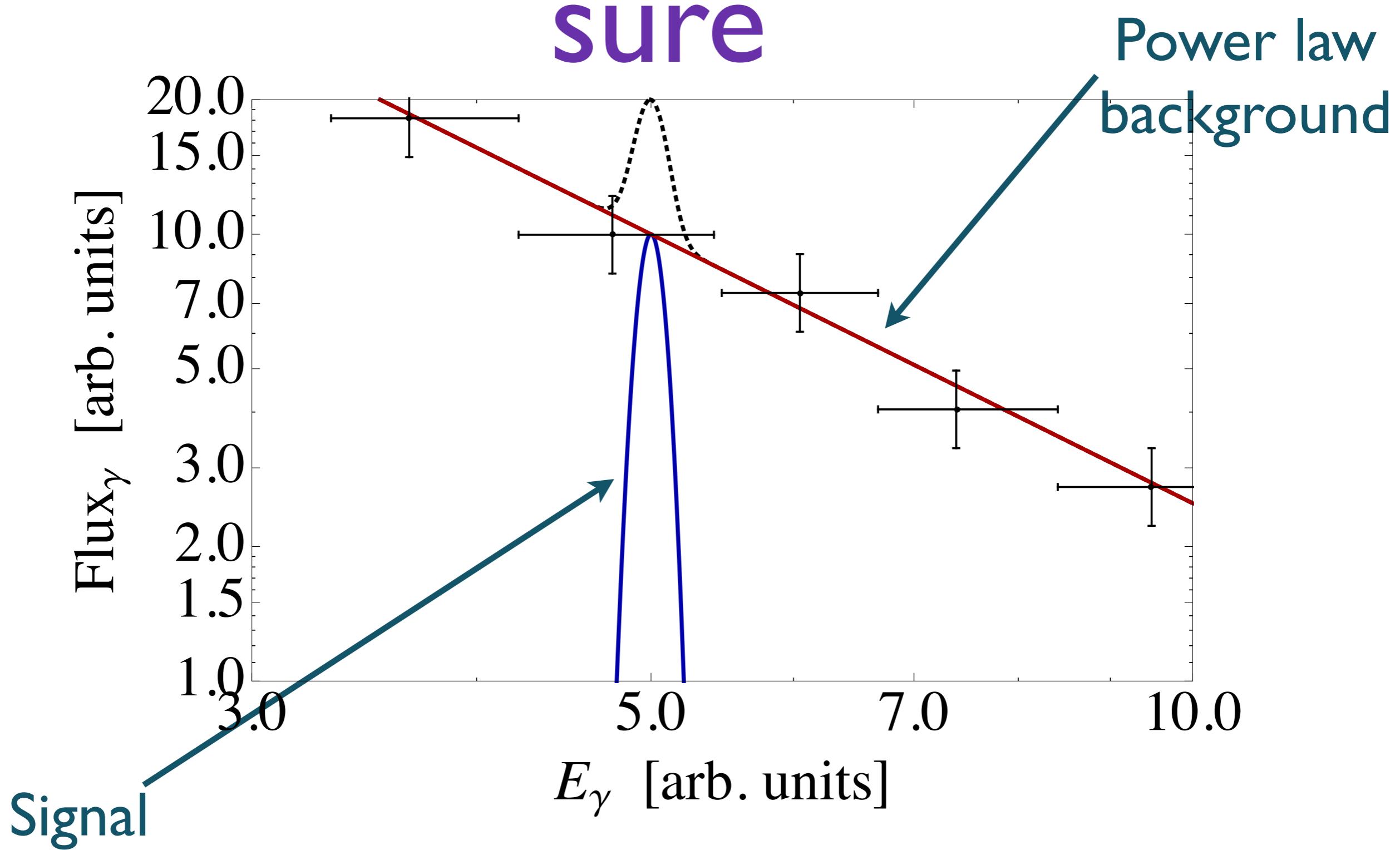


# Can we exclude this?

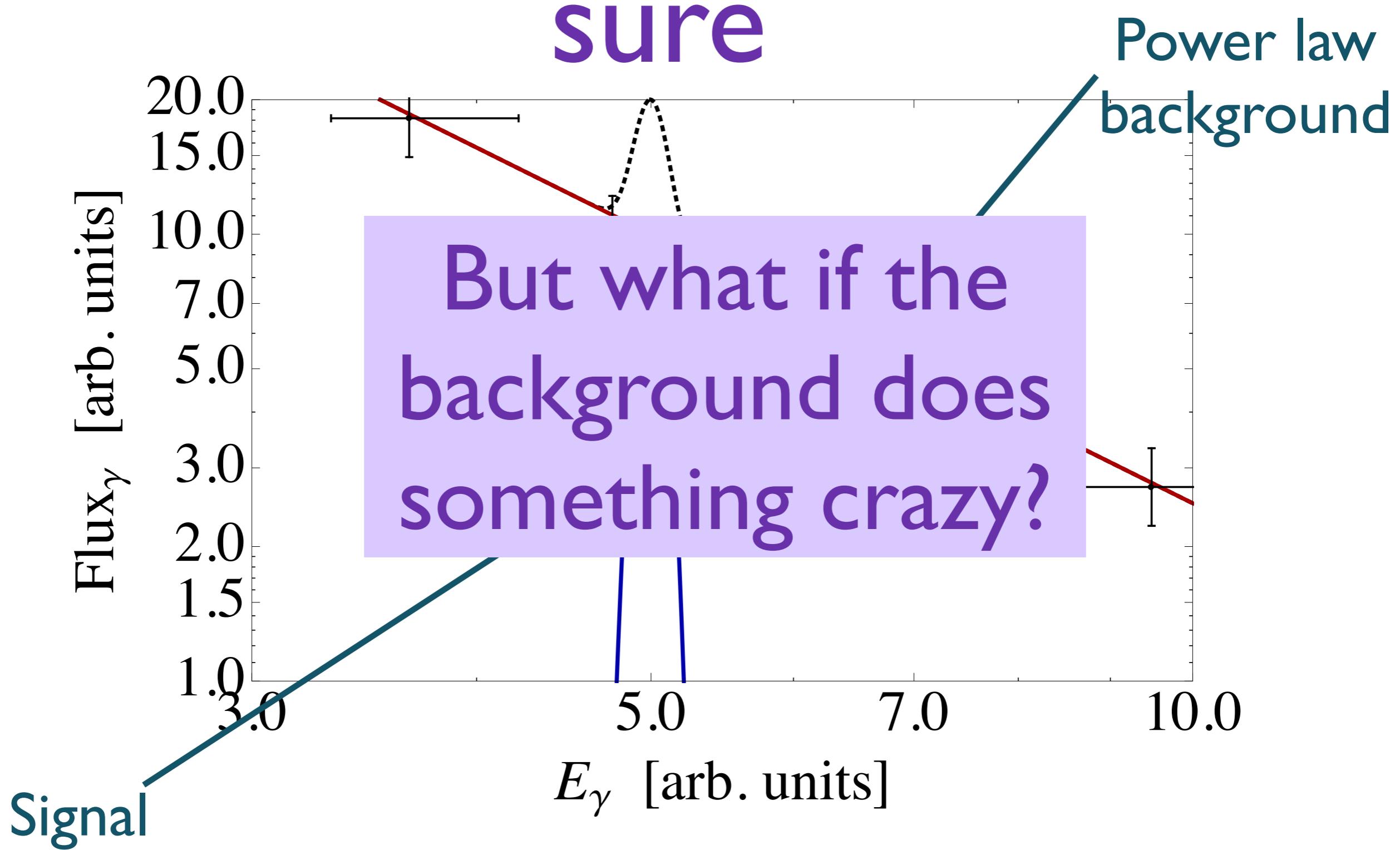


# With background model,

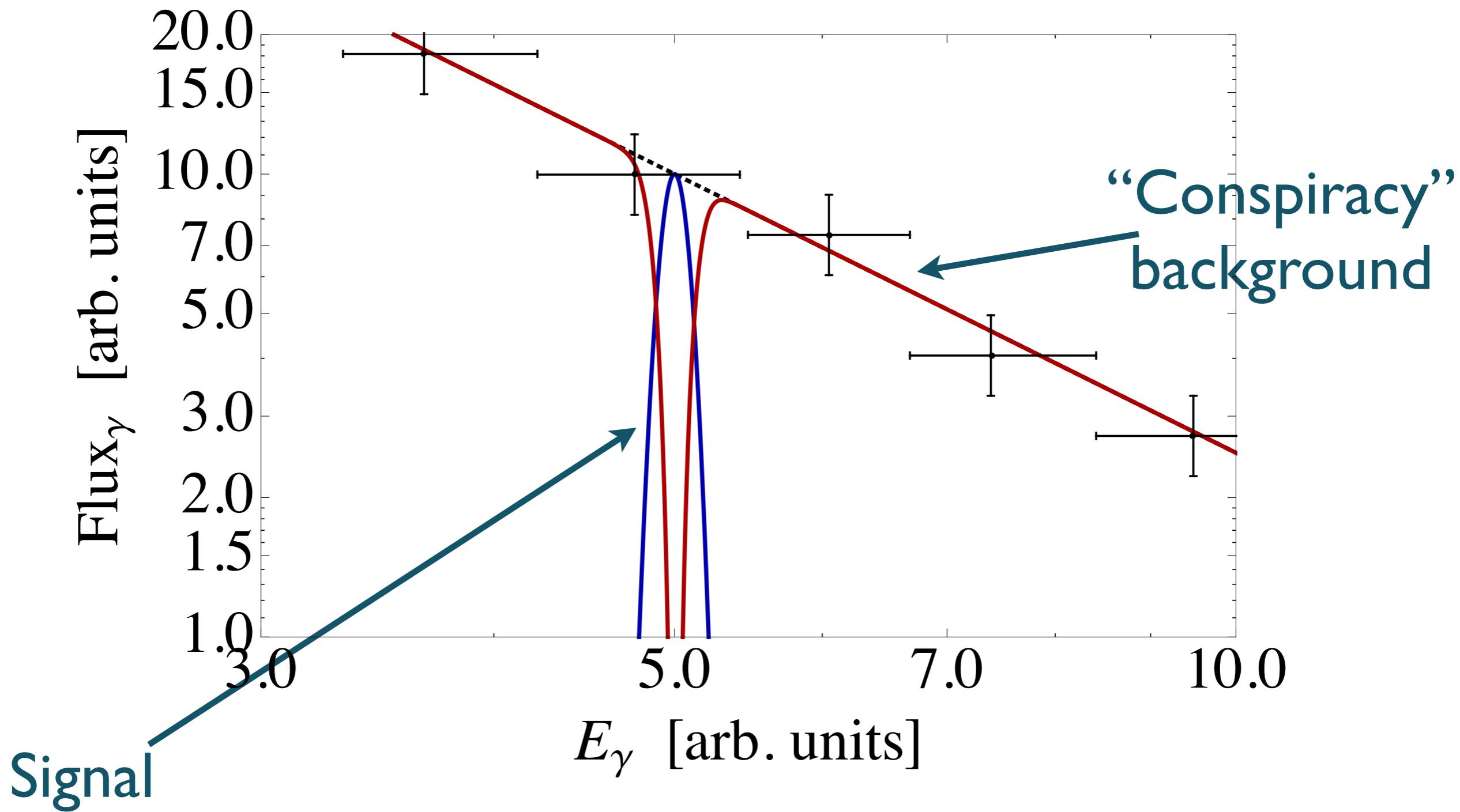
## sure



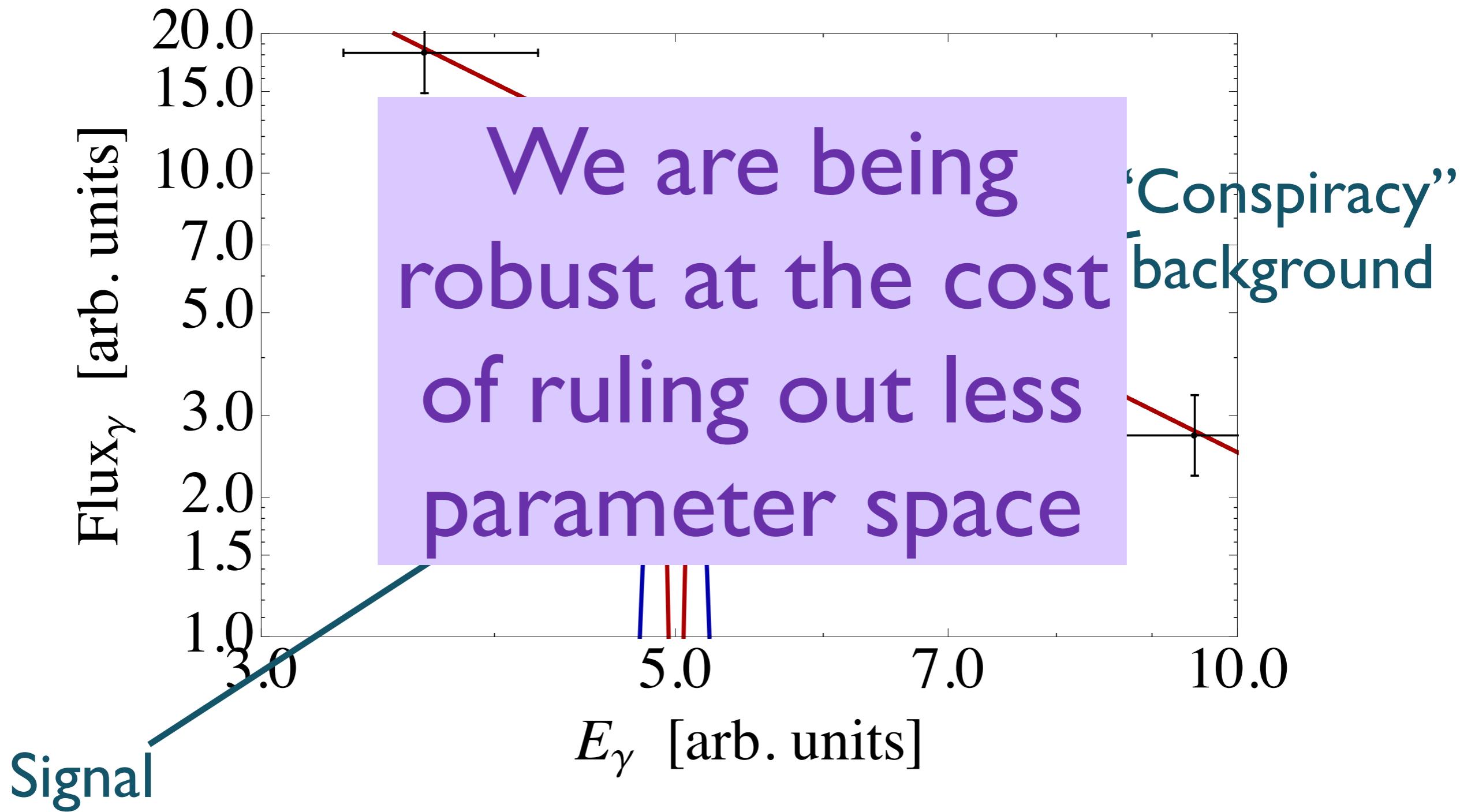
# With background model, sure



# Being conservative: Not ruled out!



# Being conservative: Not ruled out!



# Photons from DM decays

Galactic:  $\frac{d\Phi_{\gamma,G}}{dE} = \frac{r_\odot}{4\pi} \frac{\rho_\odot}{m_{\text{DM}}} \Gamma \frac{dN_\gamma}{dE} J(\Omega)$

vs. extragalactic:

$$\frac{d\Phi_{\gamma,EG}}{dE} = \frac{\Omega}{4\pi} \frac{\Gamma \Omega_{\text{DM}} \rho_c}{m_{\text{DM}} a_0 H_0} \int_0^\infty dz \frac{dN}{dE(z)} \frac{1}{\sqrt{\Omega_\Lambda + \Omega_m(1+z)^3}}$$

Galactic dominates, marginally:

$$\rho_\odot r_\odot J(\Omega) \simeq \mathcal{O}(10^{-5} \text{ GeV}^3) \quad \text{vs.} \quad \rho_{\text{DM}}/H_0 \simeq 5 \times 10^{-6} \text{ GeV}^3$$

but same order of magnitude

# Models of decaying DM

- **Hidden Photino** – SUSY + hidden U(1). [Higgs a hidden U(1); break SUSY with messengers from SM; hidden photon/photino with small mass splitting; kinetically mix hidden photon and photon of U(1)EM. Decays involving SM possible if there is a light gravitino. Topology depends on relative masses of hidden photon/photino.]
- **Sterile Neutrino** – long-lived sterile neutrino
- **RPV Gravitino** – neutrino/photino mixing. [Planck-scale suppression gives a naturally small rate for gravitino decays. Fastest decay is gravitino > photon + neutrino.]
- Plus dipole DM, dark scalar, dark pseudoscalar

# So the name of the game is...

Ignore the astro/cosmology stuff (be conservative)

Particle physics enters through  $\Gamma$  and  $dN/dE_\gamma$  only:

$dN/dE_\gamma$  is fixed by decay topology

$\Gamma$  is fixed by the model

# Dark photino model

MeV scale comes out naturally:

$$m_{\tilde{\gamma}_d}^2 = \epsilon \cdot g_d \langle D_Y \rangle \simeq (5 \text{ MeV})^2 \left( \frac{\epsilon}{10^{-8}} \right) \left( \frac{g_d}{0.2} \right) \left( \frac{\sqrt{\langle D_Y \rangle}}{50 \text{ GeV}} \right)^2$$

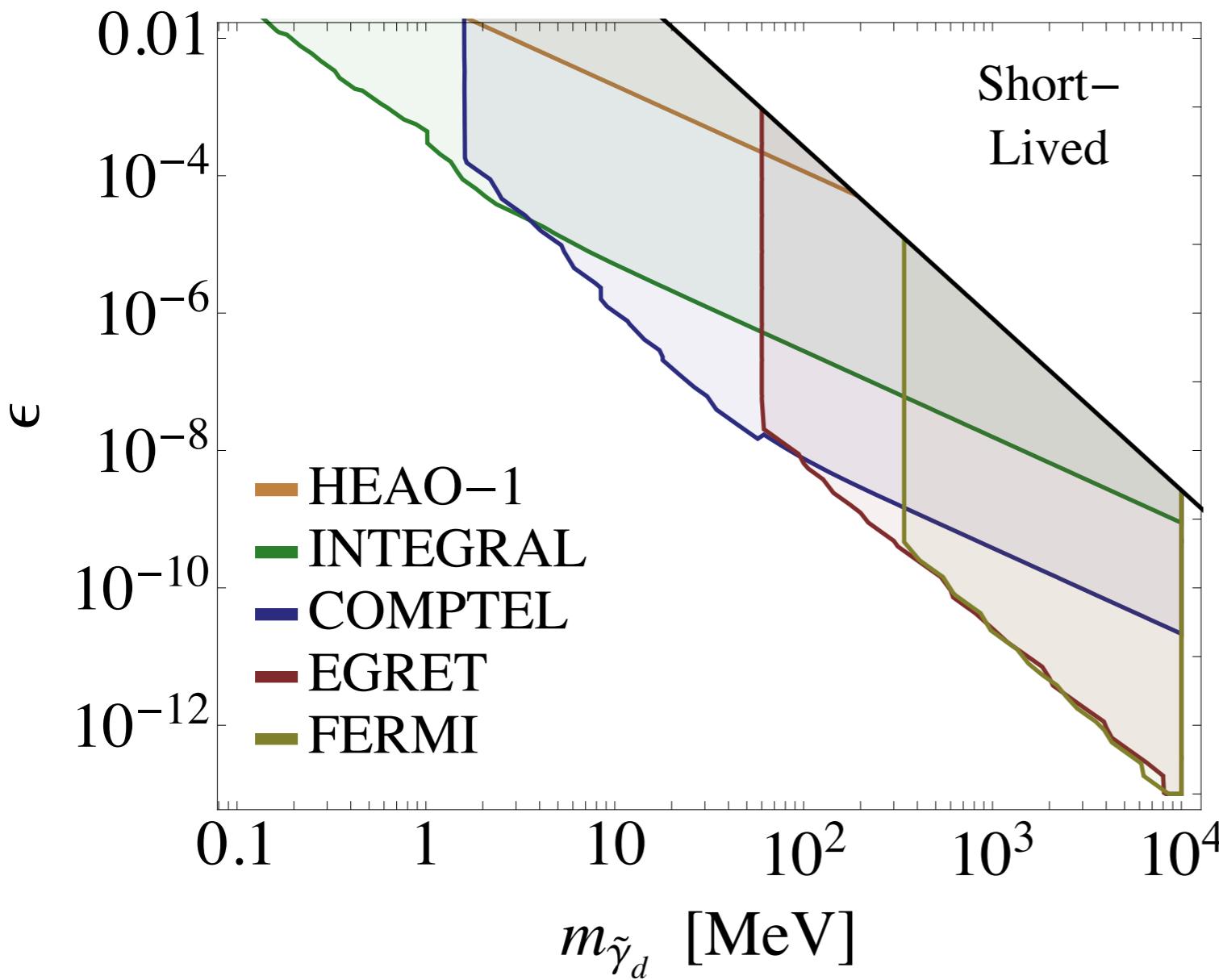
Two different decays, depending on whether the dark photon or dark photino is more massive:

$$\tau_{\tilde{\gamma}_d \rightarrow \gamma \tilde{G}} \simeq 3 \times 10^{23} \text{ sec} \left( \frac{10^{-8}}{\epsilon} \right)^2 \left( \frac{10 \text{ MeV}}{m_{\tilde{\gamma}_d}} \right)^5 \left( \frac{\sqrt{F}}{100 \text{ TeV}} \right)^4$$

$$\tau_{\tilde{\gamma}_d \rightarrow \gamma_d \tilde{G}} \simeq 3 \times 10^{20} \text{ sec} \left( \frac{1 \text{ MeV}}{m_{\tilde{\gamma}_d}} \right)^5 \left( \frac{\sqrt{F}}{10^4 \text{ TeV}} \right)^4 \left( 1 - \frac{m_{\gamma_d}^2}{m_{\text{DM}}^2} \right)^{-4}$$

# Dark photino DM, $m_{\tilde{\gamma}_d} < m_{\gamma_d}$

$$\tau_{\tilde{\gamma}_d \rightarrow \gamma \tilde{G}} \simeq 3 \times 10^{23} \text{ sec} \left( \frac{10^{-8}}{\epsilon} \right)^2 \left( \frac{10 \text{ MeV}}{m_{\tilde{\gamma}_d}} \right)^5 \left( \frac{\sqrt{F}}{100 \text{ TeV}} \right)^4$$

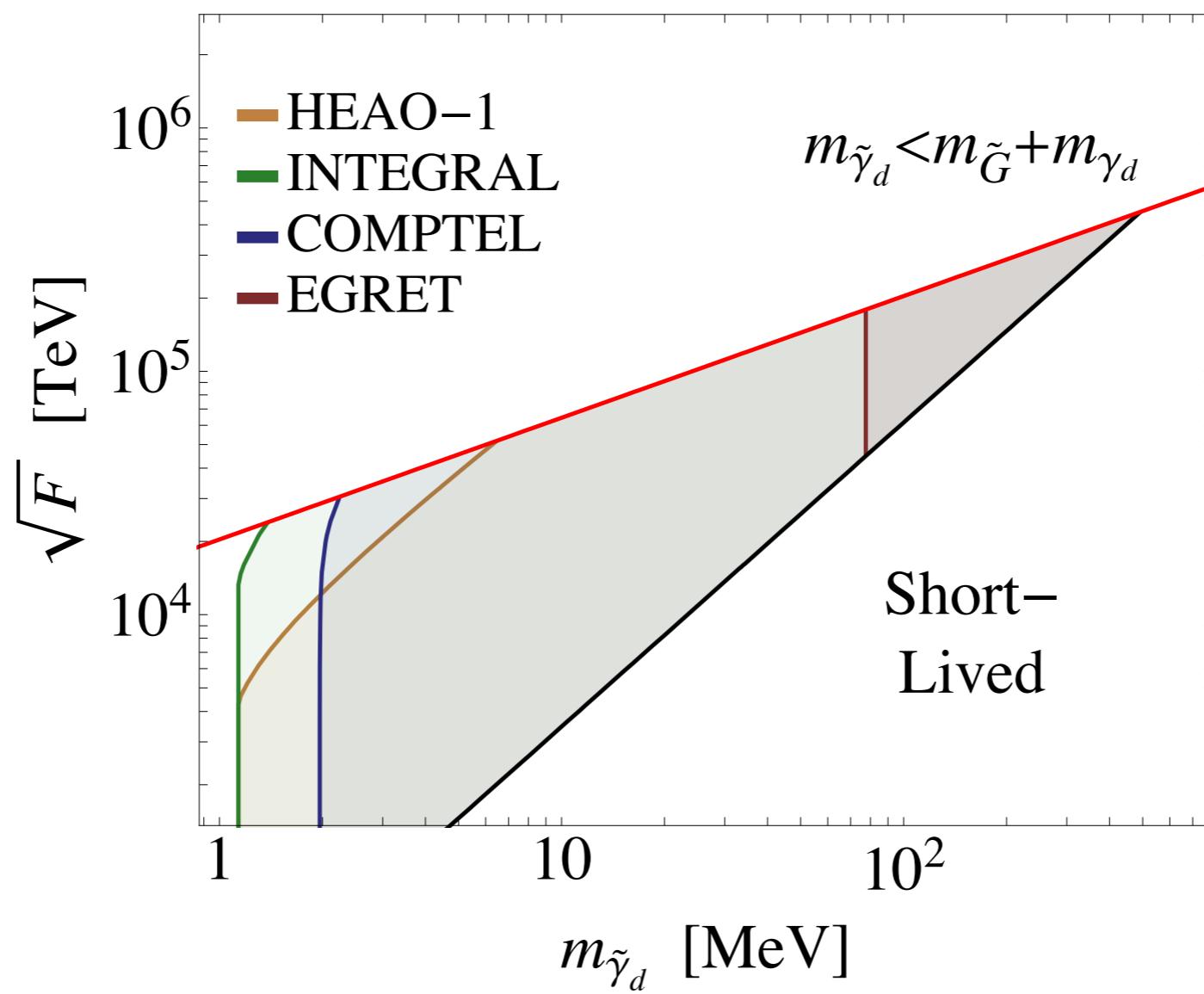


$$\sqrt{F} = 10^4 \text{ TeV}$$

# Dark photino DM, $m_{\tilde{\gamma}_d} > m_{\gamma_d}$

$$\tau_{\tilde{\gamma}_d \rightarrow \gamma_d \tilde{G}} \simeq 3 \times 10^{20} \text{ sec} \left( \frac{1 \text{ MeV}}{m_{\tilde{\gamma}_d}} \right)^5 \left( \frac{\sqrt{F}}{10^4 \text{ TeV}} \right)^4 \left( 1 - \frac{m_{\gamma_d}^2}{m_{\text{DM}}^2} \right)^{-4}$$

$\tilde{\gamma}_d \rightarrow \gamma_d \tilde{G} \rightarrow f^+ f^- \tilde{G}$



# Sterile neutrino DM

Three-body and radiative decays contribute to photon background at similar levels:

mixing angle between  $\nu_s$  and  $\nu_e$

$$\tau_{\nu_s \rightarrow \nu \gamma} \simeq 7.2 \times 10^{17} \text{ sec} \left( \frac{10 \text{ MeV}}{m_\chi} \right)^5 \left( \frac{\sin^2 2\theta}{10^{-8}} \right)^{-1}$$

$$\tau_{\nu_s \rightarrow \nu_\alpha e^+ e^-} \simeq 9.6 \times 10^{15} \text{ sec} \left( \frac{10 \text{ MeV}}{m_\chi} \right)^5 \left( \frac{\sin^2 2\theta}{10^{-8}} \right)^{-1}$$

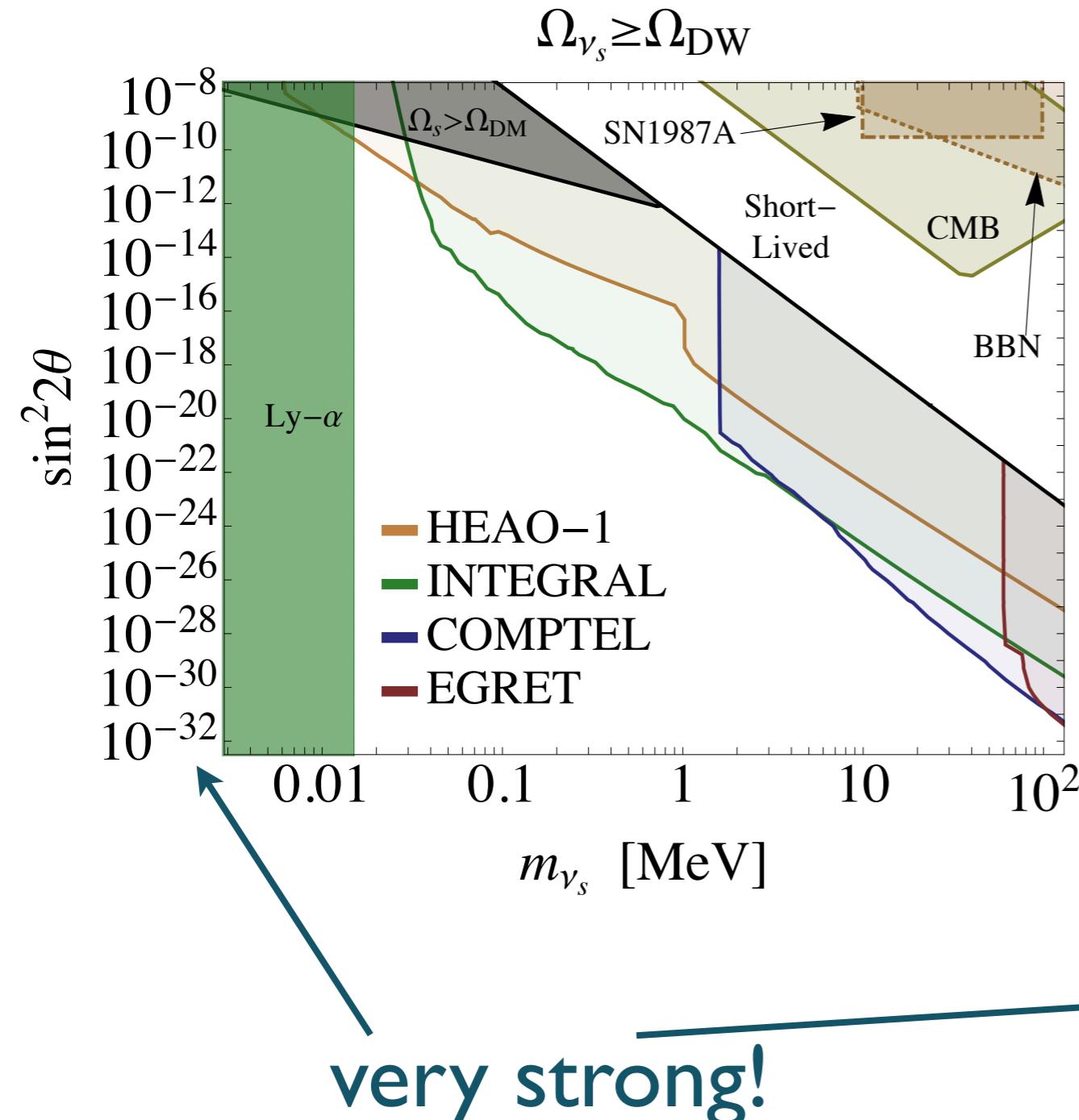
# Sterile neutrino DM

Relic abundance is model-dependent,  
but a UV-insensitive contribution  
comes from late-time oscillations  
(Dodelson-Widrow mechanism)

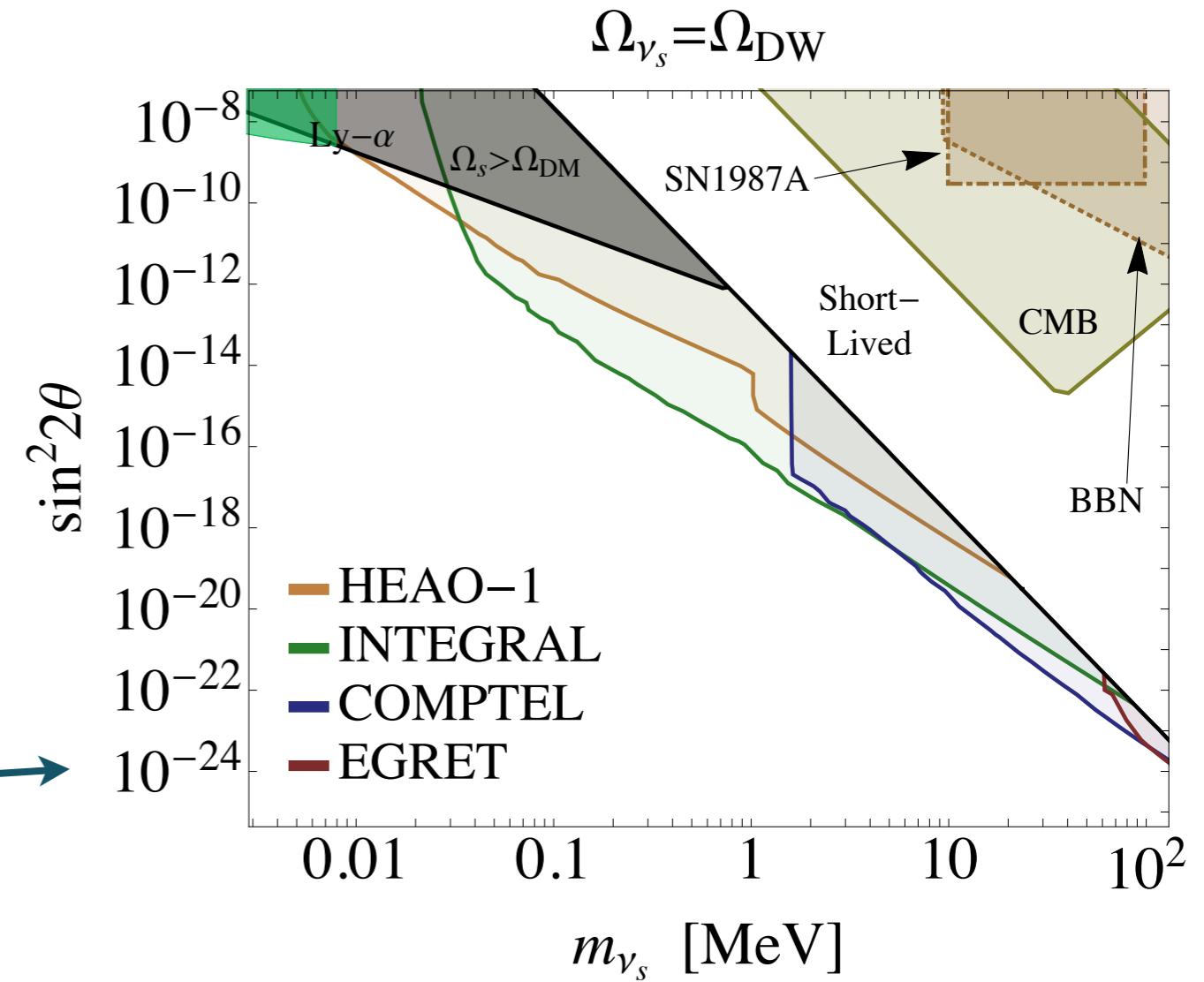
Bounds are different depending on  
whether  $\Omega_{\nu_s} = \Omega_{\text{DW}}$  or  $\Omega_{\nu_s} \geq \Omega_{\text{DW}}$

# Sterile neutrino DM

$\nu_s$  (radiative and three-body decays)

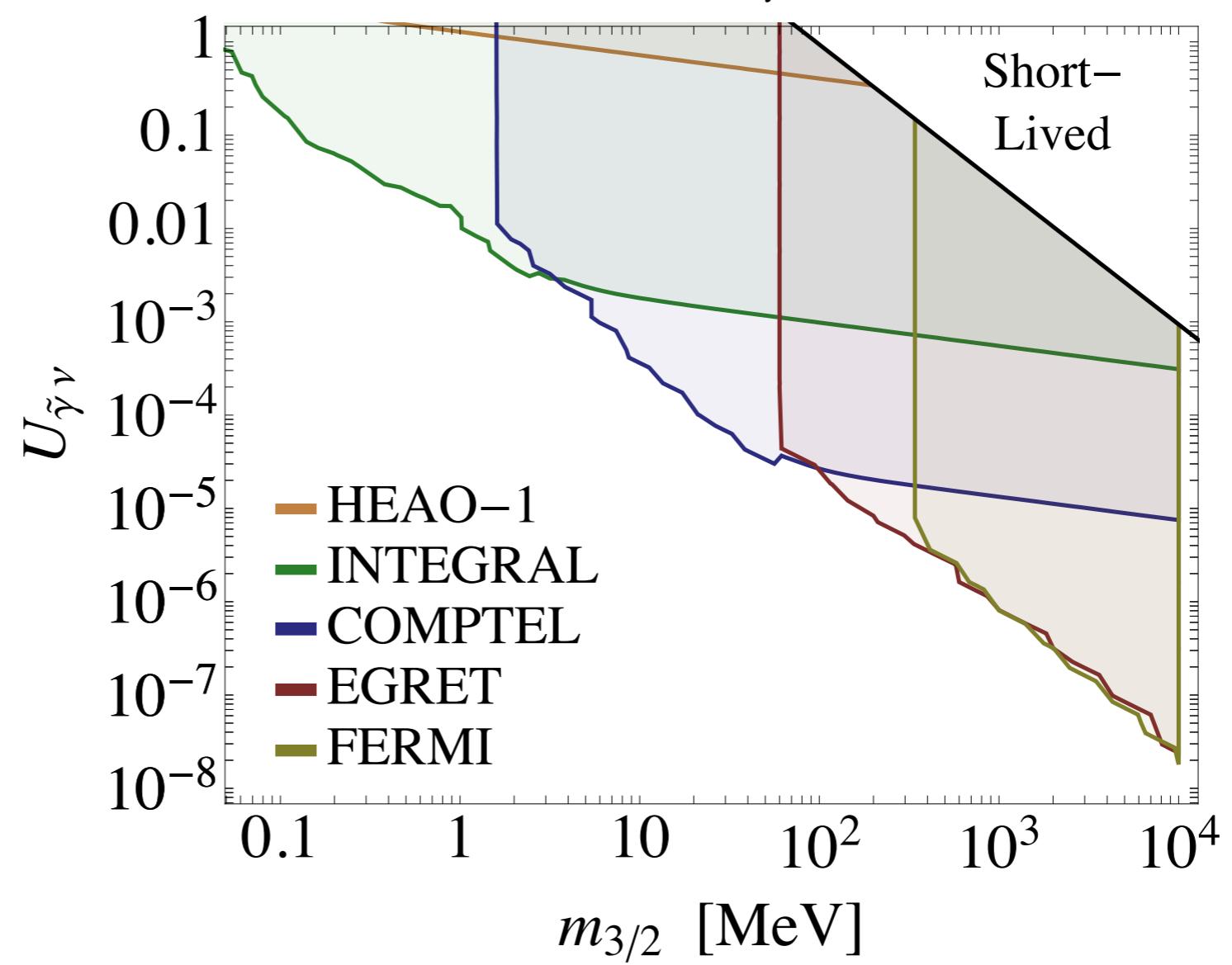
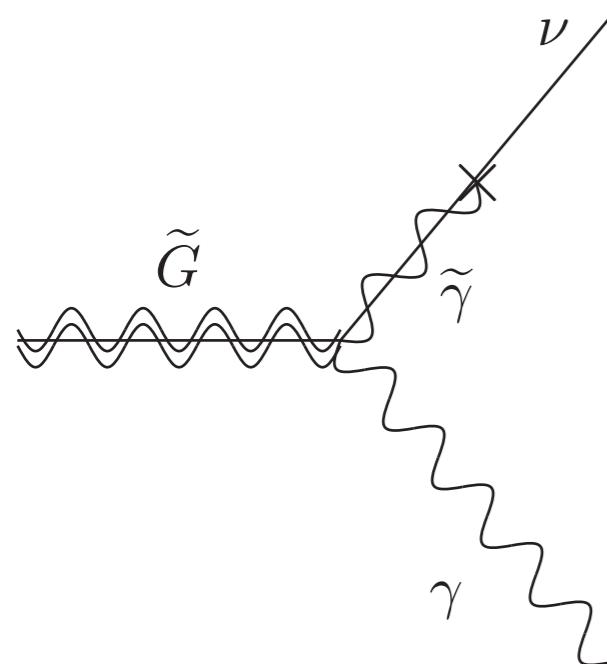


$\nu_s$  (radiative and three-body decays)



# RPV gravitino DM

$$\tau_{\tilde{G} \rightarrow \nu \gamma} = \left( \frac{1}{32\pi} |U_{\tilde{\gamma}\nu}|^2 \frac{m_{3/2}^3}{m_{\text{Pl}}^2} \right)^{-1} \simeq 3.8 \times 10^{28} \text{ sec} \left( \frac{10 \text{ MeV}}{m_{3/2}} \right)^3 \left( \frac{10^{-4}}{|U_{\tilde{\gamma}\nu}|} \right)^2$$



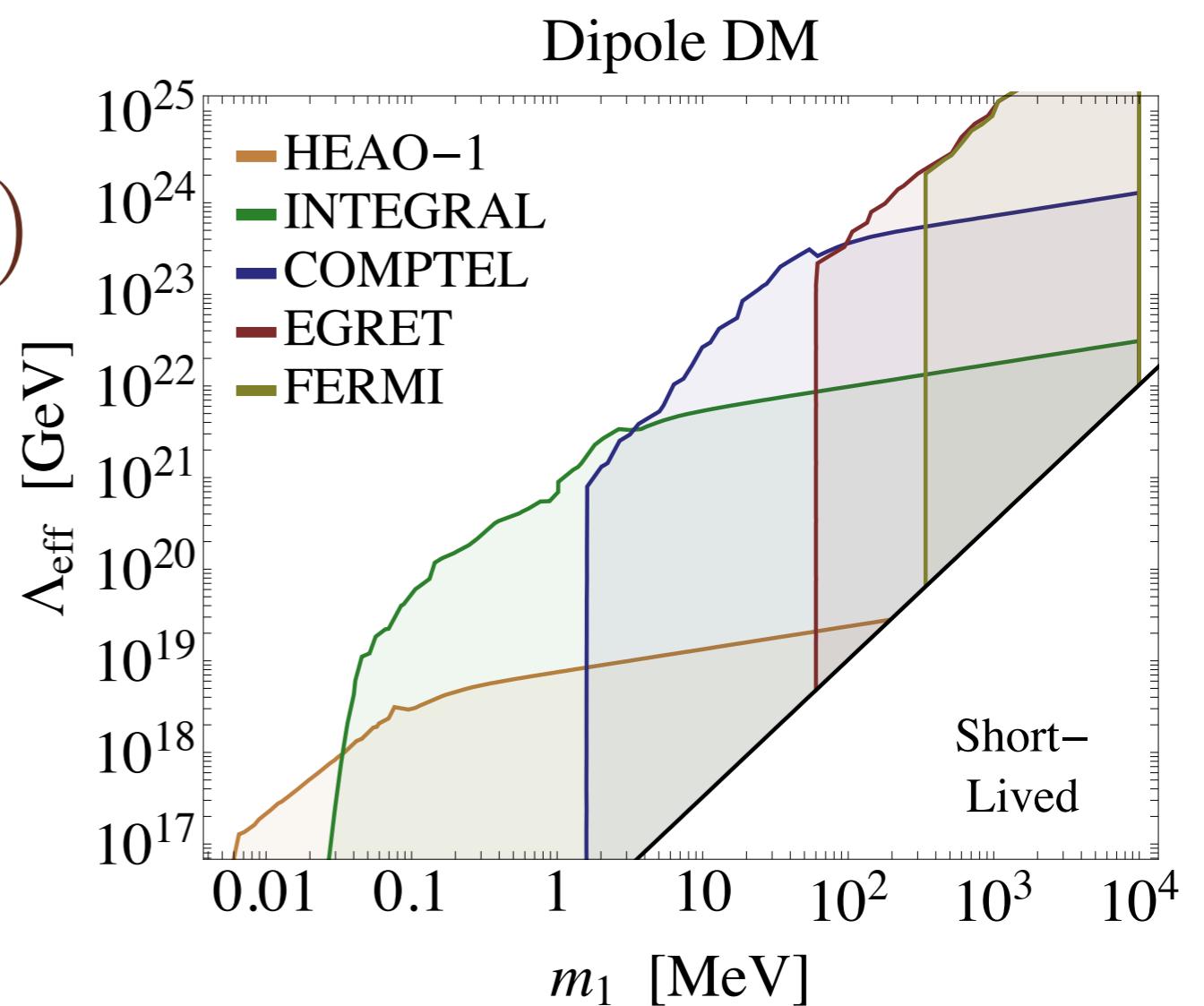
# Dipole DM

$$\mathcal{L} \supset \frac{\lambda}{\Lambda} \bar{\chi}_2 \sigma^{\mu\nu} \chi_1 F_{\mu\nu}$$

$$\tau_{\text{dipole}} \simeq 4.1 \times 10^{20} \text{ sec} \left( \frac{10 \text{ MeV}}{m_1} \right)^3 \left( \frac{\Lambda_{\text{eff}}}{10^{19} \text{ GeV}} \right)^2$$

$$(\Lambda_{\text{eff}} \equiv \Lambda/\lambda)$$

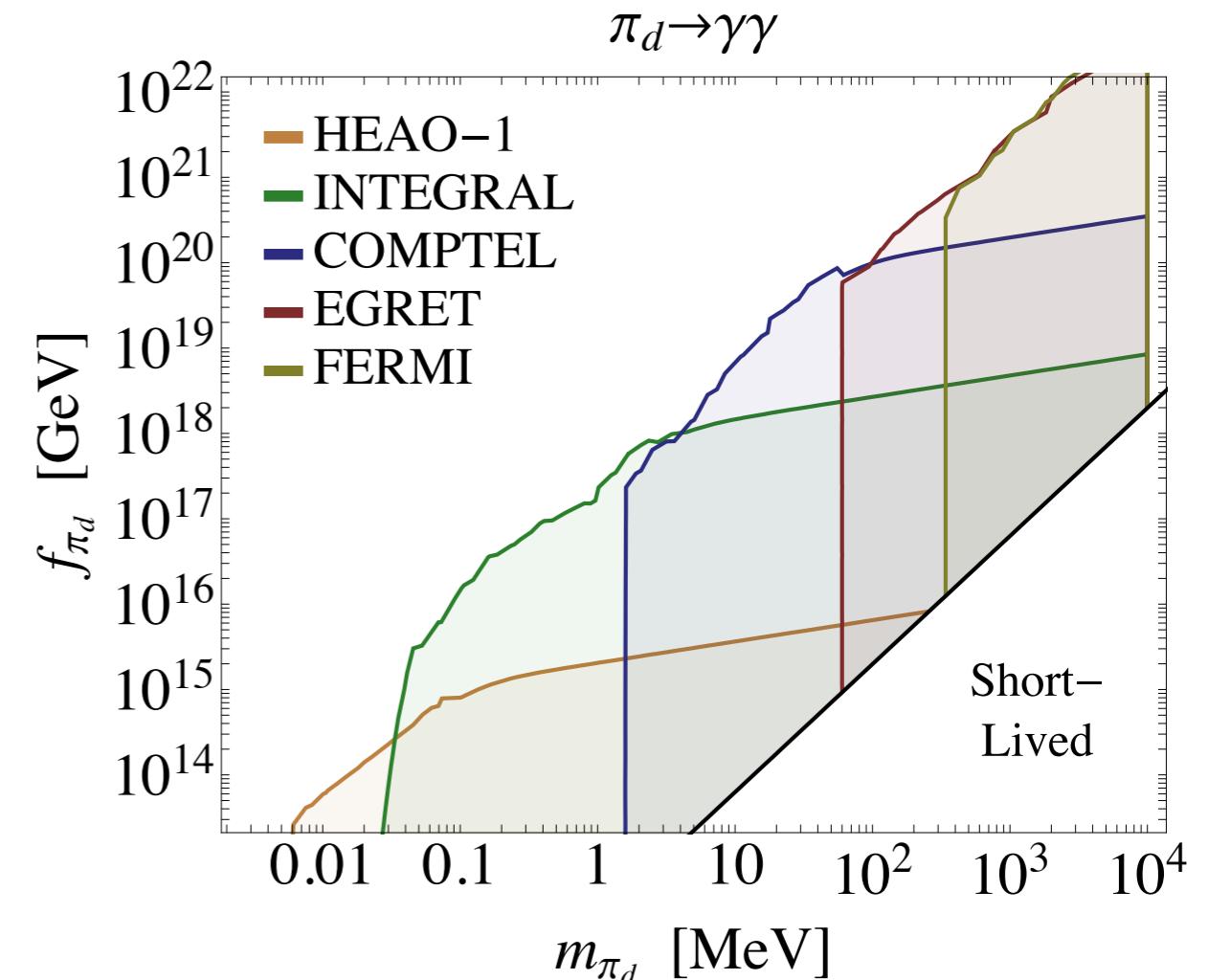
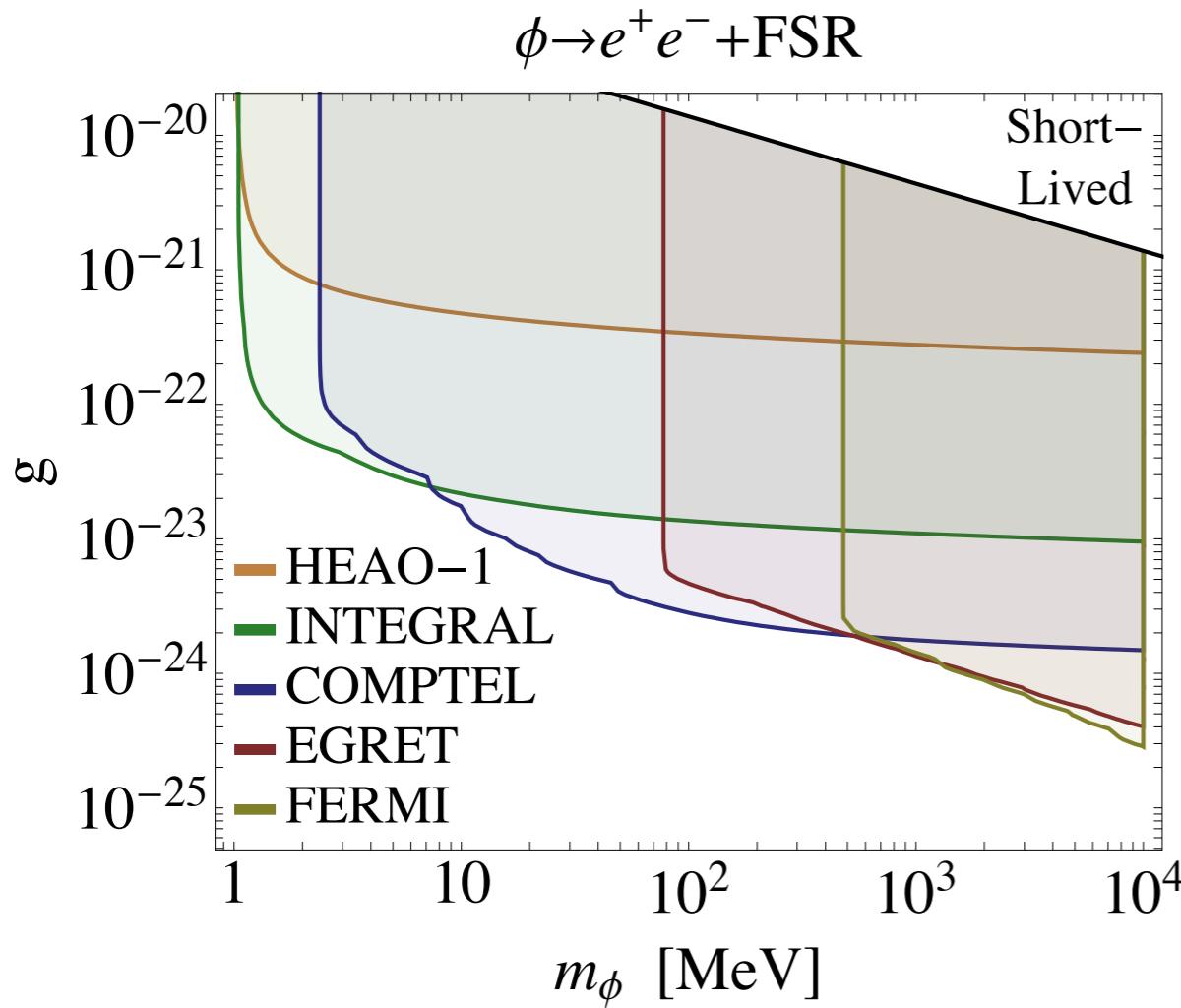
dimension 5 operator  
can be strongly constrained



# Dark (pseudo)scalars

$$\tau_{\phi \rightarrow e^+ e^-} \simeq 8.3 \times 10^{18} \text{ sec} \frac{10 \text{ MeV}}{m_\phi} \left( \frac{10^{-20}}{g_a} \right)^2$$

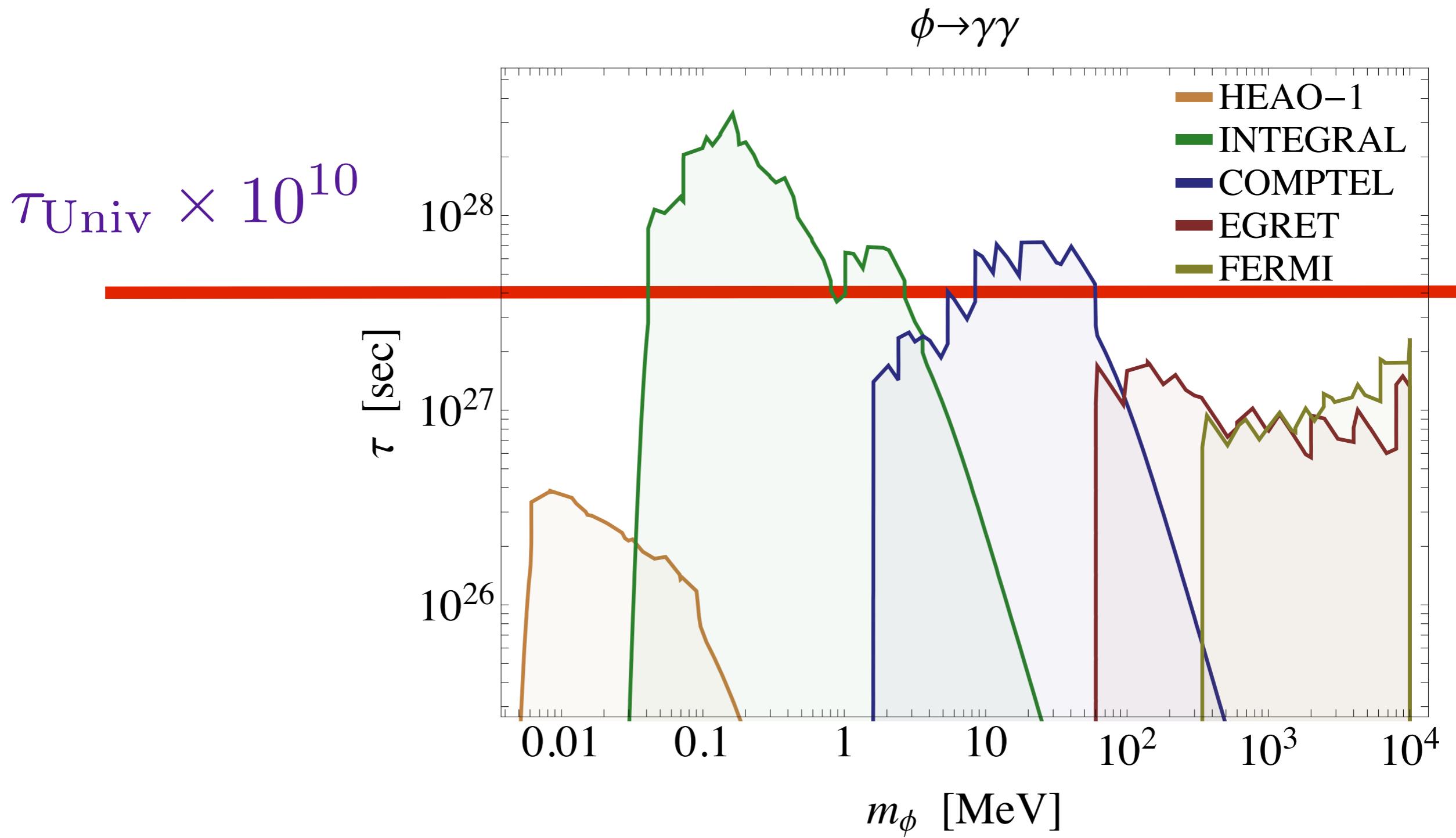
$$\tau_{\pi_d \rightarrow \gamma\gamma} \simeq 1.1 \times 10^{20} \text{ sec} \left( \frac{10 \text{ MeV}}{m_{\pi_d}} \right)^3 \left( \frac{f_{\pi_d}}{10^{15} \text{ GeV}} \right)^2$$



# Part I Halftime

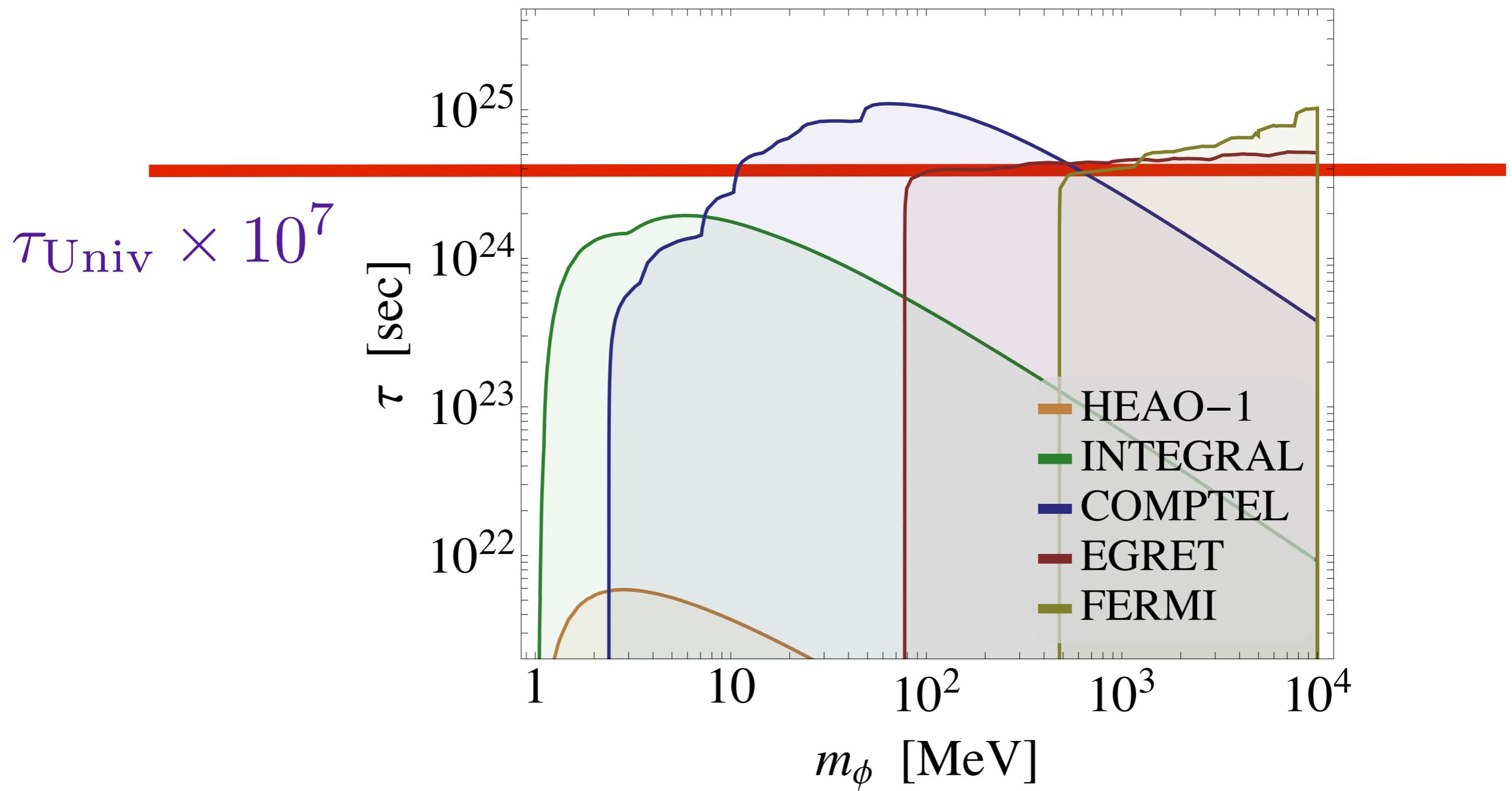
- Those were the model-dependent bounds
  - bounds on model-specific parameters (mixing angles, decay constants, etc.)
  - very strong for dimension $<6$ , non-Planck-suppressed operators
- About to show model-independent bounds
  - just the lifetime – mass plane from now on
  - lifetime bounds from 6 (FSR photons) to 10 (direct photons) orders of magnitude stronger than  $1/H_0$

# Photon line



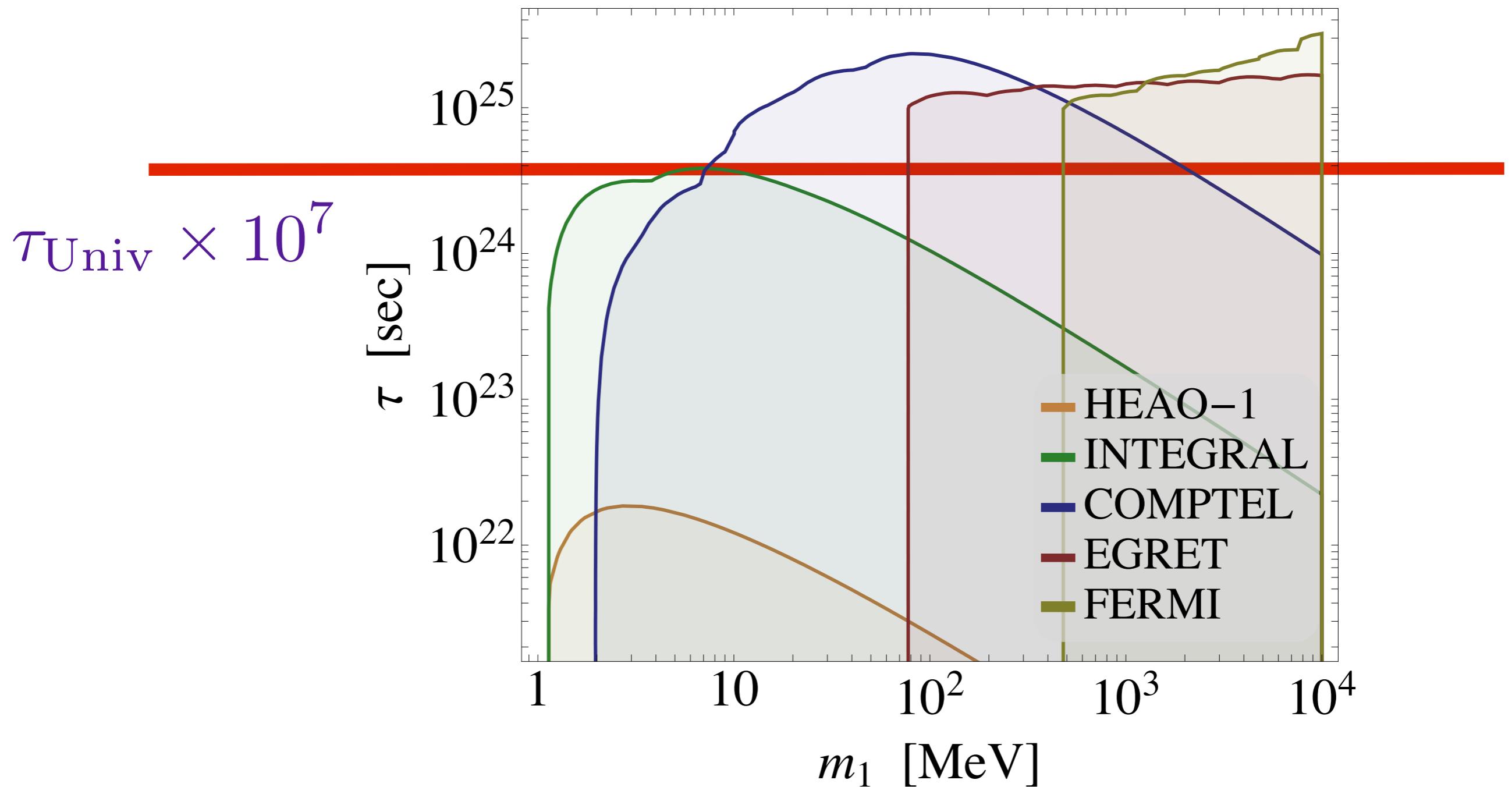
# e+ e- (FSR)

$\phi \rightarrow e^+ e^- + \text{FSR}$

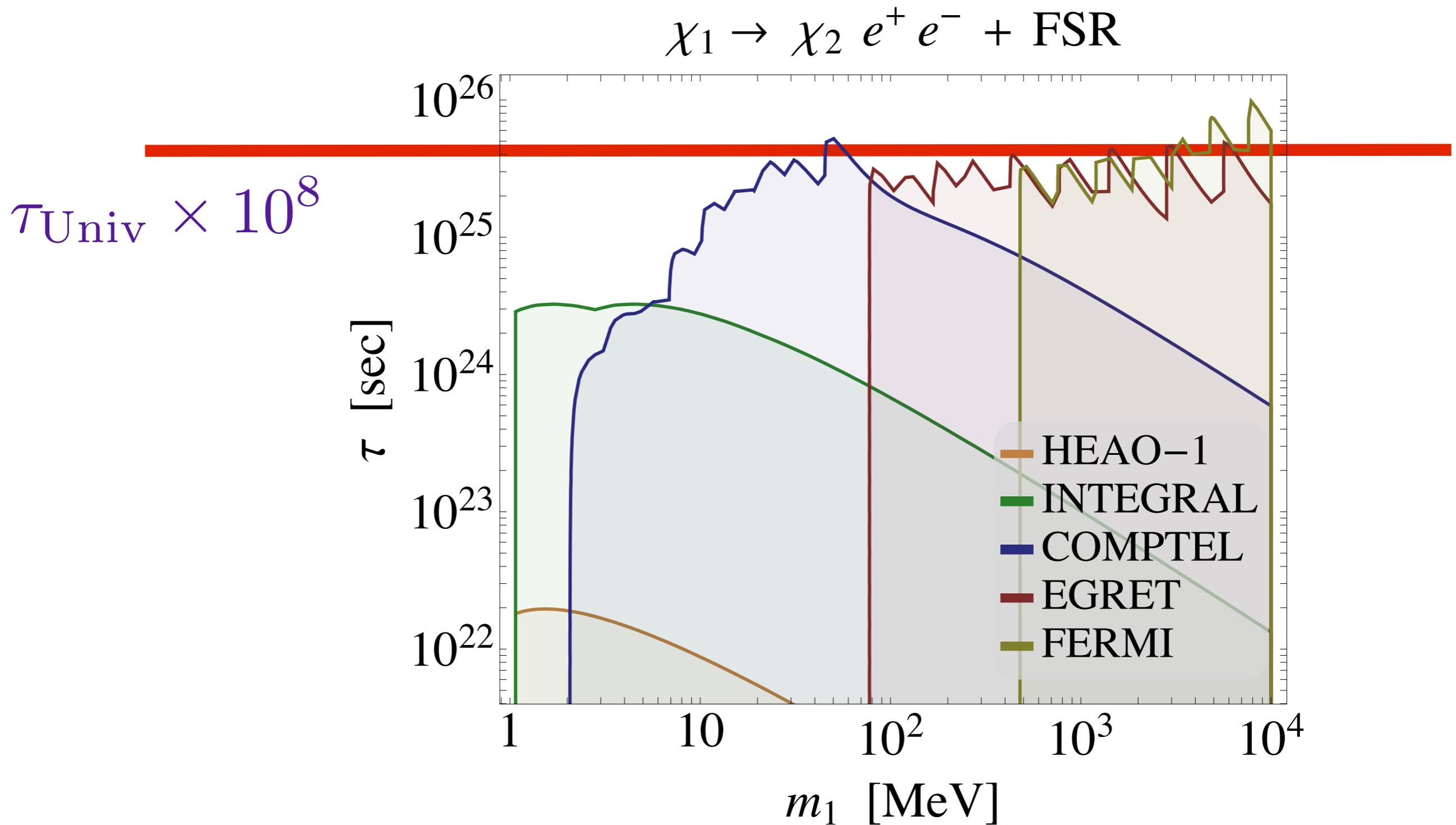


# e+ e- (FSR), boosted

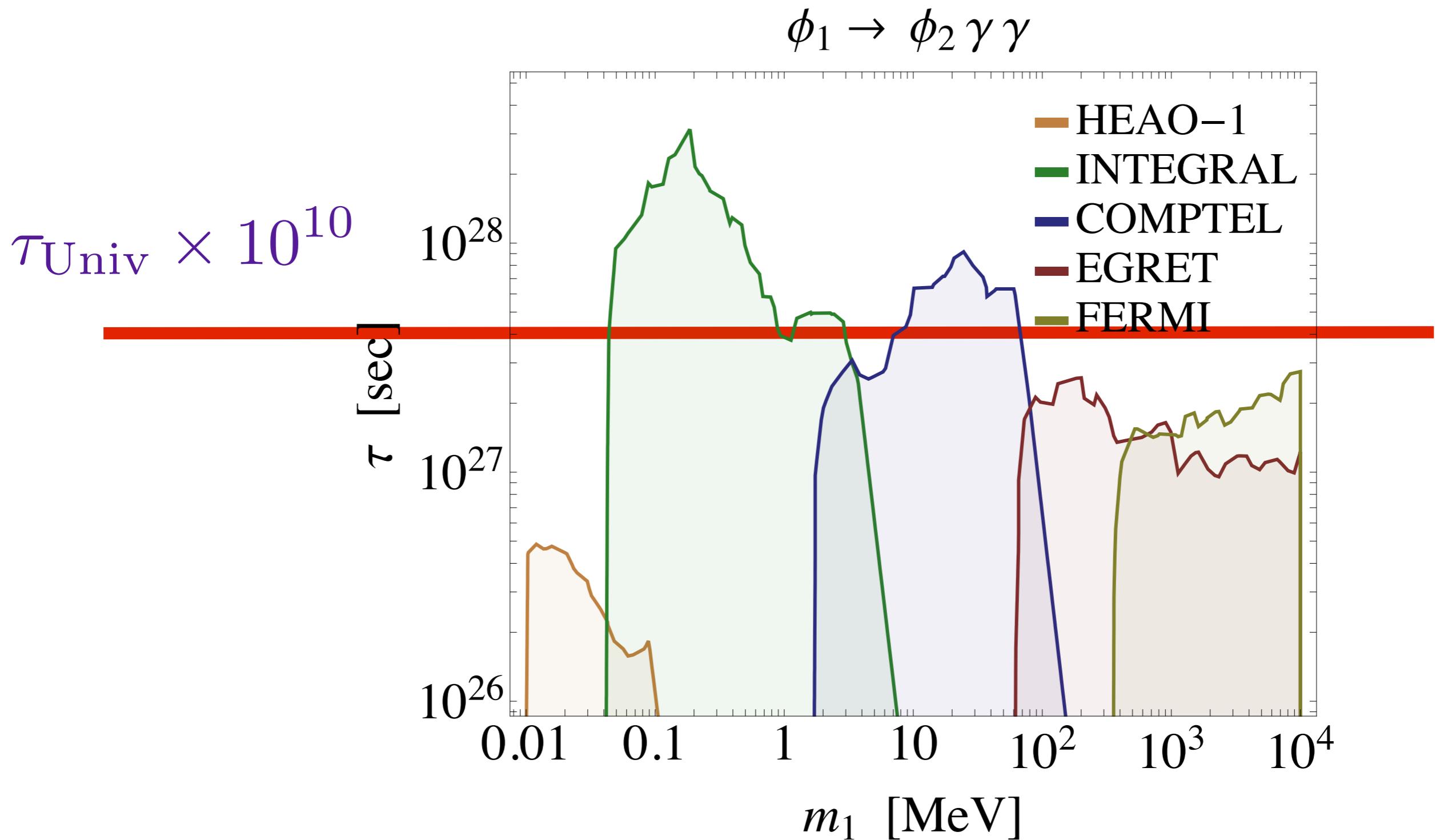
$\phi_1 \rightarrow \phi_2 \phi_3 \rightarrow \phi_2 e^+ e^-$



# e+ e- (FSR), three-body



# three-body, with photons



# Part I, Recap

- The decaying DM bounds:
  - model-specific parameters
  - model-independent parameters (lifetime, mass)
  - (very robust)

# Part II:

## “Light” (GeV-ish)

# Annihilating DM

# Part II, Outline

## Annihilating DM:

- effect of structure on DM annihilations
- relative contributions
- astrophysical backgrounds
- results

# Where does DM annihilate?

smooth galactic:  $\rho_s^2 r_s \simeq \mathcal{O}(10^{-47} \text{ GeV}^7)$

vs.

smooth extragalactic:  $\rho_{\text{DM}}^2 / H_0 \simeq \mathcal{O}(10^{-52} \text{ GeV}^7)$

First guess: mostly in the MW halo

# Structure

Structures (overdensities) dramatically change this picture because the rate of annihilations scales quadratically with density (not linearly, as for decays)

Where the DM is  
matters for annihilations!

# Structure

Critical density with which  
a sphere collapses instead  
of growing indefinitely:  $\rho_{\text{crit}} \simeq 5.5 \times \rho_{\text{av}}$

Final density after virializing:  $\rho_f = 32 \times \rho_i$

$$\delta \equiv \rho_{\text{structure}} / \rho_{\text{av}} \simeq 32 \cdot 5.5 \simeq 178$$

Improved estimate:

$$\rho_{\text{eg}}^2 \simeq \delta^2 \rho_{\text{DM}}^2 / H_0 \simeq \mathcal{O}(\text{few} \times 10^{-48} \text{ GeV}^7) \simeq \rho_s^2 r_s$$

# Different contributions

- The smooth galactic and clumpy extragalactic intensities are roughly similar in magnitude
  - this makes sense: the MW is (mostly) a DM halo!
- But there is additional structure:
  - clusters at all redshifts have subclusters (analogously: the MW has subhalos)
  - Add it all up!

# Milky Way

(from smooth distribution)  $\frac{dI_{\text{sm}}}{dE} = \frac{\langle \sigma v \rangle}{2m_\chi^2} \frac{dN_\gamma}{dE} \int_{V_{\text{obs}}} dV_{\text{MW}} \frac{\rho^2(s, b, \ell)}{4\pi s^2}$

(from satellites)

$$\frac{dI_{\text{sat}}}{dE} = \frac{\langle \sigma v \rangle}{2m_\chi^2} \frac{dN_\gamma}{dE} \int dV_{\text{MW}} \frac{1}{4\pi s^2} \times$$

satellite mass function (uncertain)

$$\times \frac{dn_{\text{sat}}(s, b, \ell, M)}{dM} \int dV_{\text{sat}} \rho_{\text{sat}}^2(M)$$

# Extragalactic

$$\frac{dI_{\text{eg}}}{dE} = \frac{\langle \sigma v \rangle}{2m_\chi^2} \int dz \frac{dN_\gamma[E(1+z)]}{dE} \frac{\bar{\rho}^2(z)}{4\pi} \times \times \frac{(1+z)^3}{H(z)} e^{-\tau[E(1+z), z]}$$
$$\bar{\rho}^2(z) = \int dM \frac{dn(M, z)}{dM} [1 + b_{\text{sh}}(M)] \times$$

boost from  
substructure  
(uncertain)

halo mass  
function  
(uncertain)

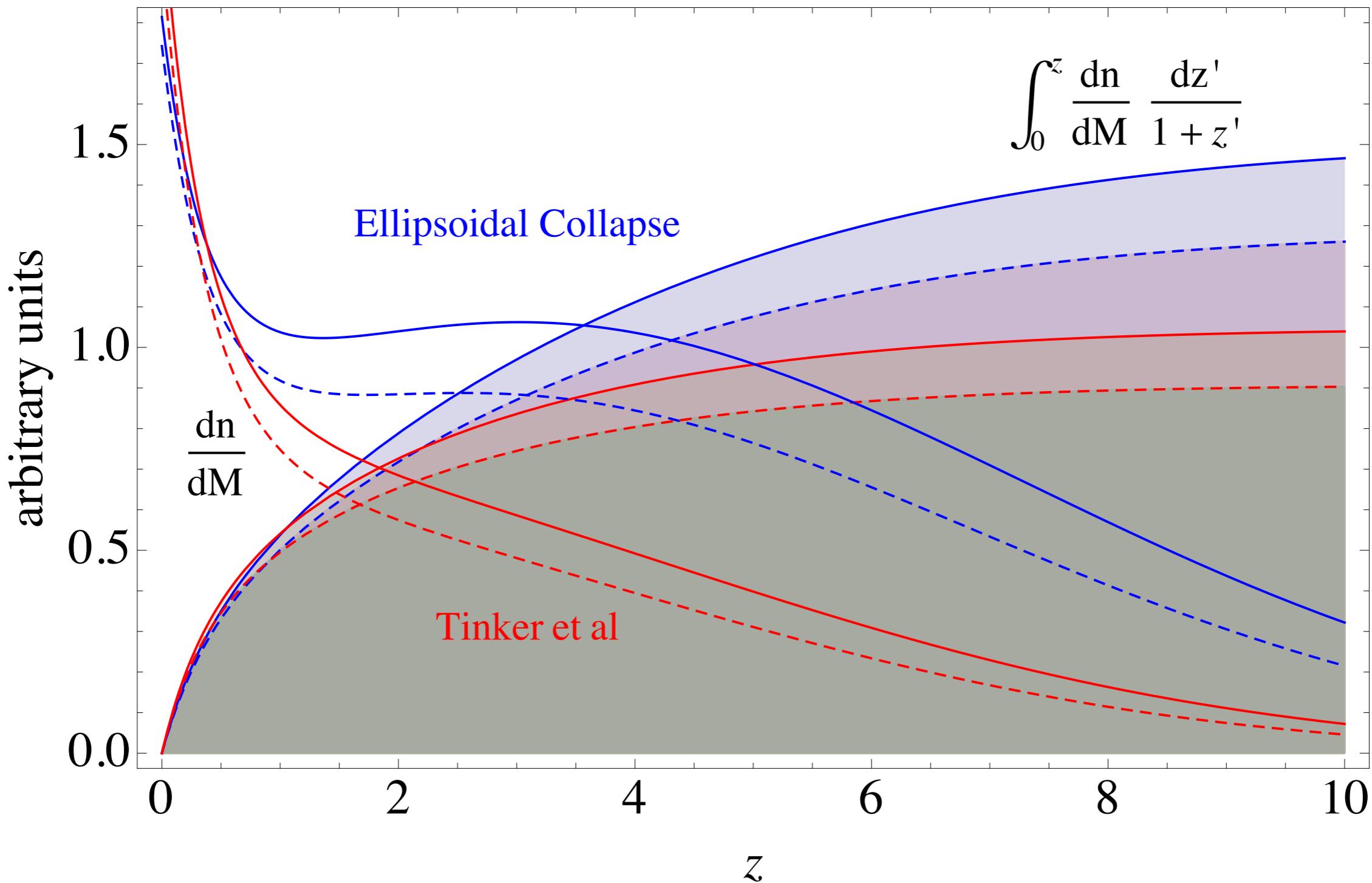
$$\times \int dV \rho_{\text{host}}^2(r, M)$$

# Uncertainties

- $dN/dE$  from PPPC DM ID (Pythia+EW corrections)
- optical depth from semi-analytic modeling  
(Gilmore, Primack, et al)
- halo mass function and subhalo boost factor
  - semi-analytic fits (Phoenix simulation)
  - Anderhalden and Diemand

# for example...

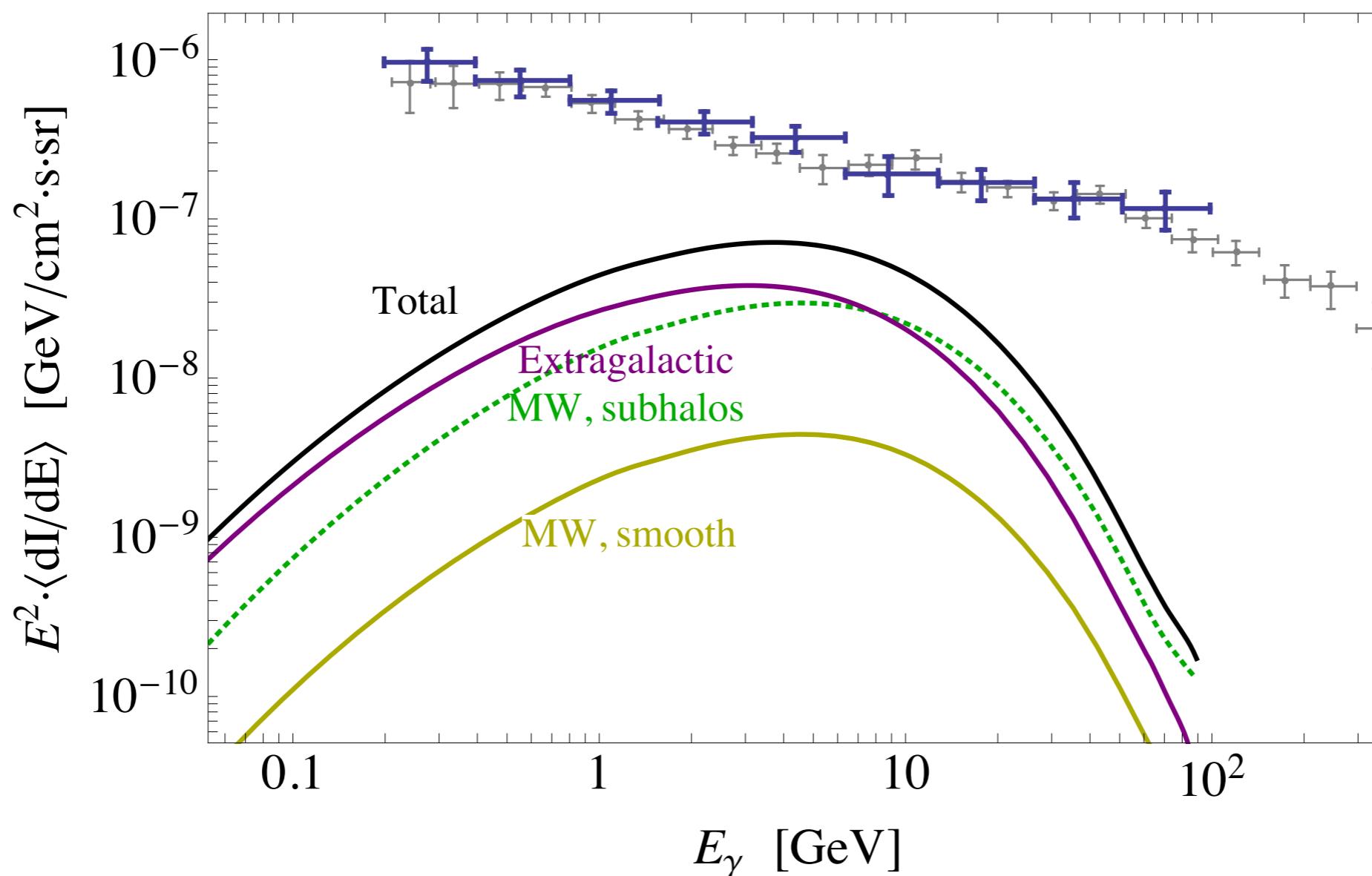
Planck+WMAP (Solid) vs. WMAP (Dashed)



# Relative intensities

$m_{\text{DM}}=100 \text{ GeV}$ ,  $\langle \sigma v \rangle = 3 \times 10^{-26} \text{ cm}^3/\text{s}$

channel: DM DM  $\rightarrow b \bar{b}$

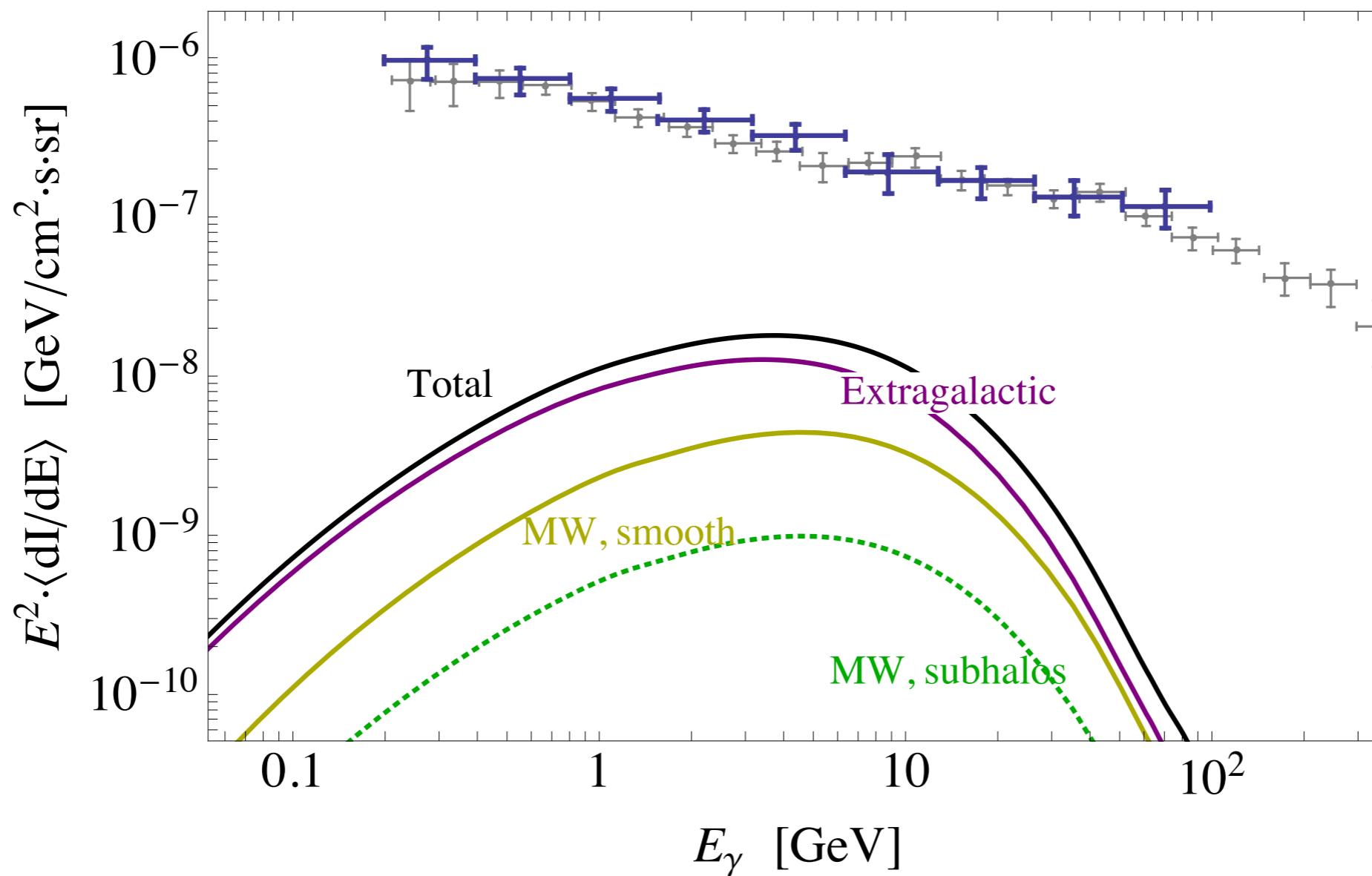


default substructure calculation

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channel: DM DM  $\rightarrow b \bar{b}$

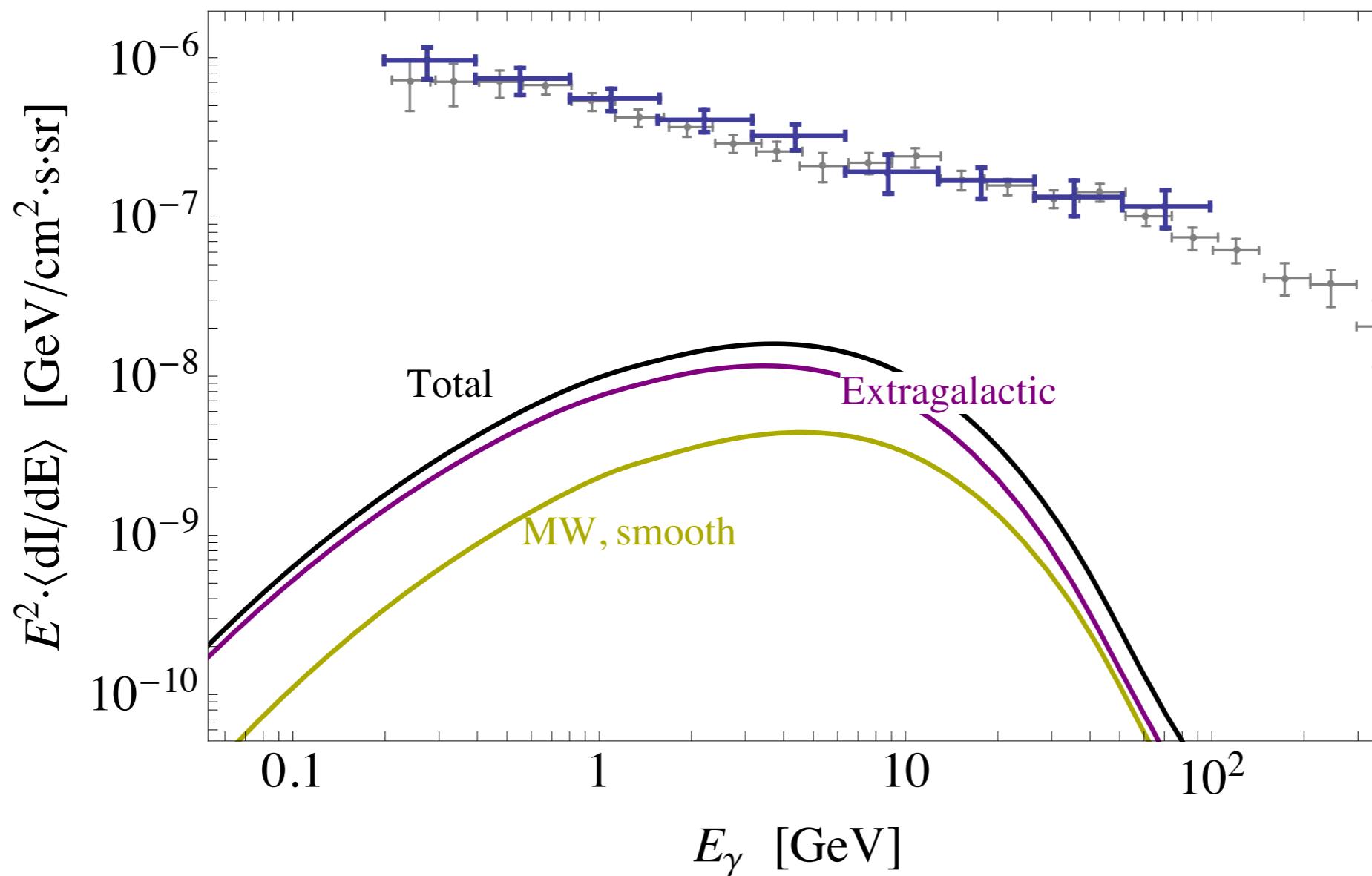


conservative substructure calculation

# Relative intensities

$m_{\text{DM}}=100 \text{ GeV}$ ,  $\langle \sigma v \rangle = 3 \times 10^{-26} \text{ cm}^3/\text{s}$

channel: DM DM  $\rightarrow b \bar{b}$

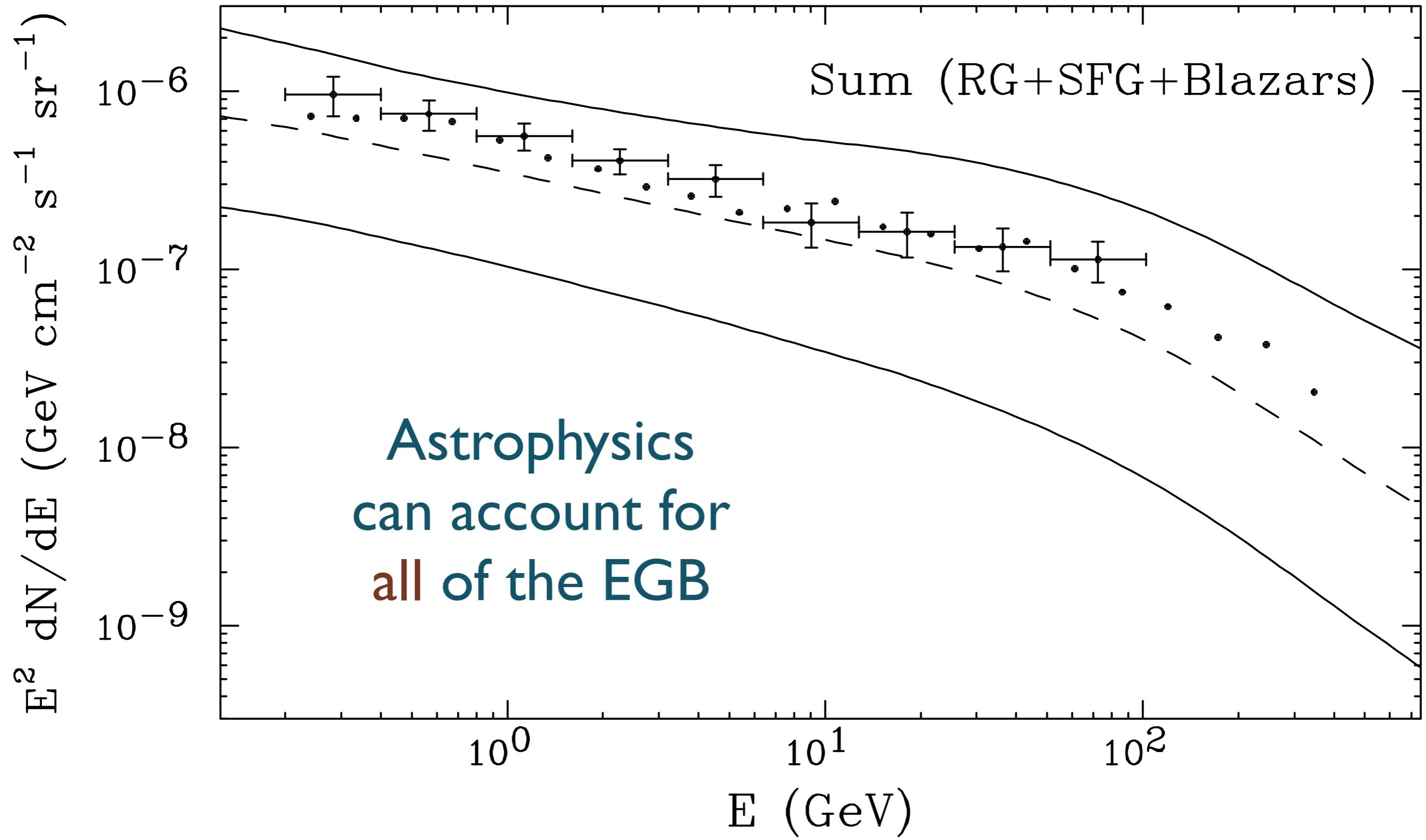


NO substructure

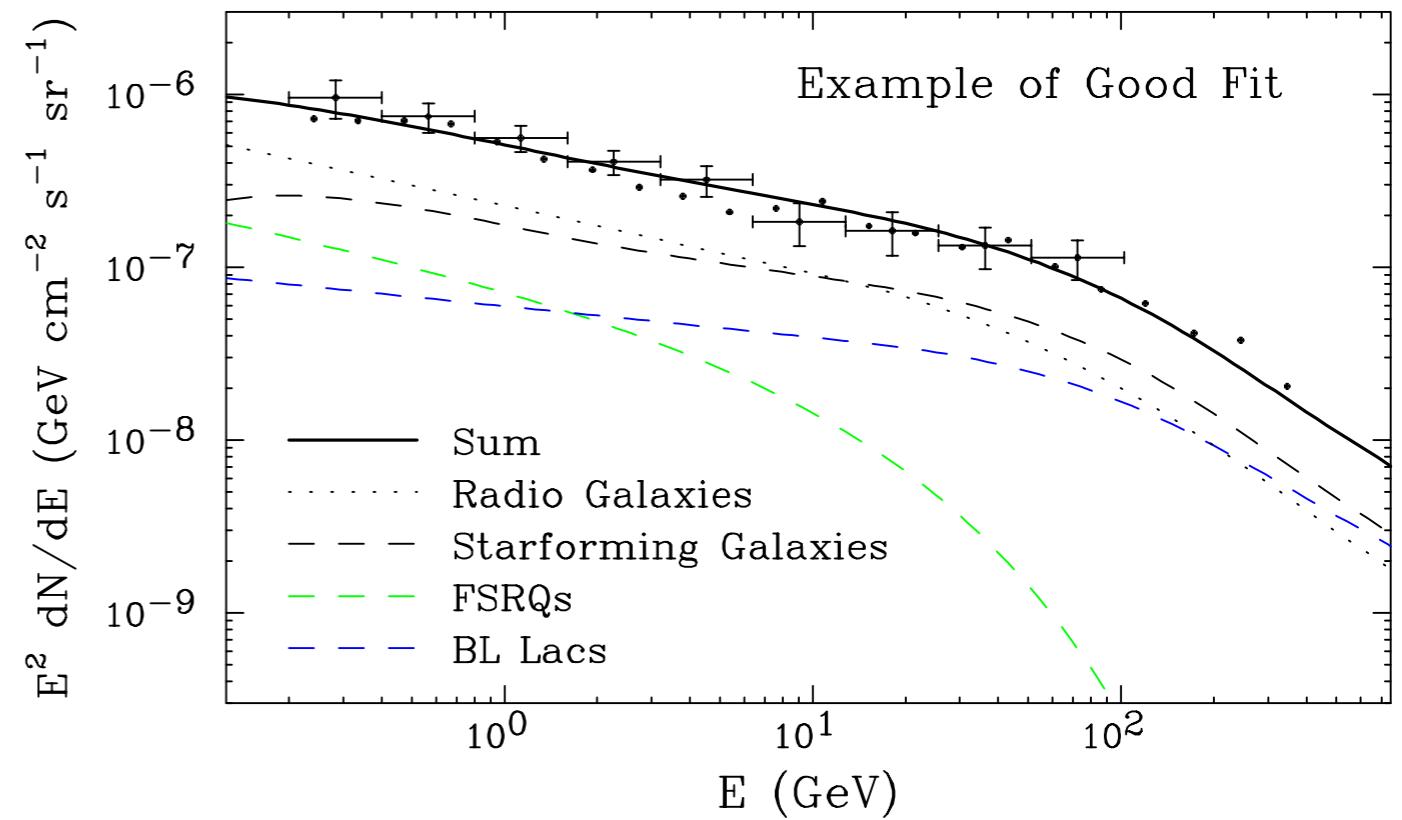
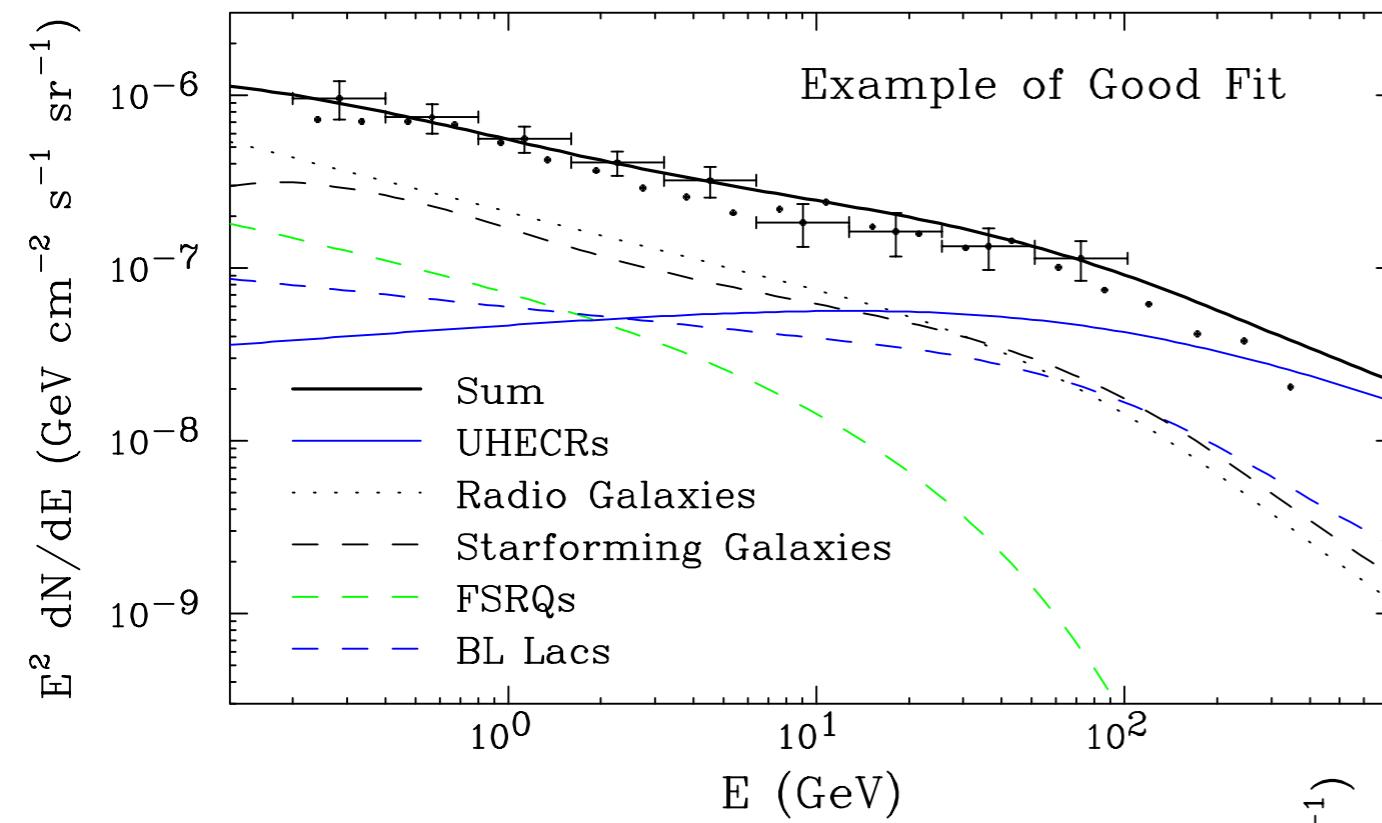
# Many backgrounds

- star forming galaxies
- blazars (resolved and unresolved)
  - radio galaxies (BL Lactaea objects, FSRQs, etc.)
- ultra-high energy cosmic rays
- millisecond pulsars (...)

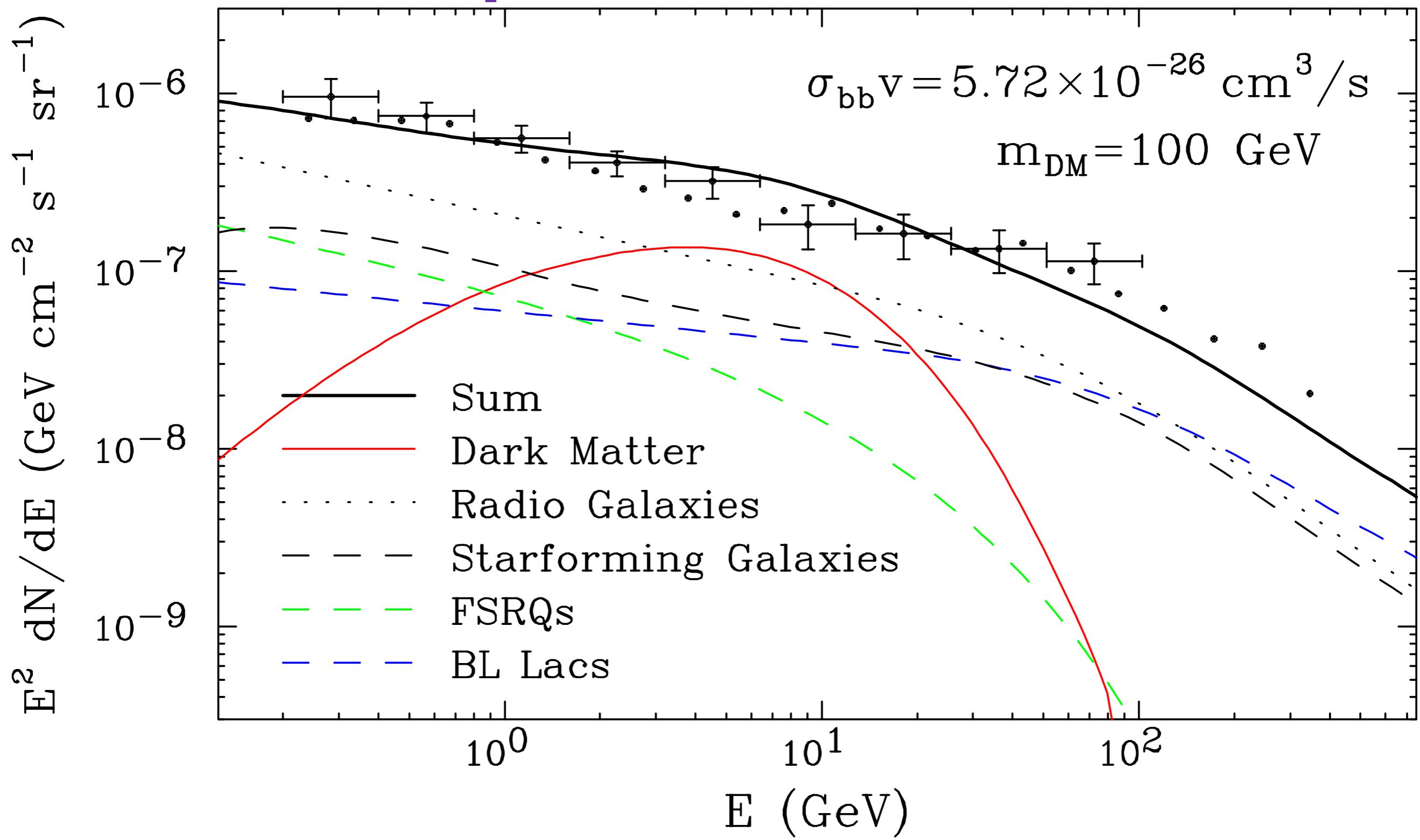
# Backgrounds



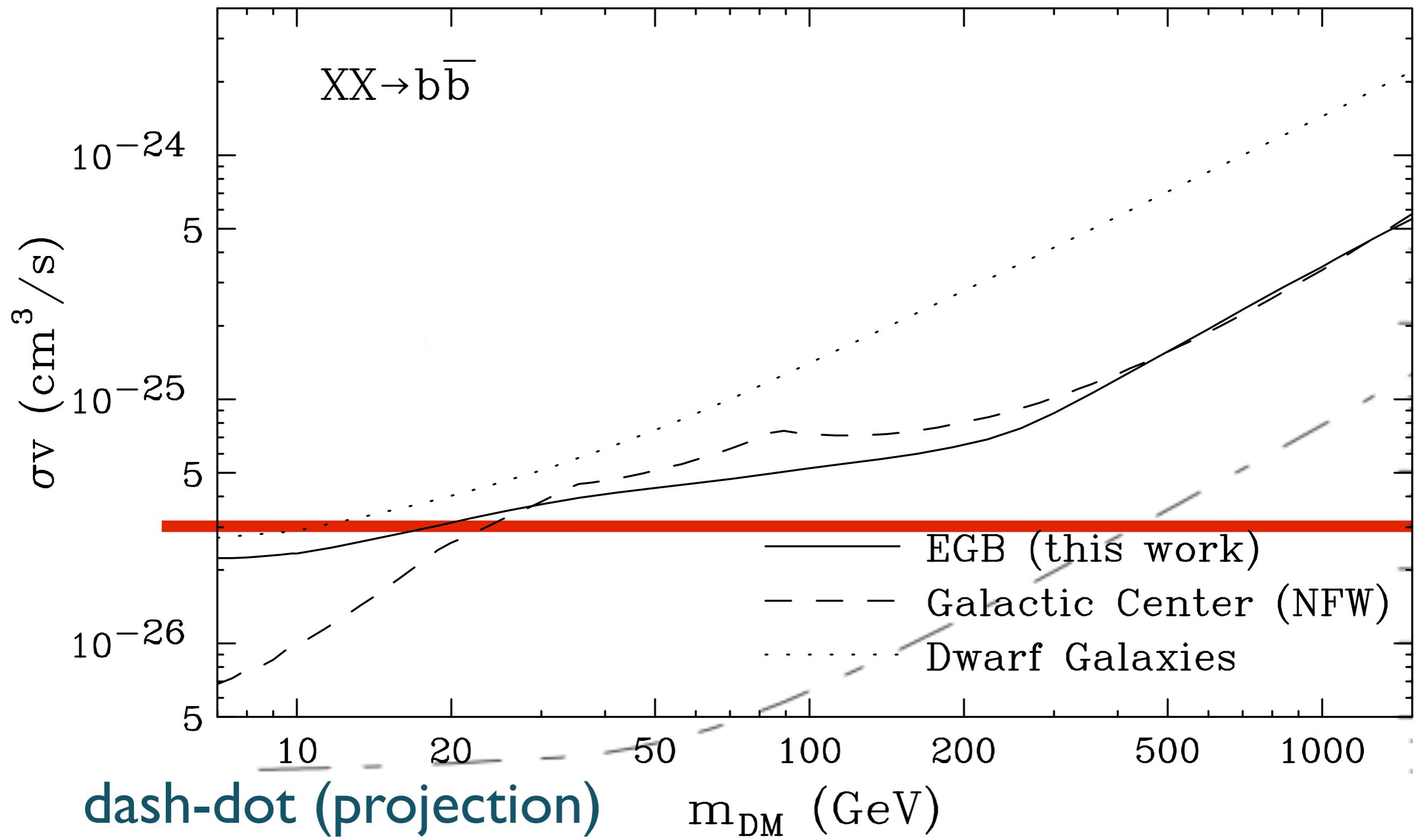
# Fits without DM



# Example fit with DM



# Cross-section bounds



# Important probe

- Complementary to observations of stacked dwarf spheroidals, galactic center, ... (similarly strong or stronger, with very different systematics)
- Very strong bounds!
  - currently bounds thermal DM up to  $\sim 20$  GeV
  - Projected to be the strongest bound by end of the Fermi mission
  - can bound thermal DM up to  $\sim 400$  GeV

# Conclusions

- Bounds on light decaying DM from the galactic diffuse background are strong (and robust!) even though observations are not DM-centric
- Looking outside the galaxy is a promising method for putting strong constraints on more massive annihilating DM

# Thank you!

Other in-progress projects I'm excited to talk about:

- thoughts on relic neutrino detection
- new work related to DM direct detection
- extensions of work on DM in neutron stars
- collider phenomenology of a light Higgs partner