

Toward a self-consistent and unitary nuclear reaction network

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Outline

- Overview: LANL light nuclear reaction program
- **Unitarity**: SBBN & beyond, reaction networks, R-matrix, ^{17}O , ^9B examples
- Recent related development: EFT \longleftrightarrow R-matrix ($a_c \rightarrow 0 \quad \forall c$)
- Future work & conclusion

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Light nuclear reaction program @ LANL

■ Motivation

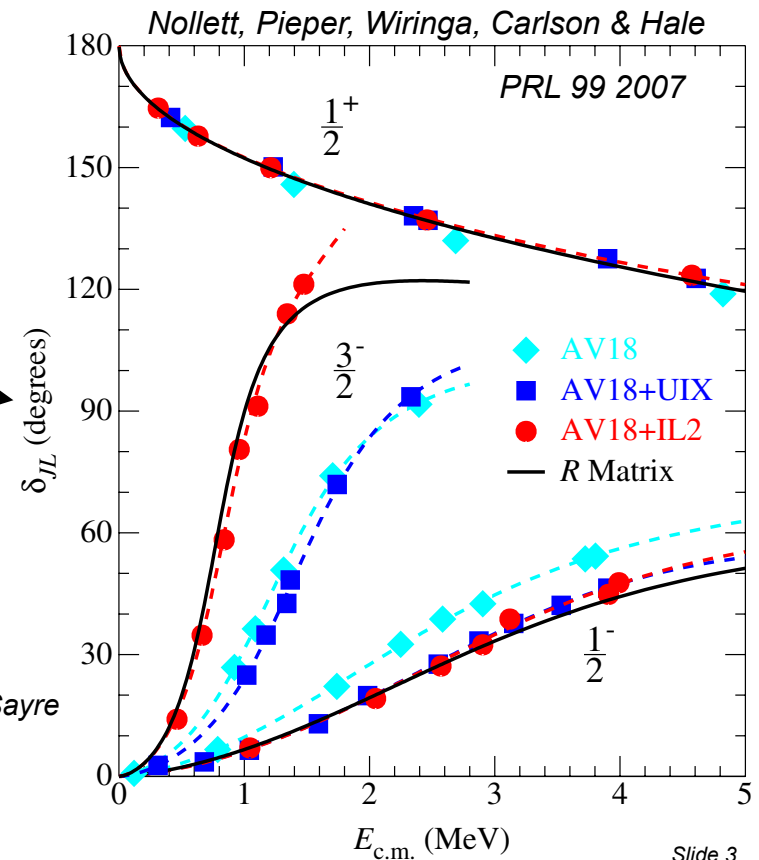
- Data sets: σ , $\sigma(\theta)$, $A_i(\theta)$, $C_{i,j}$, K_1^J , $\Sigma(\gamma)$, ... → T matrix → resonance spectrum
- **Unitary** parametrization of compound nuclear system
- Applications: **astrophysical**, nuclear security, inertial confinement fusion, **criticality safety**, charge-particle transport, nuclear data (ENDF, ENSDF)

■ Ab initio

- Variational MC; Green's function MC
- GFMC [PRL **99**, 022502 (2007)]
 - n-⁴He phase shifts
 - comparison GFMC/R-matrix
- challenge: multichannel
 - eg. $n\alpha \rightarrow n\alpha$, $n\alpha \rightarrow dt$ & $dt \rightarrow dt$

■ Phenomenology

- R matrix (2→2 body scatt/reacs)
- 3-body: isobaric models, sequential decay
 - R-Matrix description of particle energy spectra produced by low-energy T + T reactions; w/G. Hale, C. Brune (OU), D. Sayre & J. Caggiano (LLNL)



EDA Analyses of Light Systems

A	System	Channels	Energy Range (MeV)
2	N-N	p+p; n+p, γ +d	0-30 0-40
3	N-d	p+d; n+d	0-4
4	^4H ^4Li	n+t p+ ^3He	0-20
	^4He	p+t n+ ^3He d+d	0-11 0-10 0-10
5	^5He	n+ α d+t $^5\text{He}+\gamma$	0-28 0-10
	^5Li	p+ α d+ ^3He	0-24 0-1.4

Analyses of Light Systems, Cont.

A	System (Channels)
6	${}^6\text{He}$ (${}^5\text{He}+n$, $t+t$); ${}^6\text{Li}$ ($d+{}^4\text{He}$, $t+{}^3\text{He}$); ${}^6\text{Be}$ (${}^5\text{Li}+p$, ${}^3\text{He}+{}^3\text{He}$)
7	${}^7\text{Li}$ ($t+{}^4\text{He}$, $n+{}^6\text{Li}$); ${}^7\text{Be}$ ($\gamma+{}^7\text{Be}$, ${}^3\text{He}+{}^4\text{He}$, $p+{}^6\text{Li}$)
8	${}^8\text{Be}$ (${}^4\text{He}+{}^4\text{He}$, $p+{}^7\text{Li}$, $n+{}^7\text{Be}$, $p+{}^7\text{Li}^*$, $n+{}^7\text{Be}^*$, $d+{}^6\text{Li}$)
9	${}^9\text{Be}$ (${}^8\text{Be}+n$, $d+{}^7\text{Li}$, $t+{}^6\text{Li}$); ${}^9\text{B}$ ($\gamma+{}^9\text{B}$, ${}^8\text{Be}+p$, $d+{}^7\text{Be}$, ${}^3\text{He}+{}^6\text{Li}$)
10	${}^{10}\text{Be}$ ($n+{}^9\text{Be}$, ${}^6\text{He}+\alpha$, ${}^8\text{Be}+nn$, $t+{}^7\text{Li}$); ${}^{10}\text{B}$ ($\alpha+{}^6\text{Li}$, $p+{}^9\text{Be}$, ${}^3\text{He}+{}^7\text{Li}$)
11	${}^{11}\text{B}$ ($\alpha+{}^7\text{Li}$, $\alpha+{}^7\text{Li}^*$, ${}^8\text{Be}+t$, $n+{}^{10}\text{B}$); ${}^{11}\text{C}$ ($\alpha+{}^7\text{Be}$, $p+{}^{10}\text{B}$)
12	${}^{12}\text{C}$ (${}^8\text{Be}+\alpha$, $p+{}^{11}\text{B}$)
13	${}^{13}\text{C}$ ($n+{}^{12}\text{C}$, $n+{}^{12}\text{C}^*$)
14	${}^{14}\text{C}$ ($n+{}^{13}\text{C}$)
15	${}^{15}\text{N}$ ($p+{}^{14}\text{C}$, $n+{}^{14}\text{N}$, $\alpha+{}^{11}\text{B}$)
16	${}^{16}\text{O}$ ($\gamma+{}^{16}\text{O}$, $\alpha+{}^{12}\text{C}$)
17	${}^{17}\text{O}$ ($n+{}^{16}\text{O}$, $\alpha+{}^{13}\text{C}$)
18	${}^{18}\text{Ne}$ ($p+{}^{17}\text{F}$, $p+{}^{17}\text{F}^*$, $\alpha+{}^{14}\text{O}$)

LANL R-matrix light element IAEA standards: status

n+p: no new work since 2008; cross sections in pretty good shape below 30 MeV; main need is extension to higher energies (150-200 MeV), with associated covariances.

n+³He: Some new work, especially for n+³He capture; ³He(n,n)³He scattering data re-worked by Drogg and Lisowski – could be used in a new analysis of the ⁴He system.

n+⁶Li: Some new work on ⁷Li system around 2008 included new LANSCE measurements of ⁶Li(n,t)⁴He differential cross section – was included in ENDF/B-VII.1 above 1 MeV. Cross sections should be re-visited below 1 MeV, although there may not be any new data.

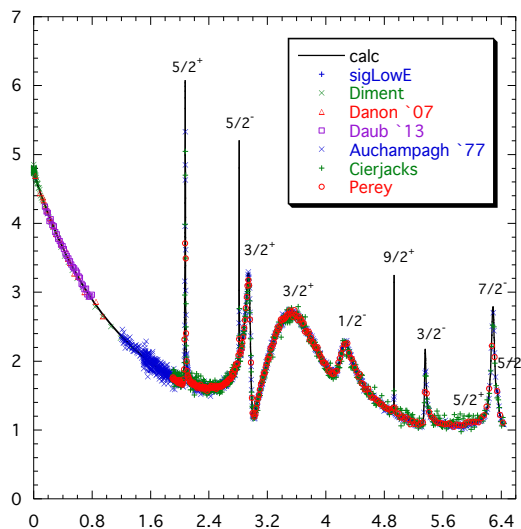
n+¹⁰B: No new work since last standards evaluation. New data from Geel, Ohio U. [including first measurement of the (n,p) cross section]. **R**-matrix analysis for ¹¹B system should be extended above 1 MeV.

n+^{12,13}C: Considerable new work in the last couple of years. New data for ¹²C(n,n' γ) from Geel, Los Alamos changed the (n,n') cross sections. Isotopic evaluations combined to make more accurate standard evaluation for C-0.

Notation: X(a,b)Y
X-target Y-recoil
a-projectile b-detected

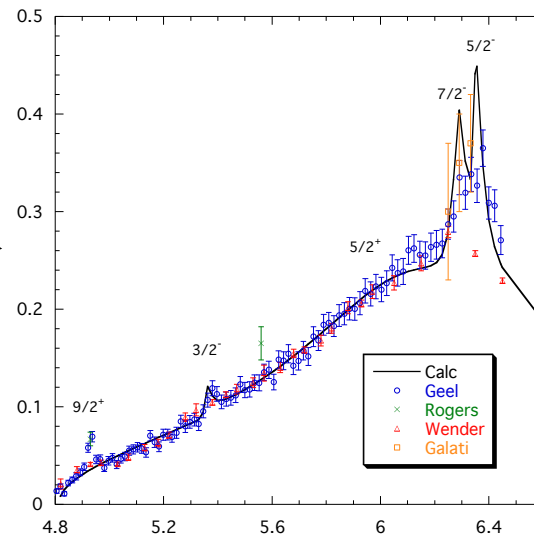
$^{13,14}\text{C}$ system analyses: σ_T (b) vs. E_n (MeV)

$n+^{12}\text{C}$ Total Cross Section

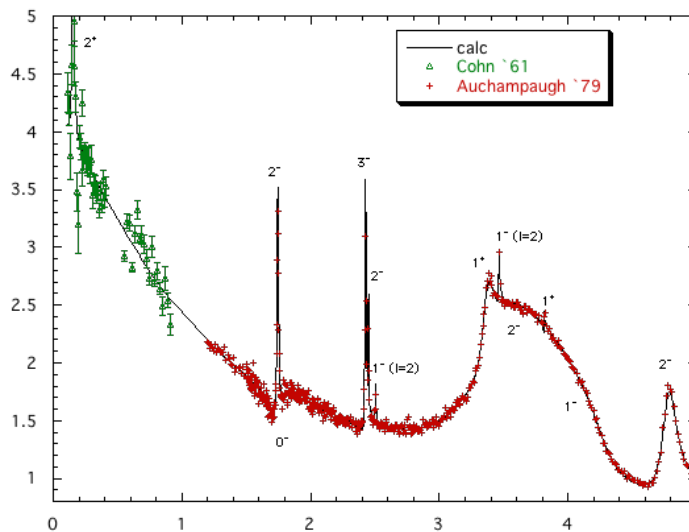


Two-channel analysis

$^{12}\text{C}(n,n')$ Cross Section



$n+^{13}\text{C}$ Total Cross Section



Single-channel analysis

Analyses by GMH/MWP

Unitary, self-consistent primordial nucleosynthesis

■ State of standard big-bang nucleosynthesis (BBN)

- d & ^4He abundances: signature success cosmology+nuc astro+astroparticle
 - but there's at least one **Lithium (^7Li) Problem** [^6Li too? See: [Lind et.al. 2013](#)]
- coming *precision* observations of d, ^4He , η , N_{eff} demand new BBN capabilities
- resolution of ^7Li problem:
 - observational/stellar astrophysics?
 - ^7Li controversial anomaly: nuclear physics solution?
 - new physics?

■ Advance BBN as a tool for precision cosmology

- incorporate **unitarity** into strong & electroweak interactions (**next slide**)
- couple **unitary reaction network (URN)** to full Boltzmann transport code
 - neutrino energy distribution function evolution/transport code
 - fully coupled to nuclear reaction network
 - calculate light primordial element abundance for non-standard BBN
 - active-sterile ν mixing
 - massive particle out-of-equilibrium decays→energetic active SM particles
- Produce tools/codes for nuc-astro-particle community: test new physics w/BBN
 - existing codes are based on Wagoner's (1969) code

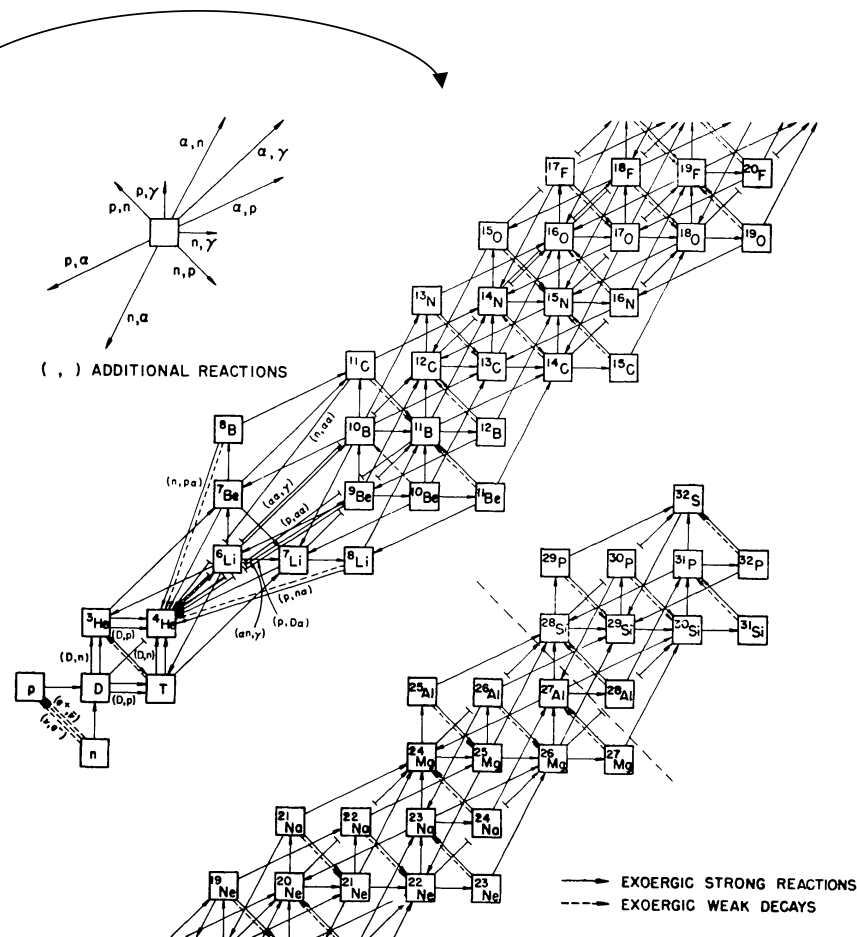
Nuclear reaction network

- **Single-process (non-unitary) analysis**

- $\sigma_{\alpha\beta}(E) \pm \delta\sigma_{\alpha\beta}(E)$ from expt
- fit form (non-res+narrow res) to $\sigma_{\alpha\beta}(E)$
- compute $\langle\sigma v\rangle(T)$ → reactivity → network
- **NB: norm. systematics can be large**
 - ^{17}O case in 2 slides

■ Multi-channel (unitary) analysis

- Construct unitary parametrization
 - R-matrix (Wigner-Eisenbud '47)
- simultaneous fit of unpolarized/pol'd scatt/reac data → determine T (or S) matrix
- determines a unitary reaction network (URN) for analyzed compound systems



Wagoner ApJSuppl '69

Formal unitarity: consequences

$$\left. \begin{aligned} \delta_{fi} &= \sum_n S_{fn}^\dagger S_{ni} \\ S_{fi} &= \delta_{fi} + 2i\rho_f T_{fi} \\ \rho_n &= \delta(H_0 - E_n) \end{aligned} \right\} \quad T_{fi} - T_{fi}^\dagger = 2i \sum_n T_{fn}^\dagger \rho_n T_{ni}$$

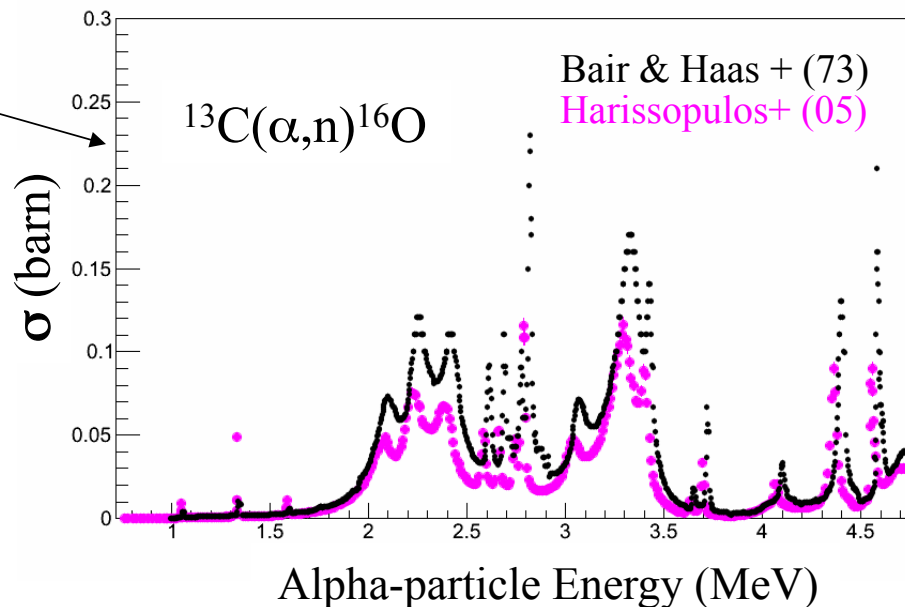
NB: **unitarity** implies optical theorem $\sigma_{\text{tot}} = \frac{4\pi}{k} \text{Im } f(0)$; but *not just* the O.T.

■ Implications of **unitarity** constraint on transition matrix

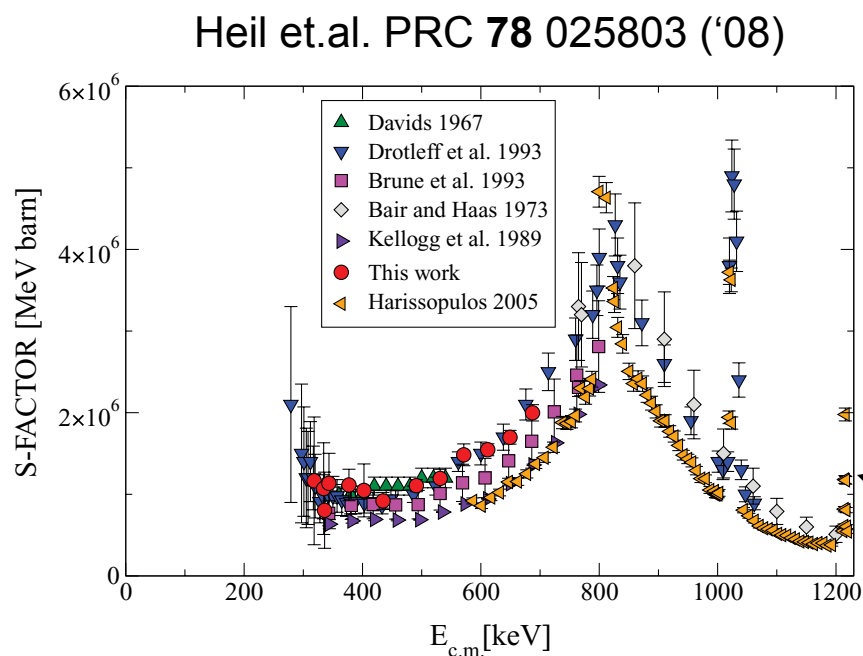
- Doesn't uniquely determine T_{ij} ; highly restrictive, however
Elastic: $\text{Im } T_{11} = -\rho_1$, $E < E_2$ (assuming T & P invariance)
Multichannel: $\text{Im } \mathbf{T} = -\rho$
- Unitarity violating transformations
 - cannot scale **any** set: $T_{ij} \rightarrow \alpha_{ij} T_{ij}$ $\alpha_{ij} \in \mathbb{R}$
 - cannot rotate **any** set: $T_{ij} \rightarrow e^{i\theta_{ij}} T_{ij}$ $\theta_{ij} \in \mathbb{R}$
 - ★ consequence of linear 'LHS' \propto quadratic 'RHS'
- Unitary parametrizations of data provide constraints that experiment may violate
 - ★ *normalization*, in particular \longrightarrow Observable $\propto \text{KF } |T_{fi}|^2$
 - ★ next slide: ^{17}O compound system

^{17}O compound system: experimental status

Recent Bochum (Harissopulos '05) measurement $^{13}\text{C}(\alpha, n)^{16}\text{O}$ vs. older ORNL (Bair & Haas '73)



Credit: S. Kunieda

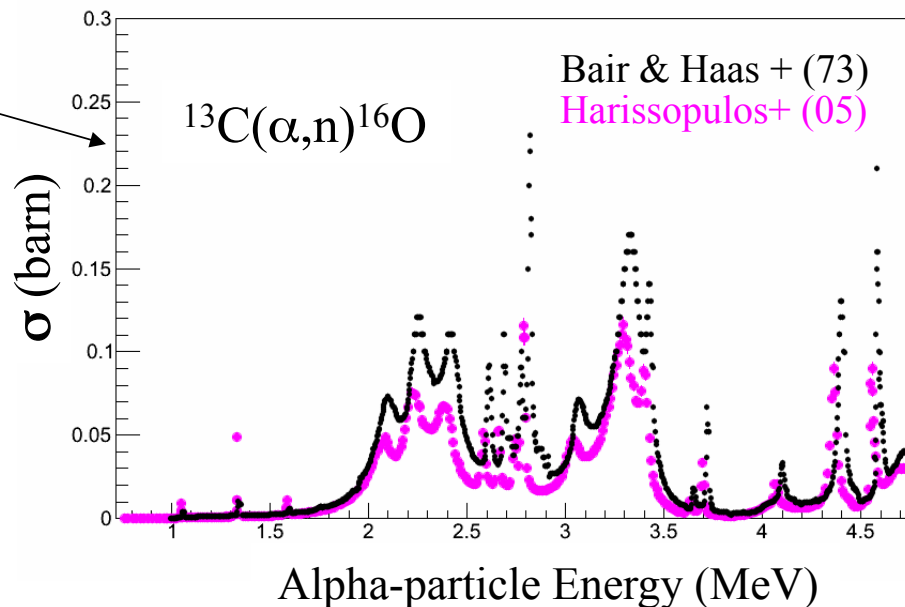


- Bochum data $2/3 \cdot \text{B\&H}(73)$
- Karlsruhe (Heil '08) consistent w/ORNL

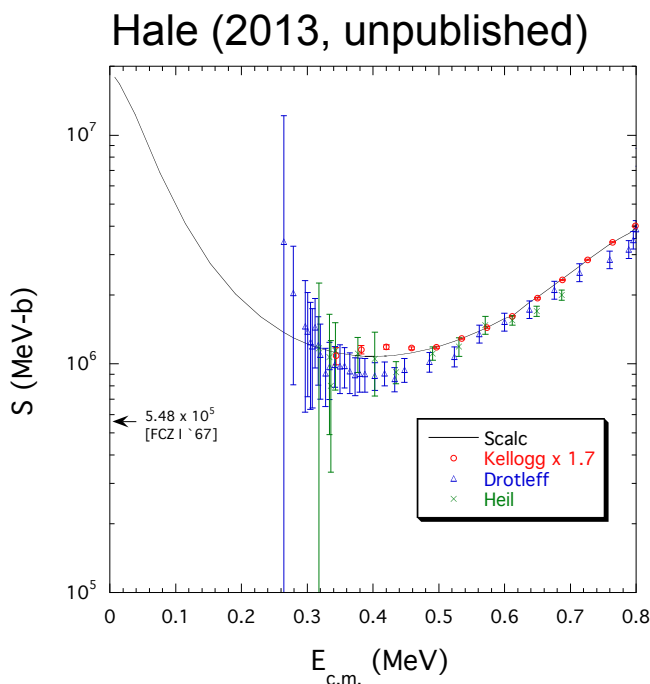
Tempting to conclude that B&H73 was right all along!

^{17}O compound system: experimental status

Recent Bochum (Harissopulos '05) measurement $^{13}\text{C}(\alpha, n)^{16}\text{O}$ vs. older ORNL (Bair & Haas '73)



Credit: S. Kunieda



- Subthreshold $\frac{1}{2}^+$
 - deep min in σ_T
 - $S(0) \gg S_{\text{FCZ67}}(0)$

Tempting to conclude that B&H73 was right all along!

Basics of R-matrix

■ Assumptions (cf. Lane & Thomas RMP '58)

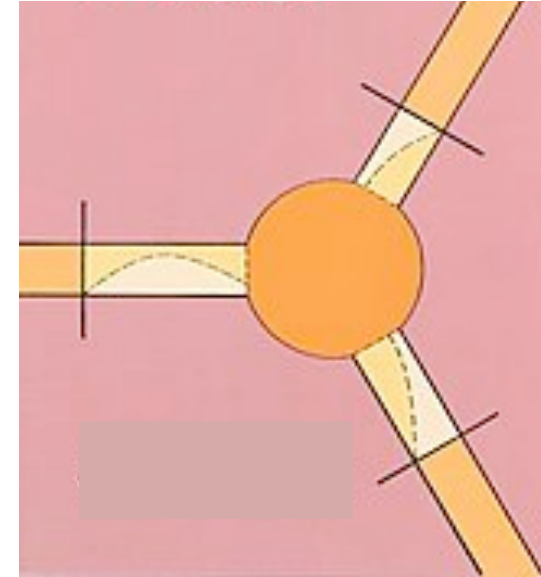
- a) Non-relativistic QM (L&T58); LANL-EDA uses rel. kin.
- b) Two-body channels only ('c'); aux. spectra code
- c) Conservation of N, Z
- d) Finite radius a_c beyond $V_{\text{pol}} \approx 0$; sharp boundaries

■ Separated pairs, "channels"

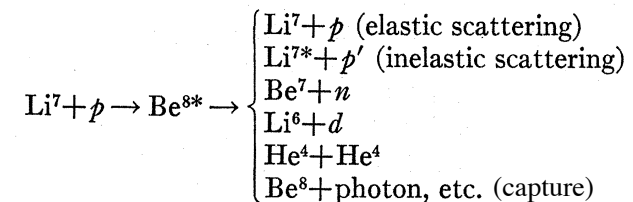
- A nucleons $\rightarrow (A_1, A_2)$
- $c = \{\alpha s_1 m_1 s_2 m_2\} \rightarrow \{\alpha(s_1 s_2) s m_s \ell m_\ell\} \rightarrow \{\alpha(s_1 s_2) s \ell, J M\}$
- Assume $a_c = a_\alpha \rightarrow$ many c have same channel in configuration space

■ Channel surface

- Consider configuration space of $3A$ dimensions
- Set of points: $\cup_c r_{\alpha(c)} = a_{\alpha(c)}$
- Surfaces coincide but assumed to have negl. prob.
- Channels are cylinders normal to channel surf.

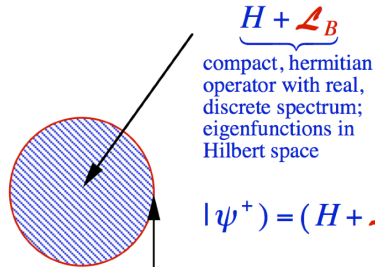


Example: ^8Be compound system



R-matrix formalism

INTERIOR (Many-Body) REGION
(Microscopic Calculations)



$$H + \mathcal{L}_B$$

$$|\psi^+\rangle = (H + \mathcal{L}_B - E)^{-1} \mathcal{L}_B |\psi^+\rangle$$

SURFACE

$$\mathcal{L}_B = \sum_c |c\rangle \left(\frac{\partial}{\partial r_c} r_c - B_c \right) \langle c|$$

$$\langle \mathbf{r}_c | c \rangle = \frac{\hbar}{\sqrt{2\mu_c a_c}} \frac{\delta(r_c - a_c)}{r_c} \left[(\phi_{s_1}^{\mu_1} \otimes \phi_{s_2}^{\mu_2})_s^\mu \otimes Y_l^m(\hat{\mathbf{r}}_c) \right]_J^M$$

$$R_{c'c} = \langle c' | (H + \mathcal{L}_B - E)^{-1} | c \rangle = \sum_\lambda \frac{\langle c' | \lambda \rangle \langle \lambda | c \rangle}{E_\lambda - E}$$

ASYMPTOTIC REGION
(S-matrix, phase shifts, etc.)

$$\langle r_{c'} | \psi^+ \rangle = -I_{c'}(r_{c'}) \delta_{c'c} + O_{c'}(r_{c'}) S_{c'c}$$

Measurements

$$\begin{aligned} \mathbf{T} &= \rho^{1/2} \mathbf{O}^{-1} \mathbf{R}_L \mathbf{O}^{-1} \rho^{1/2} - \mathbf{F} \mathbf{O}^{-1} \\ \mathbf{S} &= 1 + 2\pi i \mathbf{T} \\ \mathbf{R}_L &= [\mathbf{R}_B^{-1} - \mathbf{L} + \mathbf{B}]^{-1} \end{aligned}$$

Bloch operator $\mathcal{L}_B = \sum_c |c\rangle \left(\frac{\partial}{\partial r_c} r_c - B_c \right) \langle c|$ ensures
Hermiticity of Hamiltonian restricted to internal region

- R-matrix theory: **unitary**, multichannel parametrization of (not just resonance) data

- Interior/Exterior regions

- Interior: strong interactions
- Exterior: Coulomb/non-polarizing interactions
- Channel surface

$$\mathcal{S}_c : r_c = a_c \quad \mathcal{S} = \sum_c \mathcal{S}_c$$

- R-matrix elements

- Projections on channel surface functions $\langle \mathbf{r}_c | c \rangle$ of Green's function

$$G_B = [H + \mathcal{L}_B - E]^{-1}$$

- Boundary conditions

$$B_c = \frac{1}{u_c(a_c)} \frac{du_c}{dr_c} \Big|_{r_c=a_c}$$

Electromagnetic channels [after Newton & Hale]

■ One-photon sector of Fock space

→ Photon 'wave function'

$$\mathbf{A}_{\mathbf{k}}(\mathbf{r}) = \left(\frac{2}{\pi \hbar c} \right)^{1/2} \sum_{jm} i^j \sum_{\lambda', \lambda=e, m, 0} \mathbf{Y}_{jm}^{(\lambda')}(\hat{\mathbf{r}}) u_{\lambda' \lambda}^j(r) \mathbf{Y}_{jm}^{(\lambda)}(\hat{\mathbf{k}}) \cdot \chi$$

→ Radial part

$$\begin{aligned} u_{ee}^j &= -[f_j'(\rho) + t_{ee}^j h_j^+(\rho)] & u_{0e}^j &= -\frac{\sqrt{j(j+1)}}{\rho} [f_j(\rho) + t_{e0}^j h_j^+(\rho)] \\ u_{mm}^j &= [f_j(\rho) + t_{mm}^j h_j^+(\rho)] & u_{0m}^j &= u_{me}^j = u_{em}^j = 0 \end{aligned}$$

→ Photon channel surface functions

$$(\mathbf{r}_c|c) = \left(\frac{\hbar c}{2\rho_\gamma} \right)^{1/2} \frac{\delta(r_\gamma - a_\gamma)}{r_\gamma} \left[\phi_{s\nu} \otimes \mathbf{Y}_{jm}^{(e,m)}(\hat{\mathbf{r}}_\gamma) \right]_{JM}$$

• Photon 'mass': $\hbar k_\gamma / c$

→ R-matrix definition preserved

$$(c'|\psi) = \sum_c R_{c'c}^B (c| \frac{\partial}{\partial r_c} r_c - B_c | \psi)$$

$$\begin{aligned} \mathbf{T} &= \rho^{1/2} \mathbf{O}^{-1} \mathbf{R}_L \mathbf{O}^{-1} \rho^{1/2} - \mathbf{F} \mathbf{O}^{-1} \\ \mathbf{R}_L &= [\mathbf{R}_B^{-1} - \mathbf{L} + \mathbf{B}]^{-1} \\ \mathbf{L} &= \rho \mathbf{O}' \mathbf{O}^{-1} \\ \mathbf{F} &= \text{Im } \mathbf{O} \end{aligned}$$

Implementation in EDA

- **EDA = Energy Dependent Analysis**

→ Adjust E_λ & $\gamma_{c\lambda}$

- **Any number of two-body channels**

→ Arbitrary spins, masses, charges (incl. mass zero)

- **Scattering observables**

→ Wolfenstein trace formalism

- **Data**

→ Normalization

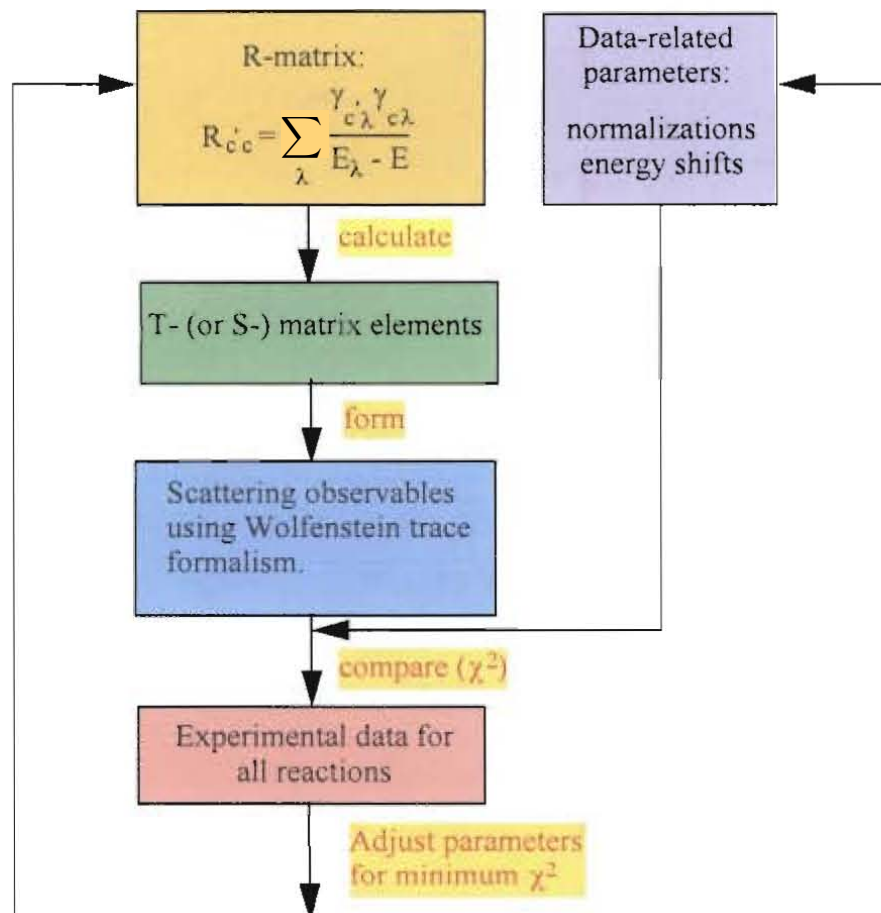
→ Energy shifts

→ Energy resolution/spread

- **Fit (rank-1 var. metric) solution**

$$\chi_{EDA}^2 = \sum_i \left[\frac{nX_i(\mathbf{p}) - R_i}{\delta R_i} \right]^2 + \left[\frac{nS - 1}{\delta S/S} \right]^2$$

- **Covariance determined**



^{17}O analysis configuration

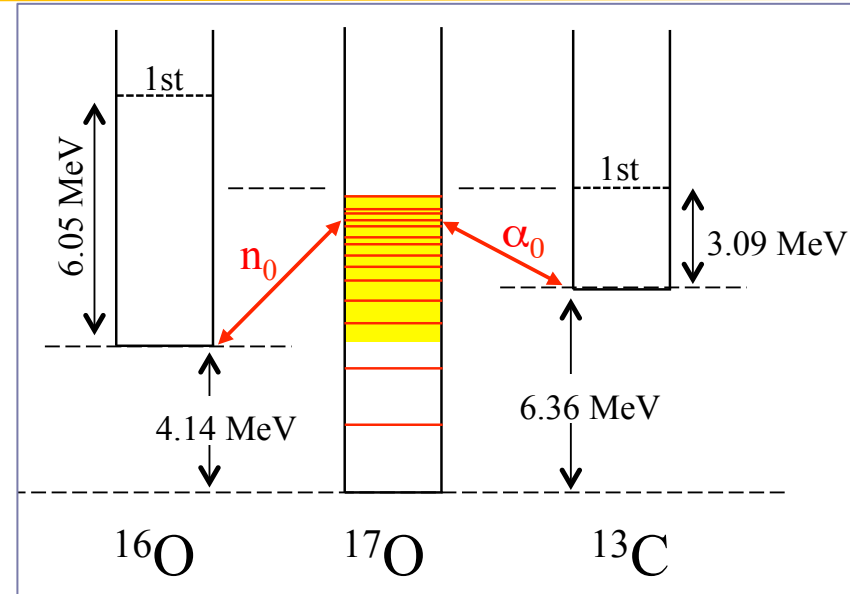
Channel	a_c (fm)	I_{max}
$n+^{16}\text{O}$	4.3	4
$\alpha+^{13}\text{C}$	5.4	5

Reaction	Energies (MeV)	# data points	Data types
$^{16}\text{O}(n,n)^{16}\text{O}$	$E_n = 0 - 7$	2718	$\sigma_T, \sigma(\theta), P_n(\theta)$
$^{16}\text{O}(n,\alpha)^{13}\text{C}$	$E_n = 2.35 - 5$	850	$\sigma_{\text{int}}, \sigma(\theta), A_n(\theta)$
$^{13}\text{C}(\alpha,n)^{16}\text{O}$	$E_\alpha = 0 - 5.4$	874	σ_{int}
$^{13}\text{C}(\alpha,\alpha)^{13}\text{C}$	$E_\alpha = 2 - 5.7$	1296	$\sigma(\theta)$
total		5738	8

R-matrix analyses support B&H73/Heil08

■ LANL R-matrix fit to Bair & Haas '73

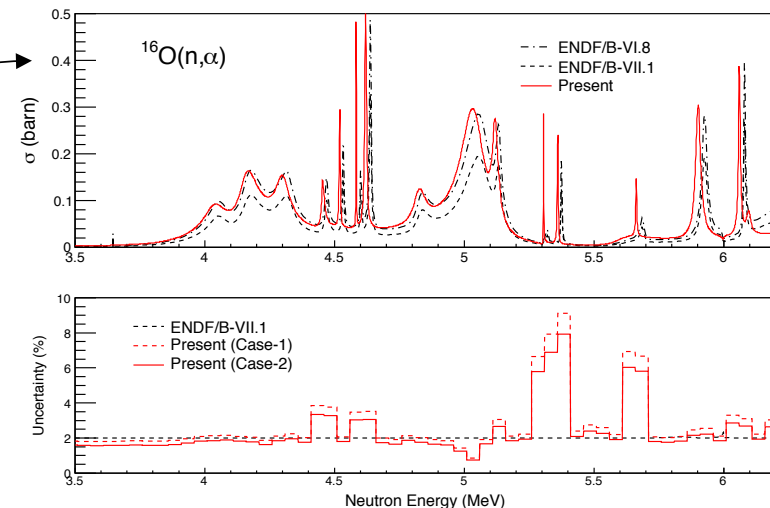
- two-channel fit: ($^{16}\text{O}, n$) & ($^{13}\text{C}, \alpha$)
 - $\ell_n = 0, \dots, 4$; $\ell_\alpha = 0, \dots, 5$
- data included: $\sigma_T(E)$
 - $^{16}\text{O}(n, n)$, $^{16}\text{O}(n, \alpha)$, $^{13}\text{C}(\alpha, n)$
 - σ_T , $\sigma(\theta)$, $P_n(\theta)$, σ_{int}
 - χ^2 minimization: normalizations float
- Test Hariss05 data
 - remove B&H73/Heil08 data
 - fix Hariss05 norm to unity
 - unable to obtain fit with realistic χ^2 (< 2.0)
 - due to tension with σ_T
 - now allow Hariss05 norm to float
 - requires scale factor of ~ 1.5 , consistent with B&H73



Credit: S. Kunieda

■ Kunieda/Kawano analysis [ND2013]

- similar to LANL R-matrix(EDA)/ENDF/B-VI.8
- with independent R-matrix code
- KK give uncertainty analysis: see ND2013 proceedings in *Nucl. Data Sh.*
- **Right to conclude B&H73 data correct on the basis of unitarity.**



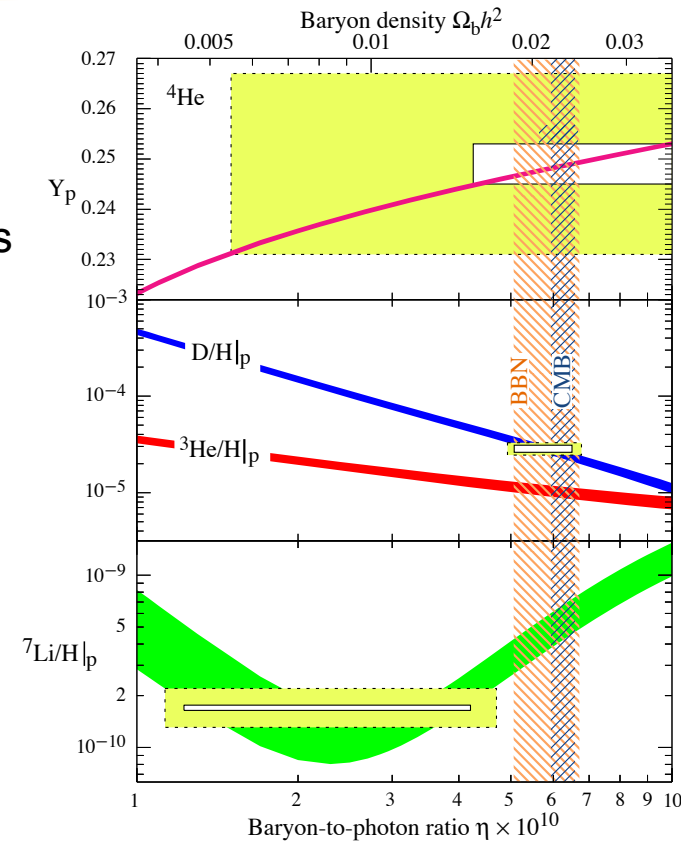
A nuclear physics solution to the BBN ${}^7\text{Li}$ problem?

■ Primordial nucleosynthesis

- Probes physics of early universe
- Big-bang nucleosynthesis: D , ${}^4\text{He}$, ${}^7\text{Li}$ abundances
- D , ${}^4\text{He}$ abundances agree with theo/expl uncertainties
- At η_{Wmap} (CMB) ${}^7\text{Li}/\text{H}|_{\text{BBN}} \sim (2.2\text{--}4.2) \cdot {}^7\text{Li}/\text{H}|_{\text{halo}}$ *
- Discrepancy $\sim 4.5\text{--}5.5\sigma$ → the “Li problem”

■ Resonant destruction ${}^7\text{Li}$

- Prod. mass 7 “well understood”; destruction not
- Cyburt & Pospelov *arXiv:0906.4373; IJMPE, 21(2012)*
 - ${}^7\text{Be}(d,p)\alpha\alpha$ & ${}^7\text{Be}(d,\gamma){}^9\text{B}$ resonant enhancement
 - Identify ${}^9\text{B}$ $E_{5/2+} \approx 16.7 \text{ MeV} \approx E_{\text{thr}}(d+{}^7\text{Be}) + 200 \text{ keV}$
 - *Near threshold*
 - $(E_r, \Gamma_d) \approx (170\text{--}220, 10\text{--}40) \text{ keV}$ solve Li problem
- Chakraborty, Fields & Olive *PRD83, 063006 (2011)*
 - More general approach: $A=8,9,10$ & 11
 - Identify as possibly important: ${}^9\text{B}$, ${}^{10}\text{B}$, ${}^{10}\text{C}$
- ‘Large’ widths
 - Both conclude “large channel radius” required



NB: both approaches assume validity of TUNL-NDG tables

^9B analysis: included data

- **$^6\text{Li}+^3\text{He}$ elastic** *Buzhinski et.al., Izv. Rossiiskoi Akademii Nauk, Ser.Fiz., Vol.43, p.158 (1979)*
 - Differential cross section
 - $1.30 \text{ MeV} < E(^3\text{He}) < 1.97 \text{ MeV}$
- **$^6\text{Li}+^3\text{He} \rightarrow \text{p}+^8\text{Be}^*$** *Elwyn et.al., Phys. Rev. C 22, 1406 (1980)*
 - Integrated cross section
 - Quasi-two-body, excited-state, summed final channel
 - $0.66 \text{ MeV} < E(^3\text{He}) < 5.00 \text{ MeV}$
- **$^6\text{Li}+^3\text{He} \rightarrow \text{d}+^7\text{Be}$** *D.W. Barr & J.S. Gilmore, unpublished (1965)*
 - Integrated cross section
 - $0.42 \text{ MeV} < E(^3\text{He}) < 4.94 \text{ MeV}$
- **$^6\text{Li}+^3\text{He} \rightarrow \gamma+^9\text{B}$** *Aleksic & Popic, Fizika 10, 273-278 (1978)*
 - Integrated cross section
 - $0.7 \text{ MeV} < E(^3\text{He}) < 0.825 \text{ MeV}$
 - New to ^9B analysis
- **New evaluation**
 - Separate $^8\text{Be}^*$ states
 - 2^+ @200 keV [16.9 MeV], 1^+ @650 keV [17.6 MeV], 1^+ @1.1 MeV [18.2 MeV]
 - $\text{n}+^8\text{B}$: $E_{\text{thresh}}(^3\text{He}) = 3 \text{ MeV}$
 - Simultaneous analysis with ^9Be mirror system

Data accessed via
EXFOR/CSISRS
database (C4 format)

R-matrix configuration in EDA code

Hadronic channels (in blue, not included)

$A_1 A_2^\pi$	${}^3\text{He}{}^6\text{Li}^+(1)$		$p{}^8\text{Be}^{*+}(2)$		$d{}^7\text{Be}^-(3)$		
$\ell \backslash S$	$\frac{3}{2}$	$\frac{1}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{1}{2}$
0	${}^4S_{3/2}$	${}^2S_{1/2}$	${}^6S_{5/2}$	${}^4S_{3/2}$	${}^6S_{5/2}$	${}^4S_{3/2}$	${}^2S_{1/2}$
1	${}^4P_{5/2,3/2,1/2}$	${}^2P_{3/2,1/2}$	${}^6P_{7/2,5/2,3/2}$	${}^4P_{5/2,3/2,1/2}$	${}^6P_{7/2,5/2,3/2}$	${}^4P_{5/2,3/2,1/2}$	${}^2P_{3/2,1/2}$
2	${}^4D_{7/2,5/2,3/2,1/2}$	${}^2D_{5/2,3/2}$	${}^6D_{9/2,7/2,5/2,3/2,1/2}$	${}^4D_{7/2,5/2,3/2,1/2}$	${}^6D_{9/2,7/2,5/2,3/2,1/2}$	${}^4D_{7/2,5/2,3/2,1/2}$	${}^2D_{5/2,3/2}$
$E_{\text{thr}}(\text{CM, MeV})$	16.6		16.7		16.5		

Electromagnetic channel: $\gamma + {}^9\text{B} \rightarrow E_1^{3/2}, M_1^{5/2}, M_1^{3/2}, M_1^{1/2}, E_1^{5/2}, E_1^{1/2}$

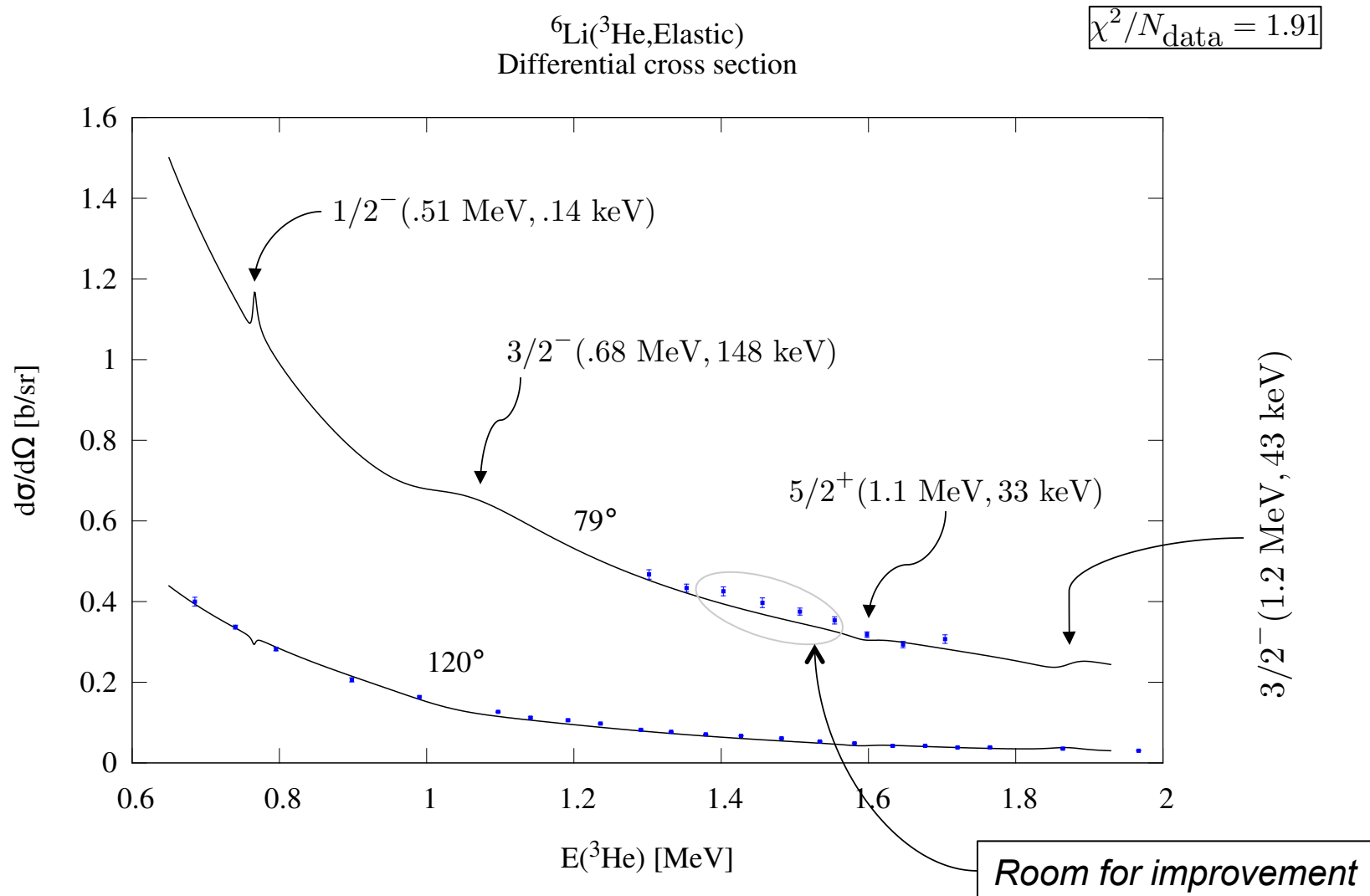
Full model space:
state number;
channel pair;
LS; J; channel
radius [fm]

1	1 4s 3/2	7.50000000f	20	1 4p 1/2	7.50000000f
2	1 4d 3/2	7.50000000f	21	1 2p 1/2	7.50000000f
3	1 2d 3/2	7.50000000f	22	2 4p 1/2	5.50000000f
4	2 4s 3/2	5.50000000f	23	3 2s 1/2	7.00000000f
5	3 6p 3/2	7.00000000f	24	4 M1 1/2	50.00000000f
6	3 4p 3/2	7.00000000f	25	1 4d 7/2	7.50000000f
7	3 2p 3/2	7.00000000f	26	3 6p 7/2	7.00000000f
8	4 E1 3/2	50.00000000f	27	1 4d 5/2	7.50000000f
9	1 4p 5/2	7.50000000f	28	1 2d 5/2	7.50000000f
10	2 6p 5/2	5.50000000f	29	2 6s 5/2	5.50000000f
11	2 4p 5/2	5.50000000f	30	3 6p 5/2	7.00000000f
12	3 6s 5/2	7.00000000f	31	3 4p 5/2	7.00000000f
13	4 M1 5/2	50.00000000f	32	4 E1 5/2	50.00000000f
14	1 4p 3/2	7.50000000f	33	1 4d 1/2	7.50000000f
15	1 2p 3/2	7.50000000f	34	1 2s 1/2	7.50000000f
16	2 6p 3/2	5.50000000f	35	3 4p 1/2	7.00000000f
17	2 4p 3/2	5.50000000f	36	3 2p 1/2	7.00000000f
18	3 4s 3/2	7.00000000f	37	4 E1 1/2	50.00000000f
19	4 M1 3/2	50.00000000f	38	2 6p 7/2	5.50000000f

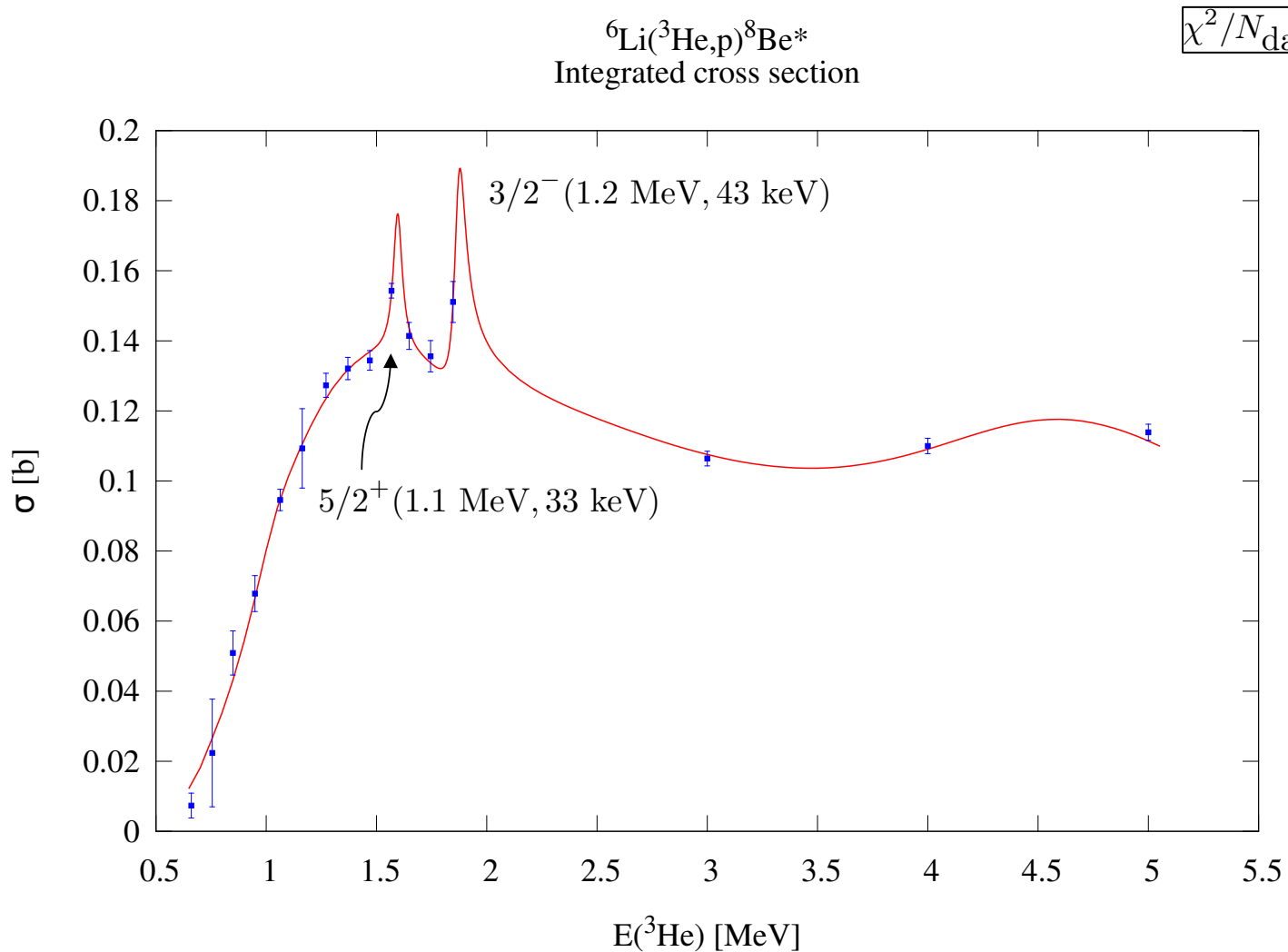
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Slide 21

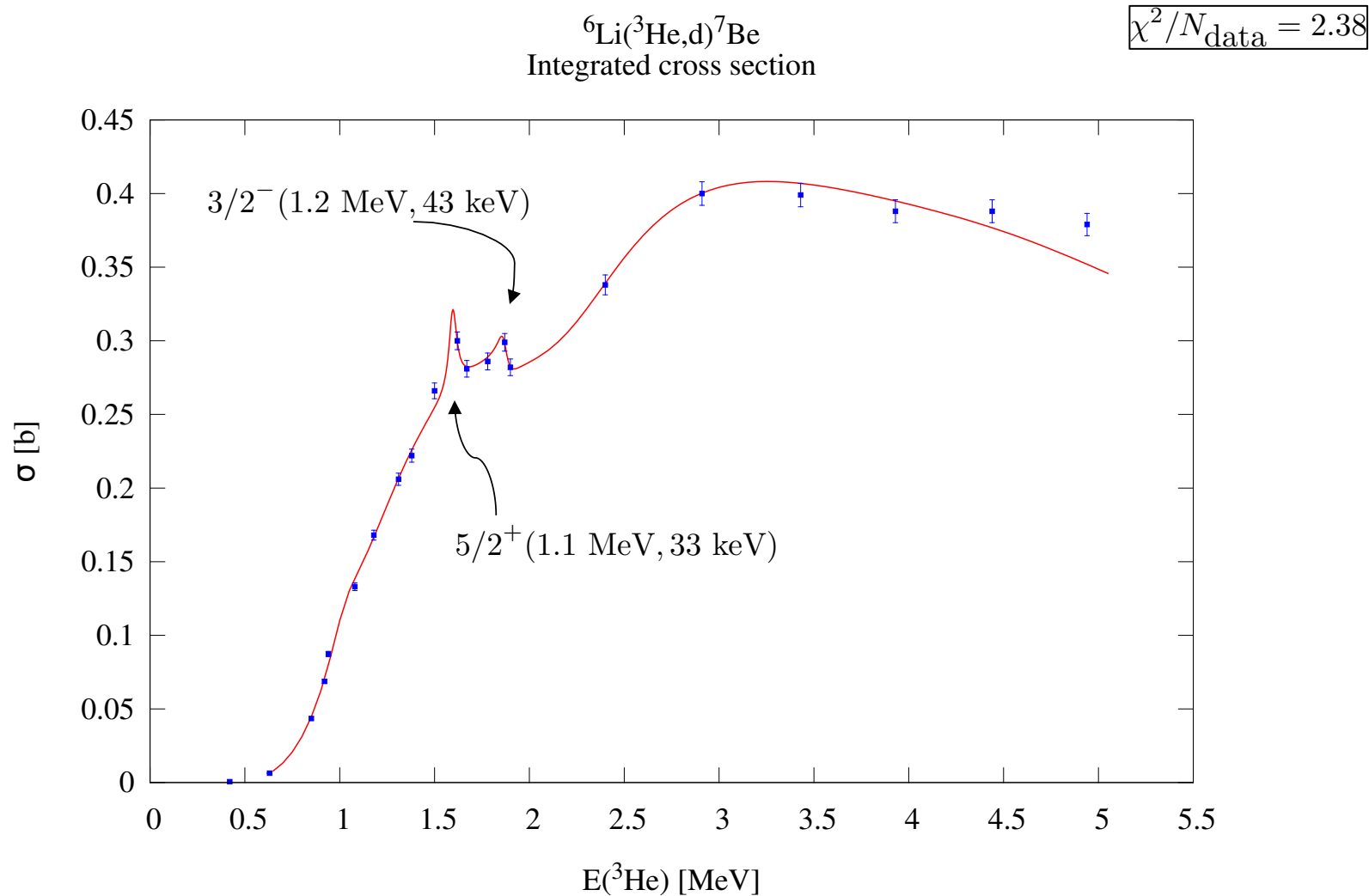
Observable fit: ${}^3\text{He}+{}^6\text{Li}$ elastic DCS



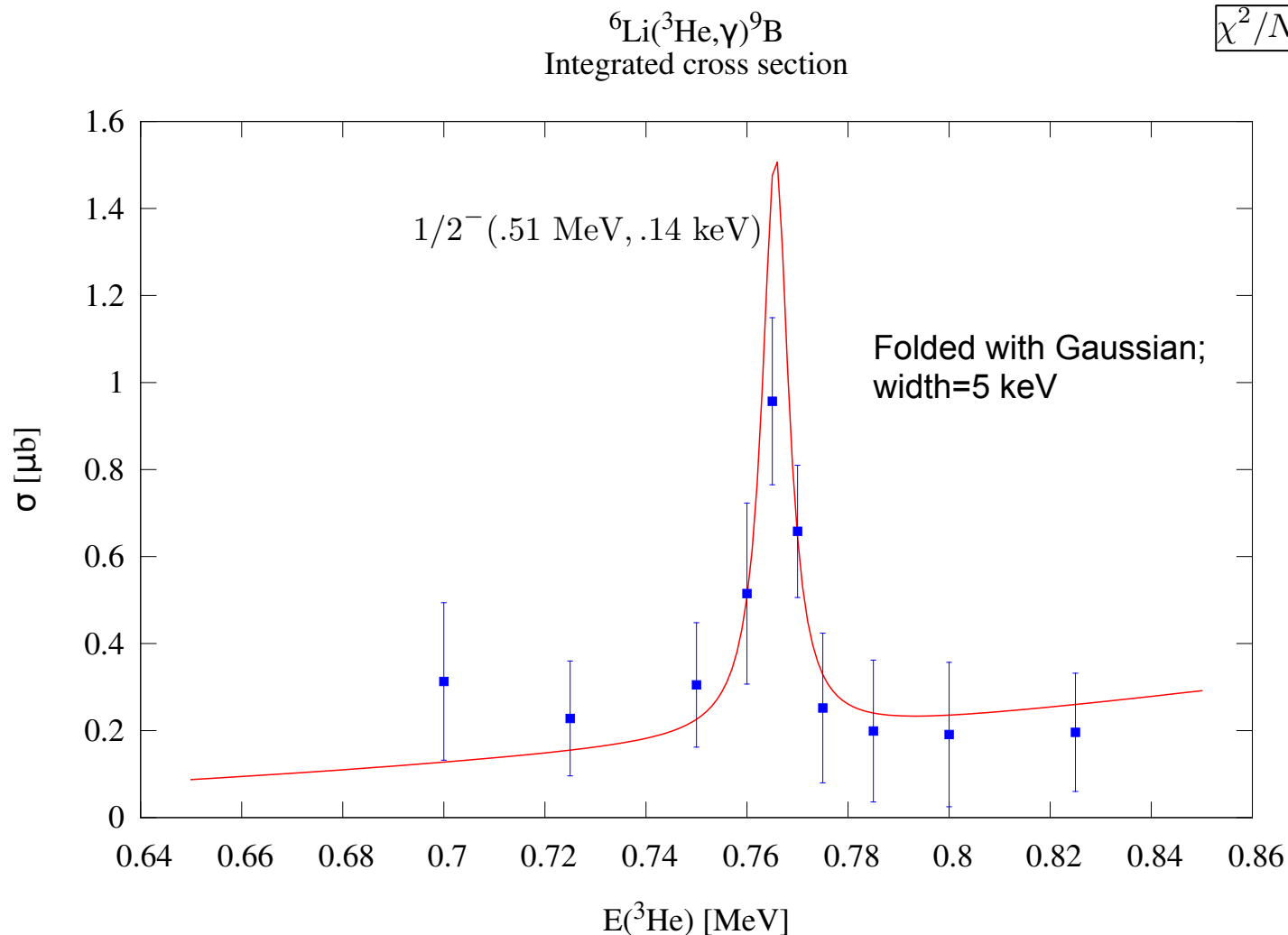
Observable fit: ${}^6\text{Li}({}^3\text{He},p){}^8\text{Be}^*$ integrated x-sec



Observable fit: ${}^6\text{Li}({}^3\text{He},d){}^7\text{Be}$ integrated x-sec



Observable fit: ${}^6\text{Li}({}^3\text{He},\gamma){}^9\text{B}$ integrated x-sec



Analysis result: resonance structure

Ex (MeV)	Jpi	Gamma (keV)	Er (MeV)	ImEr (MeV)	E (3He)	Strength
16.46539	1/2-	768.46	-.1369	-0.3842	-0.2054	0.06 weak
17.11317	1/2-	0.14	0.5109	-0.6771E-04	0.7664	1.00 strong
17.20115	5/2-	871.63	0.5989	-0.4358	0.8984	0.40 weak
17.28086	3/2-	147.78	0.6785	-0.0739	1.0178	0.77 strong
17.66538	5/2+	33.33	1.0631	-0.0167	1.5947	0.98 strong
17.83619	7/2+	2036.21	1.2339	-1.0181	1.8509	0.15 weak
17.84773	3/2-	42.52	1.2454	-0.0213	1.8681	0.97 strong
18.04821	3/2+	767.11	1.4459	-0.3836	2.1689	0.54 weak
18.42292	1/2+	5446.32	1.8206	-2.7232	2.7309	0.03 weak
18.67716	1/2-	10278.41	2.0749	-5.1392	3.1124	0.15 weak
19.60923	3/2-	1478.22	3.0069	-0.7391	4.5104	0.52 weak

S-matrix pole & residue *Hale, Brown, Jarmie PRL 59 '87*

$$\mathcal{E}_{\lambda'\lambda} = E_{\lambda} \delta_{\lambda'\lambda} - \sum_c \gamma_{c\lambda'} [L_c(E) - B_c] \gamma_{c\lambda}$$

$$E_0 = E_r - i\Gamma/2 \quad \text{residue: } i\rho_0\rho_0^T$$

**NB: no strong resonance seen
~100 keV of $^3\text{He}+^6\text{Li}$ threshold**

$$\text{strength} = \frac{1}{\Gamma} \rho_0^\dagger \rho_0 = \frac{1}{\Gamma} \sum_c \Gamma_c$$

$$\rho_{0c} = \left(\frac{2k_{0c}a_c}{N} \right)^{1/2} \mathcal{O}_c^{-1}(k_{0c}a_c) \sum_{\lambda} \gamma_{c\lambda} (\lambda|\mu_0)$$

$$N = \sum_{\lambda'\lambda} (\lambda|\mu_0)(\lambda'|\mu_0) \left[\delta_{\lambda'\lambda} + \sum_c \gamma_{c\lambda'} \frac{\partial L_c}{\partial E} \Big|_{E=E_0} \gamma_{c\lambda} \right]$$

$$L_c = r_c \frac{\partial \mathcal{O}_c}{\partial r_c} \mathcal{O}_c^{-1} \Big|_{r_c=a_c}$$

Analysis result: resonance structure

Ex (MeV)	J _{pi}	Gamma (keV)	Er (MeV)	ImEr (MeV)	E(3He)	Strength
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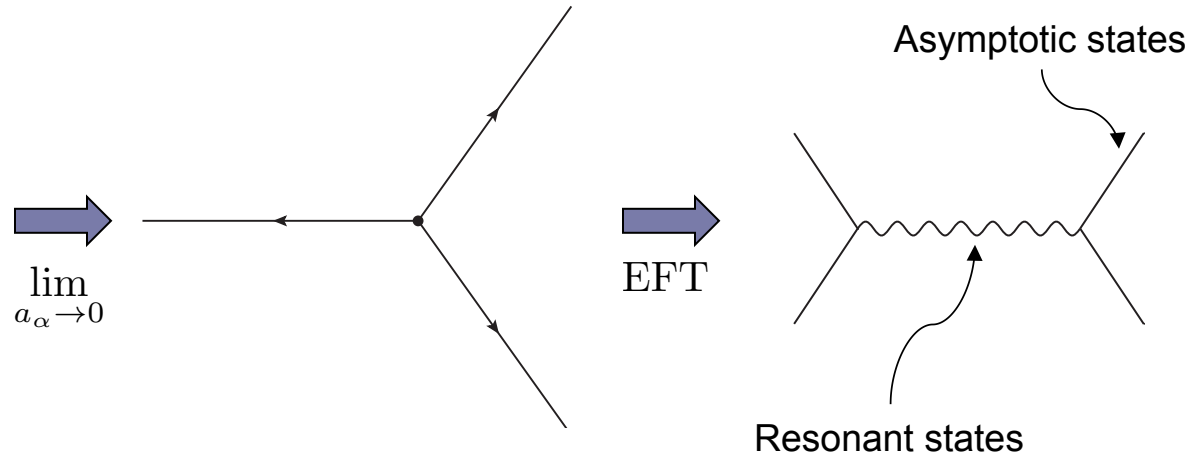
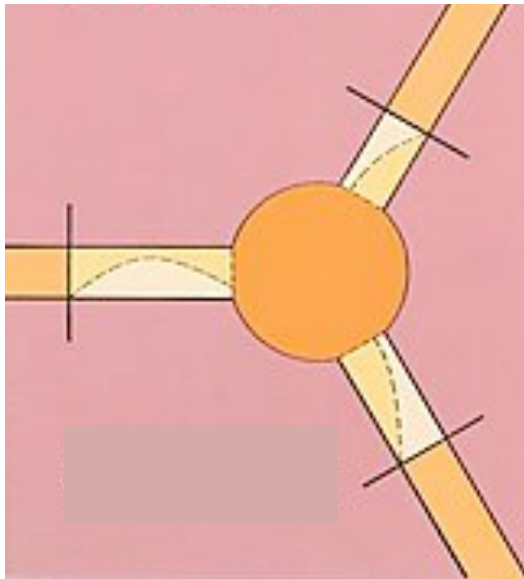
TUNL-NDG/ENSDF
parameters

**NB: no strong resonance seen
~100 keV of $^3\text{He}+^6\text{Li}$ threshold**

E_x^a (MeV \pm keV)	$J^\pi; T$	$\Gamma_{\text{c.m.}}$ (keV)	Decay
16.024 \pm 25	$T = (\frac{1}{2})$	180 \pm 16	$(\gamma, ^3\text{He})$ p, d, ^3He
16.71 \pm 100 ^h	$(\frac{5}{2}^+); (\frac{1}{2})$		
17.076 \pm 4	$\frac{1}{2}^-; \frac{3}{2}$	22 \pm 5	
17.190 \pm 25		120 \pm 40	
17.54 \pm 100 ^{h,i}	$(\frac{7}{2}^+); (\frac{1}{2})$		
17.637 \pm 10 ⁱ		71 \pm 8	p, d, $^3\text{He}, \alpha$

The zero channel radius limit of R-matrix theory?!

Brown & Hale, *Phys. Rev. C*, arXiv:1308.0348
Hale, Brown & Paris, *Phys. Rev. C*, arXiv:1308.0349



■ ‘Conventional’ R-matrix (left)

→ Consider the limit $a_\alpha \rightarrow 0$

■ Continued R-matrix theory (center)

→ Interior region: set of measure zero

■ Effective field theory* (right)

→ local Lagrangian (non-relativistic) field theory of stable and unstable particles

■ Next slides:

- Consider 1S_0 np scattering
- study a_c dependence of R-matrix parameters
- Consider $dt \rightarrow n\alpha$ R-matrix in zero channel radius limit
- Compare $dt \rightarrow n\alpha$ R-matrix to EFT

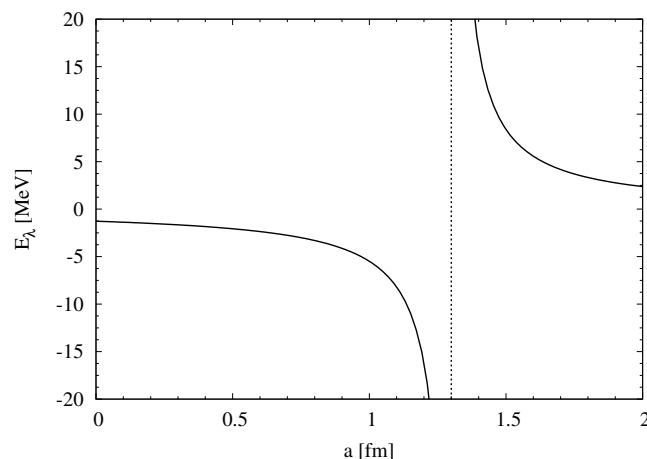
**Old fashioned EFT w/o power counting*

Zero channel radius limit: 1S_0 np scattering

$$S = e^{-2ika} \frac{1 + ikaR}{1 - ikaR} \quad R = \frac{\gamma_\lambda^2(a)}{E_\lambda(a) - E} \quad k \cot \delta = \frac{E_\lambda - E + kg^2 \tan ka}{g^2 - (E_\lambda - E) \frac{1}{k} \tan ka}$$

$$\rho = ka \quad k = \sqrt{2\mu E} \quad g^2 = a\gamma_\lambda^2$$

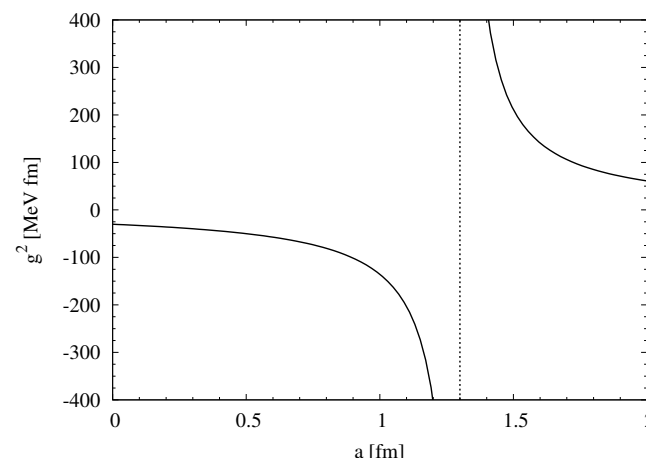
Effective range expansion: $a_0 = a - \frac{g^2}{E_\lambda} \quad r_0 = \frac{2a^3 E_\lambda^2/3 - 2a^2 E_\lambda g^2 + 2ag^4 - g^2 \hbar^2/\mu}{(g^2 - aE_\lambda)^2}$



$$E_\lambda(a) = \frac{\hbar^2(a_0 - a)}{2\mu[r_0 a_0^2/2 - a^3/3 - aa_0(a_0 - a)]}$$

Pole position: $a_p = a_0 + \left\{ a_0^3 \left[\frac{3r_0}{2a_0} - 1 \right] \right\}^{1/3}$

$$g^2(a) = (a - a_0)E_\lambda(a) = -\frac{\hbar^2(a_0 - a)^2}{2\mu[r_0 a_0^2/2 - a^3/3 - aa_0(a_0 - a)]}$$

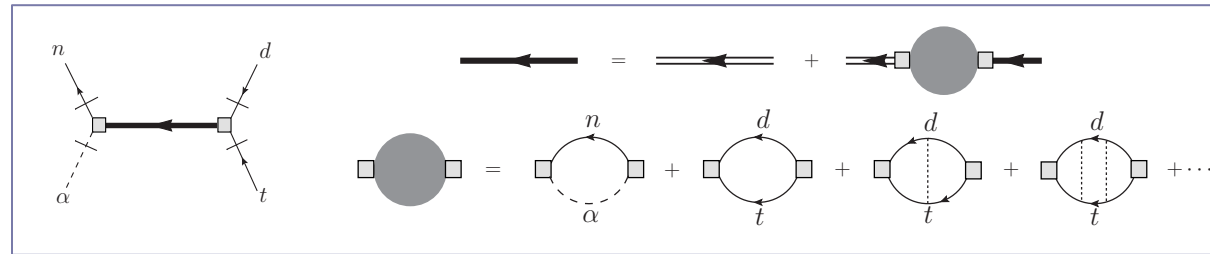


Effective field theory \leftrightarrow R matrix: $dt \rightarrow n\alpha$

Exactly soluble EFT with 'wrong-sign' free Lagrangian and DOF:

Particle	Spin	Operators	Mass	Binding
Alpha	0^+	$\phi_\alpha^\dagger(\mathbf{r}, t), \phi_\alpha(\mathbf{r}, t)$	$m_\alpha = 2m_p + 2m_n$	ϵ_α
Deuteron	1^+	$\phi_d^\dagger(\mathbf{r}, t), \phi_d(\mathbf{r}, t)$	$m_d = m_p + m_n$	ϵ_d
Neutron	$\frac{1}{2}^+$	$\psi_n^\dagger(\mathbf{r}, t), \psi_n(\mathbf{r}, t)$	m_n	$\epsilon_n \equiv 0$
Triton	$\frac{1}{2}^+$	$\psi_t^\dagger(\mathbf{r}, t), \psi_t(\mathbf{r}, t)$	$m_t = m_p + 2m_n$	ϵ_t
$^5\text{He}^*$	$\frac{3}{2}^+$	$\psi_*^\dagger(\mathbf{r}, t), \psi_*(\mathbf{r}, t)$	$m_* = 2m_p + 3m_n$	ϵ_*

$$\mathcal{L}_{\text{int}} = g_{dt} \left[\psi_*^\dagger \cdot \phi_d \psi_t + \psi_t^\dagger \phi_d^\dagger \cdot \psi_* \right] + g_{n\alpha} \left[\psi_*^\dagger \cdot \Psi_{n\alpha} + \Psi_{n\alpha}^\dagger \cdot \psi_* \right]$$



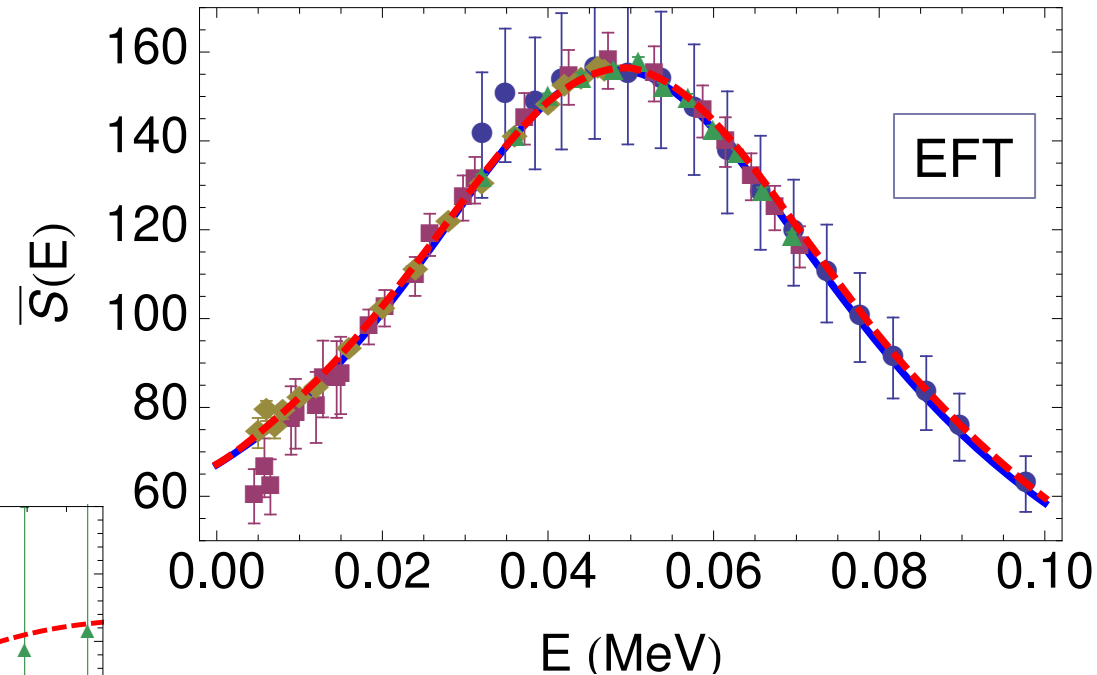
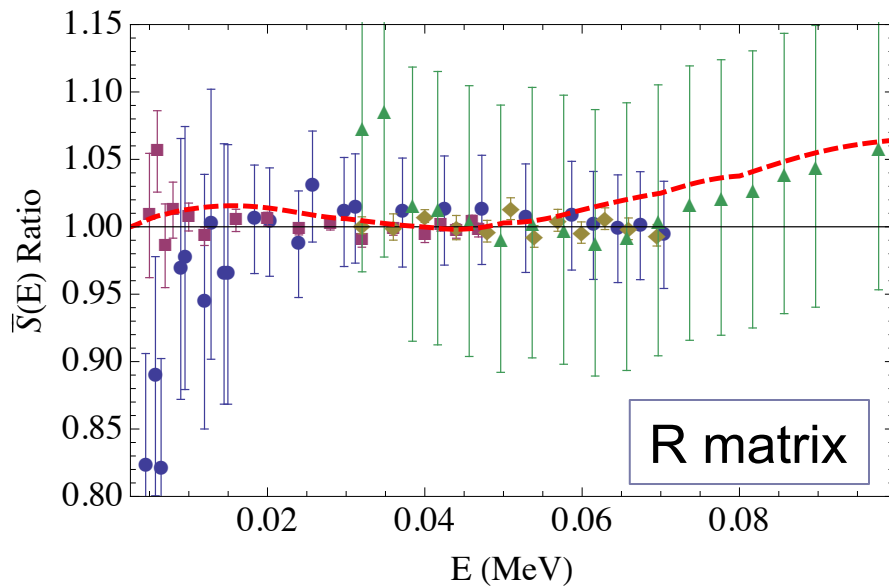
$$\sigma_{dt \rightarrow n\alpha} = \frac{8}{9} 4\pi m_{n\alpha} \frac{p_{n\alpha}^5}{v_{dt}} \frac{g_{dt}^2}{4\pi} \frac{g_{n\alpha}^2}{4\pi} \left| \psi_{\mathbf{p}_{dt}}^{(0)}(0) \right|^2 \left[\frac{p_{dt}^2}{2m_{dt}} - E_* - \frac{g_{dt}^2}{4\pi} \Delta(W) \right]^2 + \left[\frac{g_{dt}^2}{4\pi} 2m_{dt} p_{dt} \left| \psi_{\mathbf{p}_{dt}}^{(C)}(0) \right|^2 + \frac{g_{n\alpha}^2}{4\pi} \frac{2}{3} m_{n\alpha} p_{n\alpha}^5 \right]^2 \Big|^{-2}$$

Identical to R-matrix (d,t)&(n,α) in the limit $a_d, a_n \rightarrow 0$

$$\gamma_d^2 = -\frac{g_d^2}{2\pi} \frac{\mu_d}{\hbar^2 a_d} \text{ and } \gamma_n^2 = -\frac{g_n^2}{6\pi} \frac{\mu_n}{\hbar^2 a_n^5}$$

Effective field theory \leftrightarrow R matrix: $dt \rightarrow n\alpha$

- Dim'less astrophys factor vs E_d (CM)
- Solid-blue: 3 par EFT fit $\chi^2/\text{dof} \sim 0.8$
- Dashed-red: Bosch & Hale (1992)
 - 2665 data/117 pars/ $\chi^2/\text{dof} \sim 1.6$



- Four parameter fit @ finite a
- Divide out Bosch-Hale92 fit
- Unphysical reduced widths $\gamma_c^2 < 0$ $a < 2$ fm

Summary, findings & future work

Summary/findings

- Provided overview of current work in the LANL light nuclear reaction program
- Emphasize the utility of multichannel, unitary parametrization of light nuc data
 - ^{17}O norm issue: are Bair & Haas '73 data conclusive?
 - ^9B resonance spectrum: no resonances in ^9B that reside within ~ 200 (~ 100) keV of the $d+^7\text{Be}$ ($^3\text{He}+^6\text{Li}$) threshold with 'large' widths 10—40 keV
 - Appears to rule out scenarios considered by *Cybert & Pospelov (2009)* that low-lying, robust resonance in ^9B could explain the "Li problem"

Near-term, Future Work

- Complete $^{13,14}\text{C}$ analyses
- NN up to 200 MeV
- Improvements in the ^9B analysis: more channels; incorporate $p+^8\text{Be}^*$ angular data; proper treatment three-body final states
- Extend EFT—R-matrix approach to multichannel, multilevel problems

Supplementary material

Additional slides follow

Motivation

■ Cross section evaluation & resonance structure

→ *Nucl. Phys. A745, 155, 2004(2011)*

E_x^a (MeV \pm keV)	$J^\pi; T$	$\Gamma_{c.m.}$ (keV)	Decay
16.024 ± 25	$T = (\frac{1}{2})$	180 ± 16	$(\gamma, {}^3\text{He})$ $p, d, {}^3\text{He}$
16.71 ± 100^h	$(\frac{5}{2}^+); (\frac{1}{2})$		
17.076 ± 4	$\frac{1}{2}^-; \frac{3}{2}$	22 ± 5	
17.190 ± 25		120 ± 40	
$17.54 \pm 100^{h,i}$	$(\frac{7}{2}^+); (\frac{1}{2})$		
17.637 ± 10^i		71 ± 8	$p, d, {}^3\text{He}, \alpha$

■ Astrophysical applications

→ Big bang nucleosynthesis

- Nuclear physics solution to ${}^7\text{Li}$ predicted overproduction problem? (*cf. Hoyle*)
- Details next slide.

■ Purpose within Los Alamos Nat. Lab programmatic

→ Continue the R-matrix program for various end-users

→ Ongoing/upcoming analysis releases: ${}^7\text{Be}$, ${}^{13}\text{C}$ [*G. Hale Tues. Session GA 2*], ${}^{14}\text{C}$, ${}^{17}\text{O}$, ...

Implementation in EDA

■ EDA = Energy Dependent Analysis

→ Adjust E_λ & $\gamma_{c\lambda}$

■ Any number of two-body channels

→ Arbitrary spins, masses, charges (incl. mass zero)

■ Scattering observables

→ Wolfenstein trace formalism

■ Data

→ Normalization

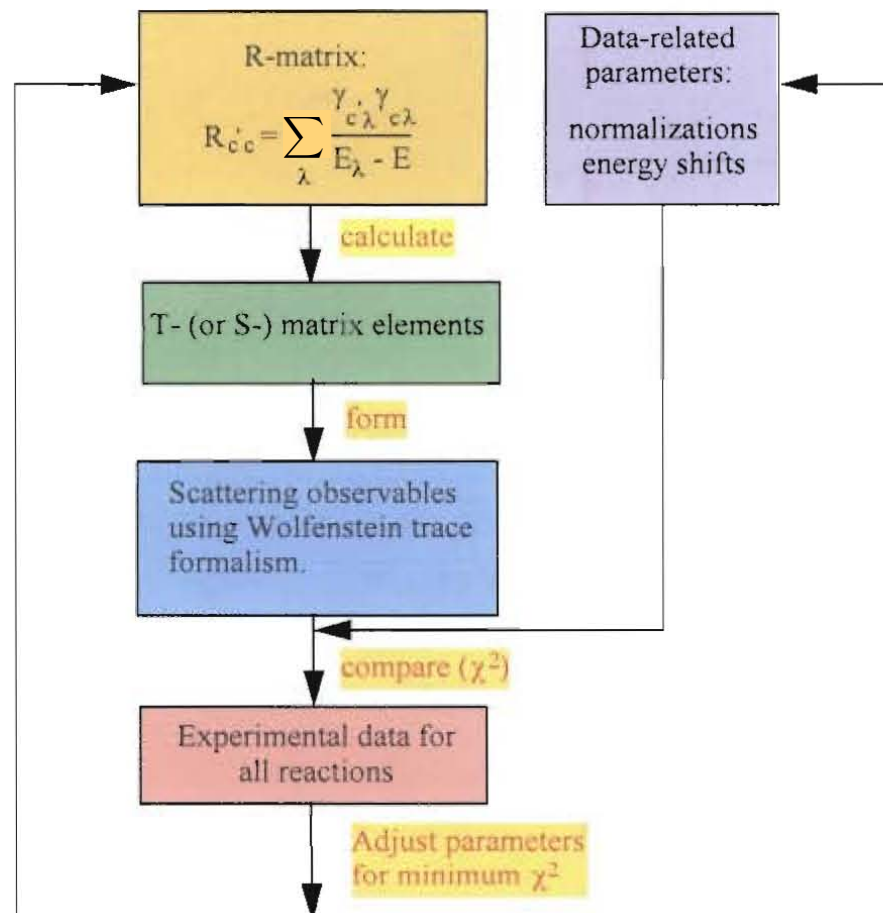
→ Energy shifts

→ Energy resolution/spread

■ Fit solution

$$\chi_{EDA}^2 = \sum_i \left[\frac{nX_i(\mathbf{p}) - R_i}{\delta R_i} \right]^2 + \left[\frac{nS - 1}{\delta S/S} \right]^2$$

■ Covariance determined



Electromagnetic channels

■ One-photon sector of Fock space

→ Photon 'wave function'

$$\mathbf{A}_{\mathbf{k}}(\mathbf{r}) = \left(\frac{2}{\pi \hbar c} \right)^{1/2} \sum_{jm} i^j \sum_{\lambda', \lambda=e, m, 0} \mathbf{Y}_{jm}^{(\lambda')}(\hat{\mathbf{r}}) u_{\lambda' \lambda}^j(r) \mathbf{Y}_{jm}^{(\lambda)}(\hat{\mathbf{k}}) \cdot \chi$$

→ Radial part

$$\begin{aligned} u_{ee}^j &= -[f_j'(\rho) + t_{ee}^j h_j^{+'}(\rho)] & u_{0e}^j &= -\frac{\sqrt{j(j+1)}}{\rho} [f_j(\rho) + t_{e0}^j h_j^+(\rho)] \\ u_{mm}^j &= [f_j(\rho) + t_{mm}^j h_j^+(\rho)] & u_{0m}^j &= u_{me}^j = u_{em}^j = 0 \end{aligned}$$

→ Photon channel surface functions

$$(\mathbf{r}_c|c) = \left(\frac{\hbar c}{2\rho_\gamma} \right)^{1/2} \frac{\delta(r_\gamma - a_\gamma)}{r_\gamma} \left[\phi_{s\nu} \otimes \mathbf{Y}_{jm}^{(e,m)}(\hat{\mathbf{r}}_\gamma) \right]_{JM}$$

• Photon 'mass': $\hbar k_\gamma / c$

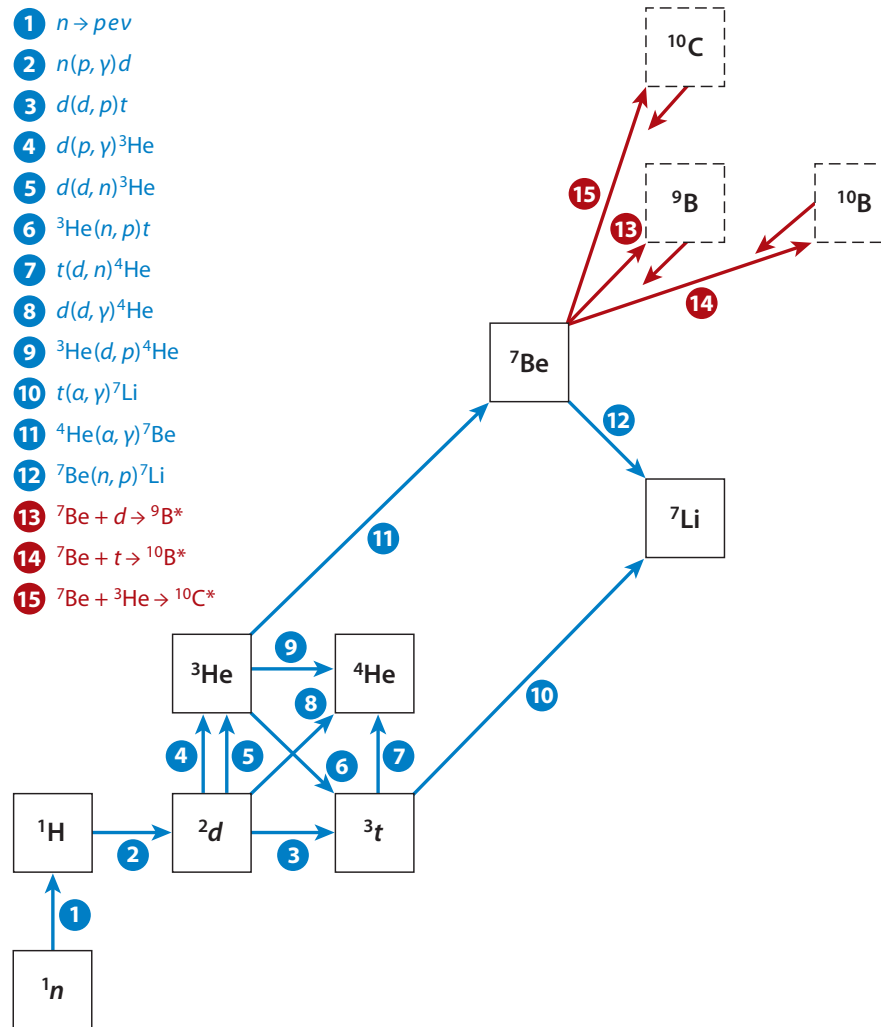
→ R-matrix definition preserved

$$(c'|\psi) = \sum_c R_{c'c}^B (c | \frac{\partial}{\partial r_c} r_c - B_c | \psi)$$

$$\begin{aligned} \mathbf{T} &= \rho^{1/2} \mathbf{O}^{-1} \mathbf{R}_L \mathbf{O}^{-1} \rho^{1/2} - \mathbf{F} \mathbf{O}^{-1} \\ \mathbf{R}_L &= [\mathbf{R}_B^{-1} - \mathbf{L} + \mathbf{B}]^{-1} \\ \mathbf{L} &= \rho \mathbf{O}' \mathbf{O}^{-1} \\ \mathbf{F} &= \text{Im } \mathbf{O} \end{aligned}$$

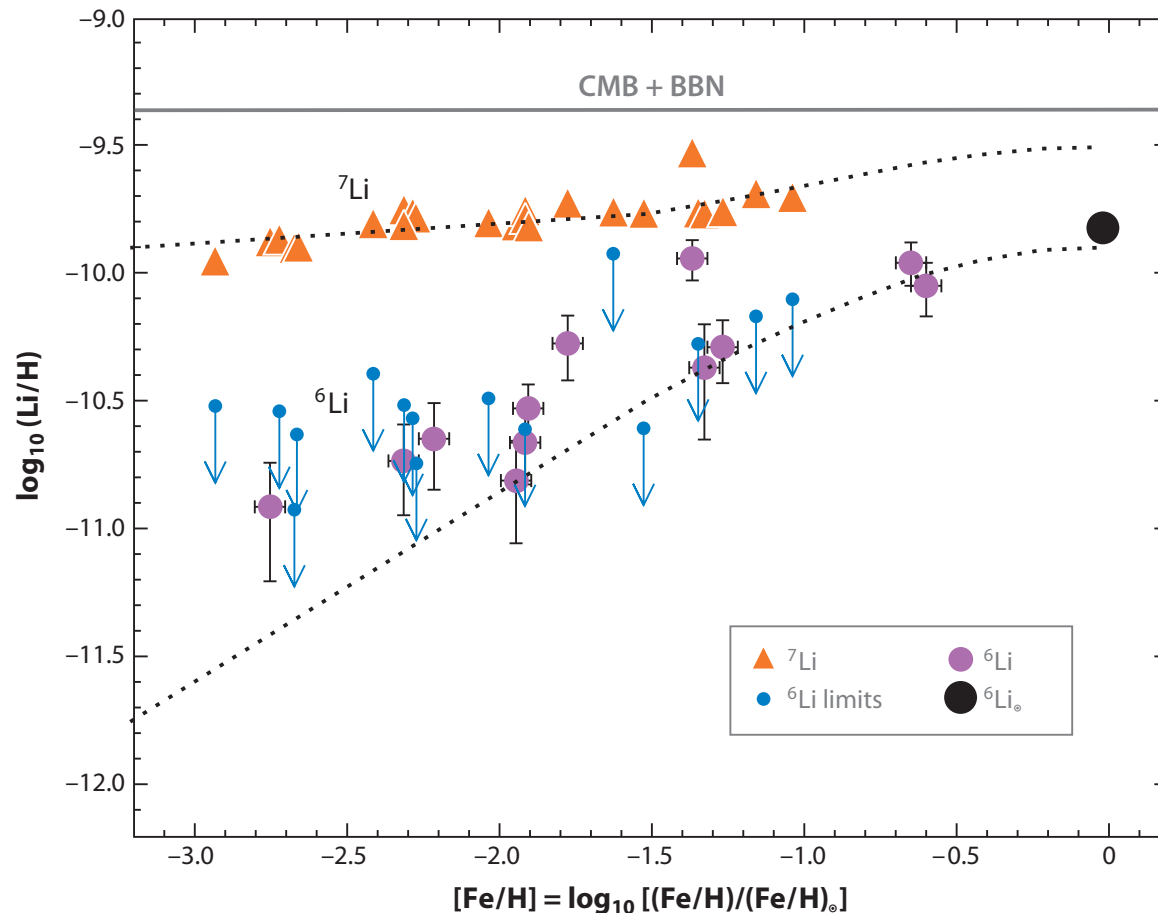
BBN reaction network (simplified)

■ **Fields** *Annu. Rev. Nucl. Part. Sci.* 2011. 61:47–68



Spite Plateau

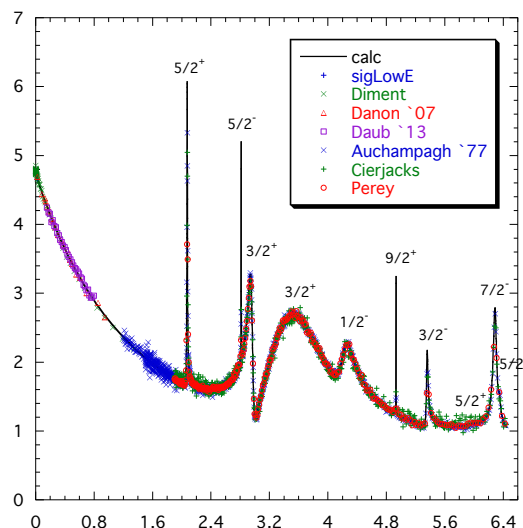
■ Measurement of primordial ${}^7\text{Li}$ from low-metallicity halo dwarf stars



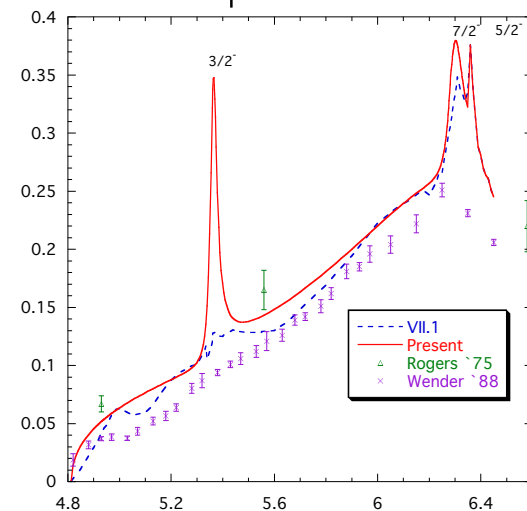
Asplund M, et al. Astrophys. J. 644:229 (2006)

$^{13,14}\text{C}$ system analyses: σ_T (b) vs. E_n (MeV)

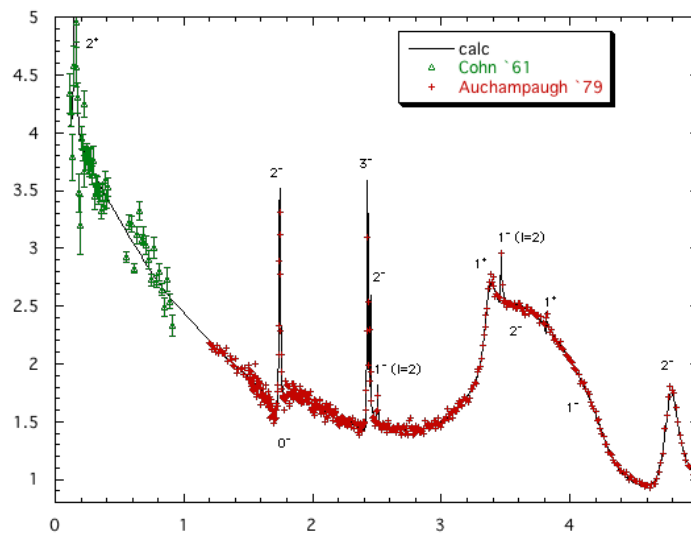
$n+^{12}\text{C}$ Total Cross Section



$^{12}\text{C}(n,n_1)^{12}\text{C}^*$ Cross Section



$n+^{13}\text{C}$ Total Cross Section



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